

# PARTICLE ACCELERATION AT RELATIVISTIC SHOCKS: AGN JET TSs & PW TSs

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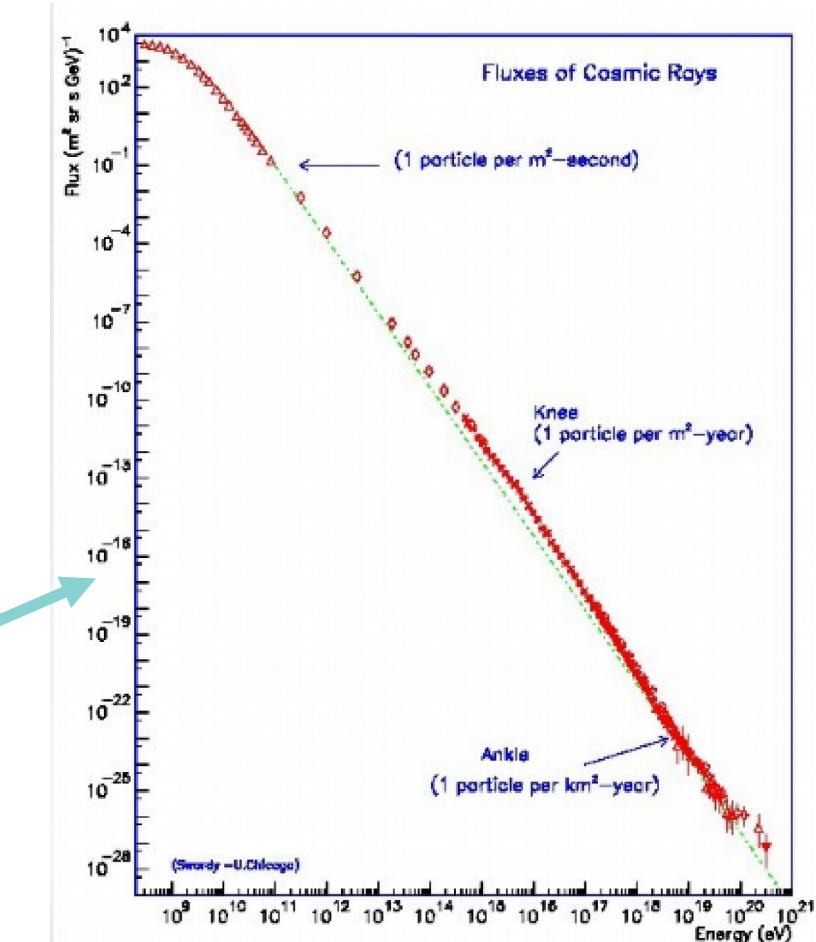
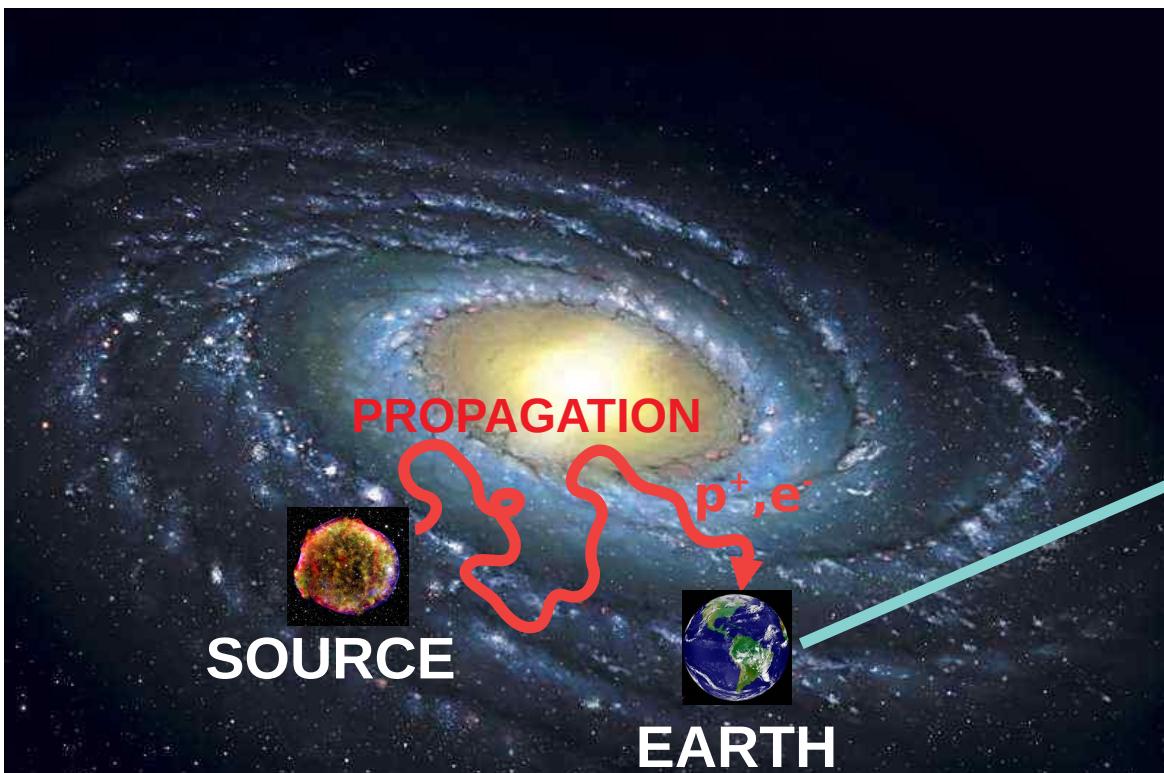
上海交通大学  
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Tsung-Dao Lee Institute

A&A 676, A23 (2023); arXiv:2303.12636

# Cosmic-rays & their secondaries :

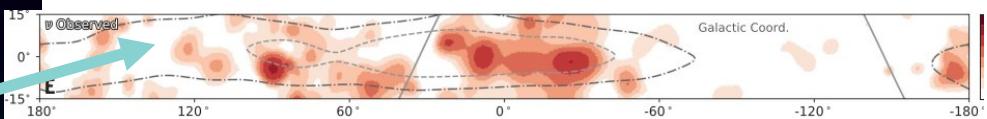
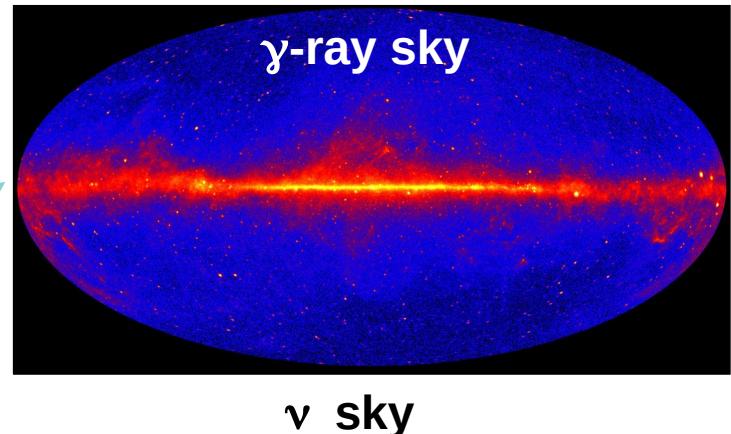
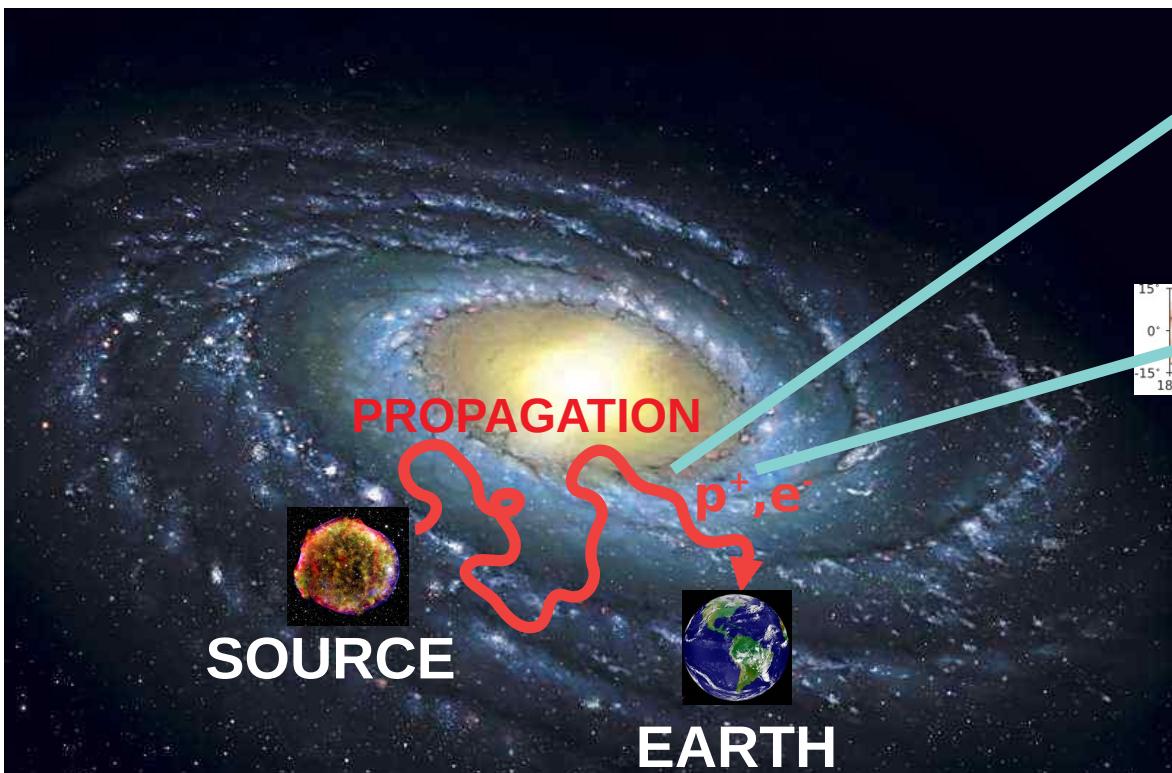
High-energy particles ( $p^+, e^-$ ,...) up to  $10^{20}$  eV.



# Cosmic-rays & their secondaries :

Produce  $\gamma$ -rays, neutrinos:

High-energy particles ( $p^+, e^-$ ,...) up to  $10^{20}$  eV.

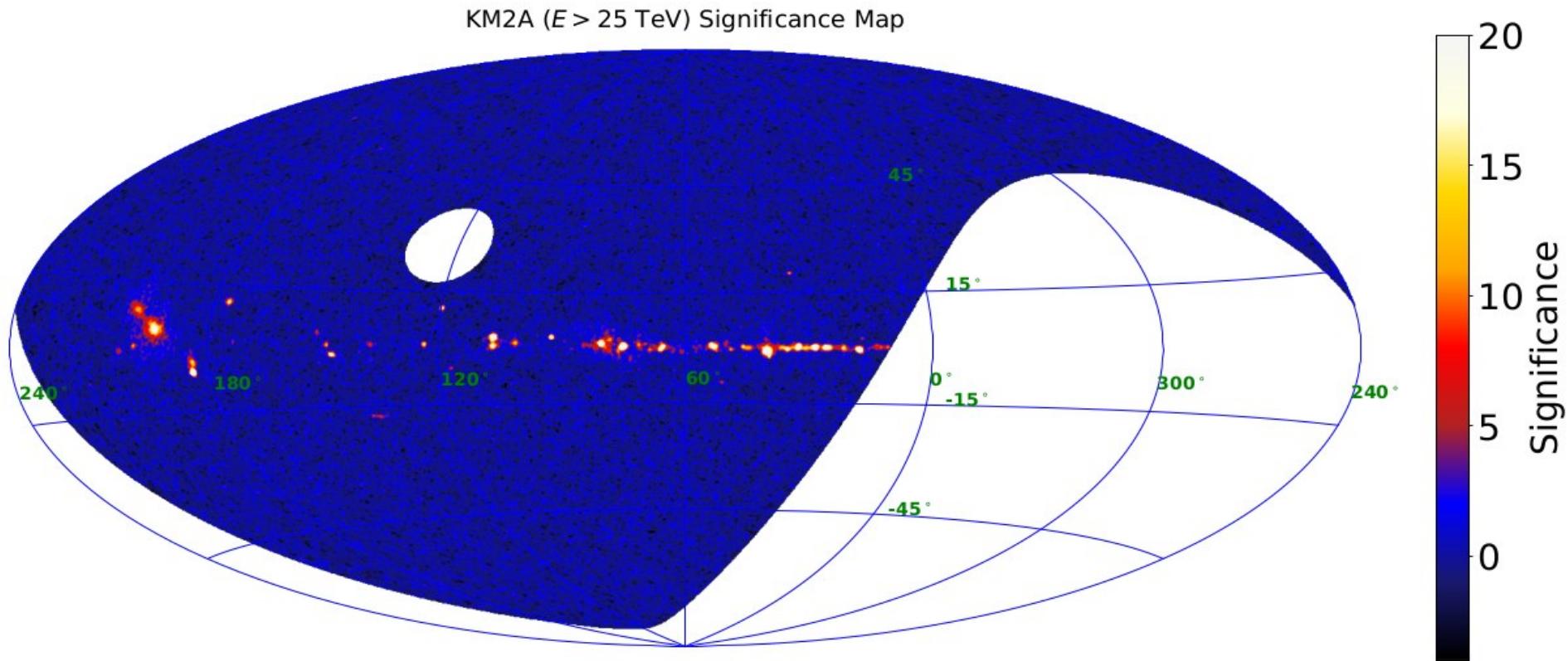


# PeV Gamma-Ray Astronomy: LHAASO



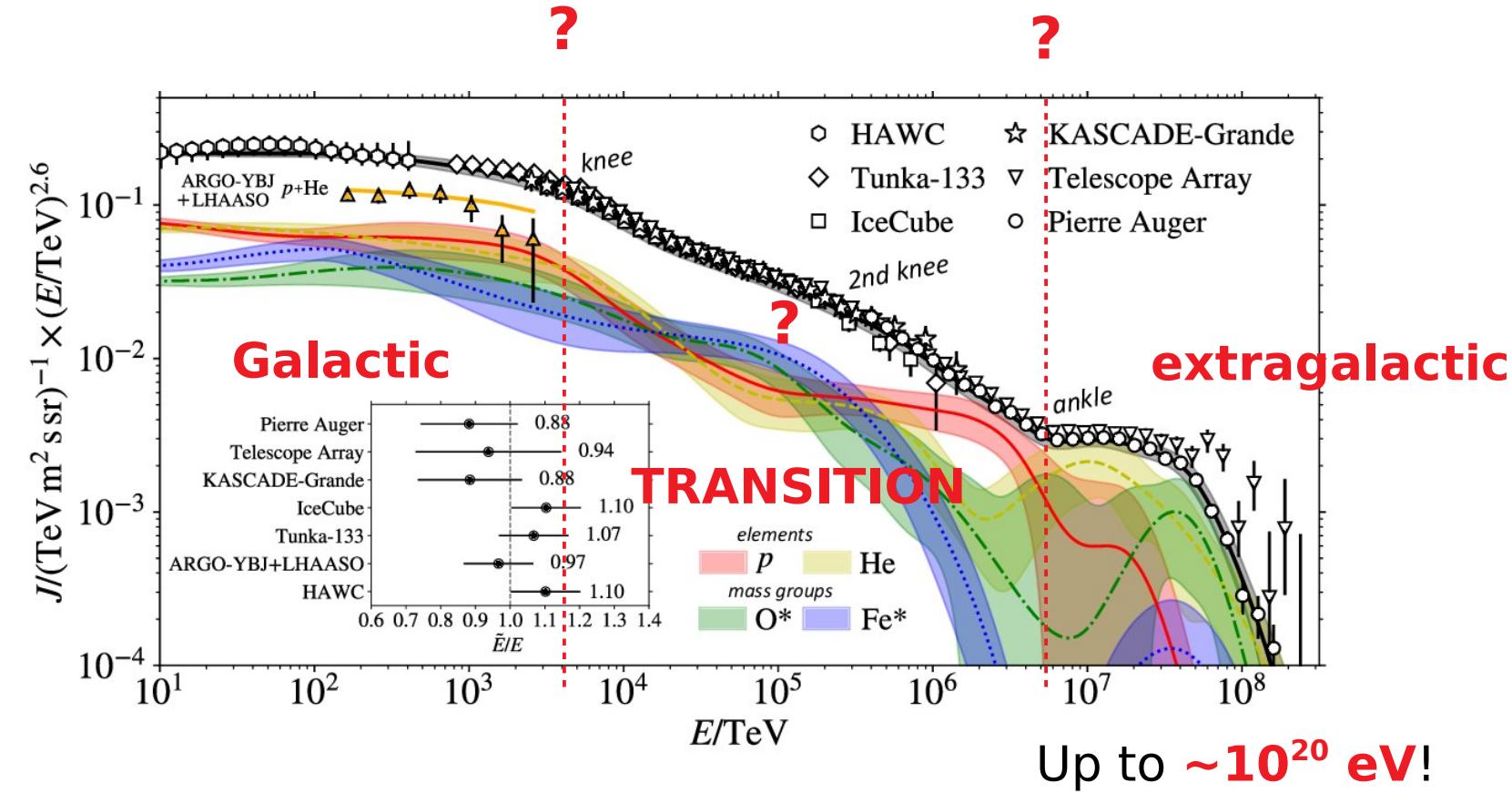
# Sky at 25 TeV – 1 PeV with LHAASO

LHAASO Collaboration

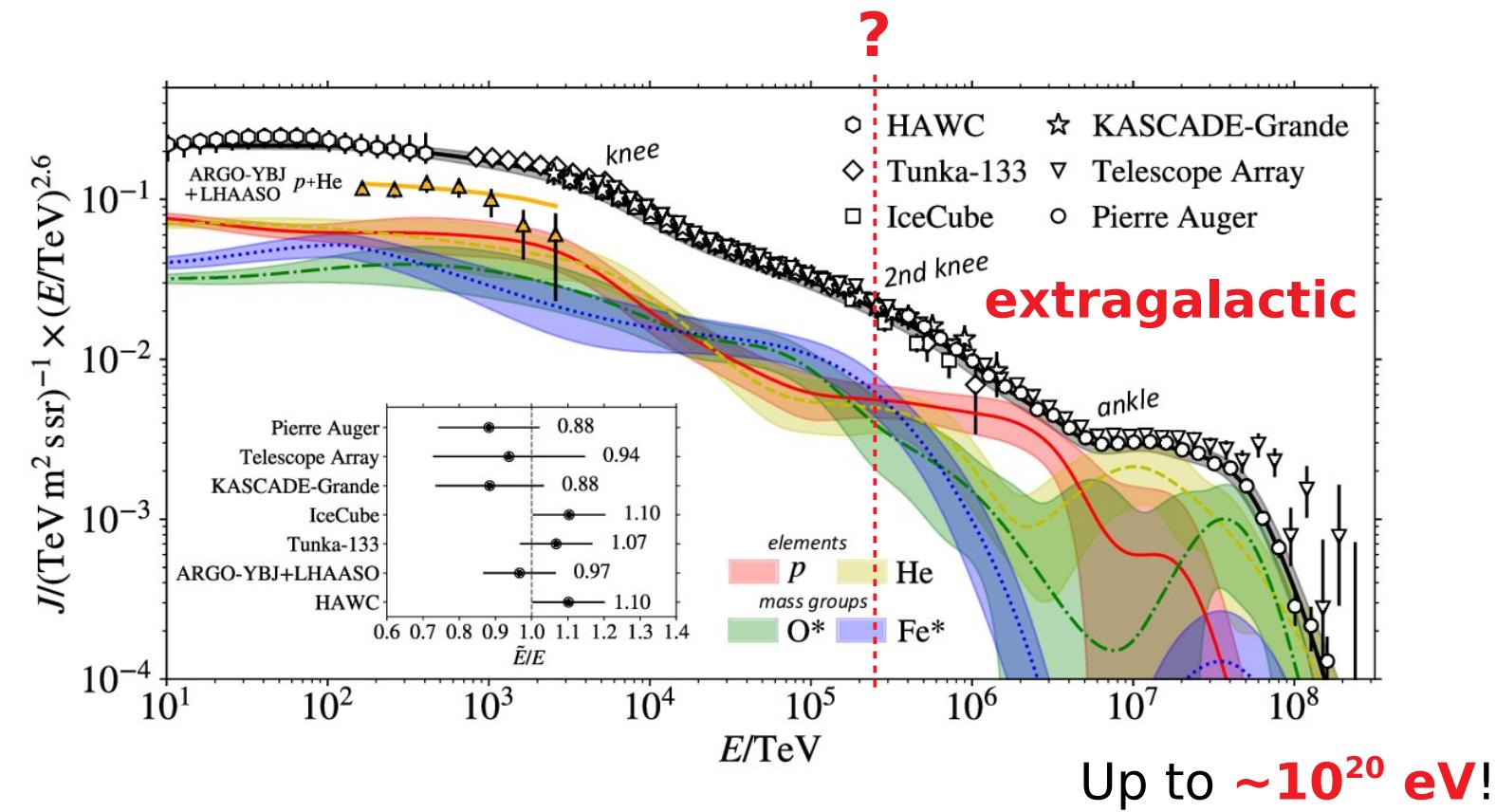


Hadronic PeVatrons: Stellar clusters ?; Many PWNe (leptonic) → Relativistic shock

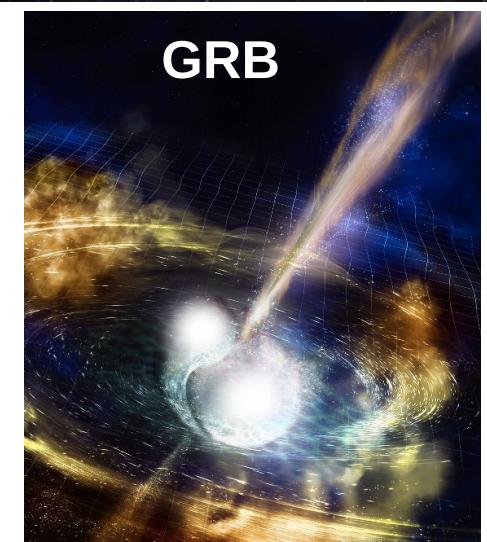
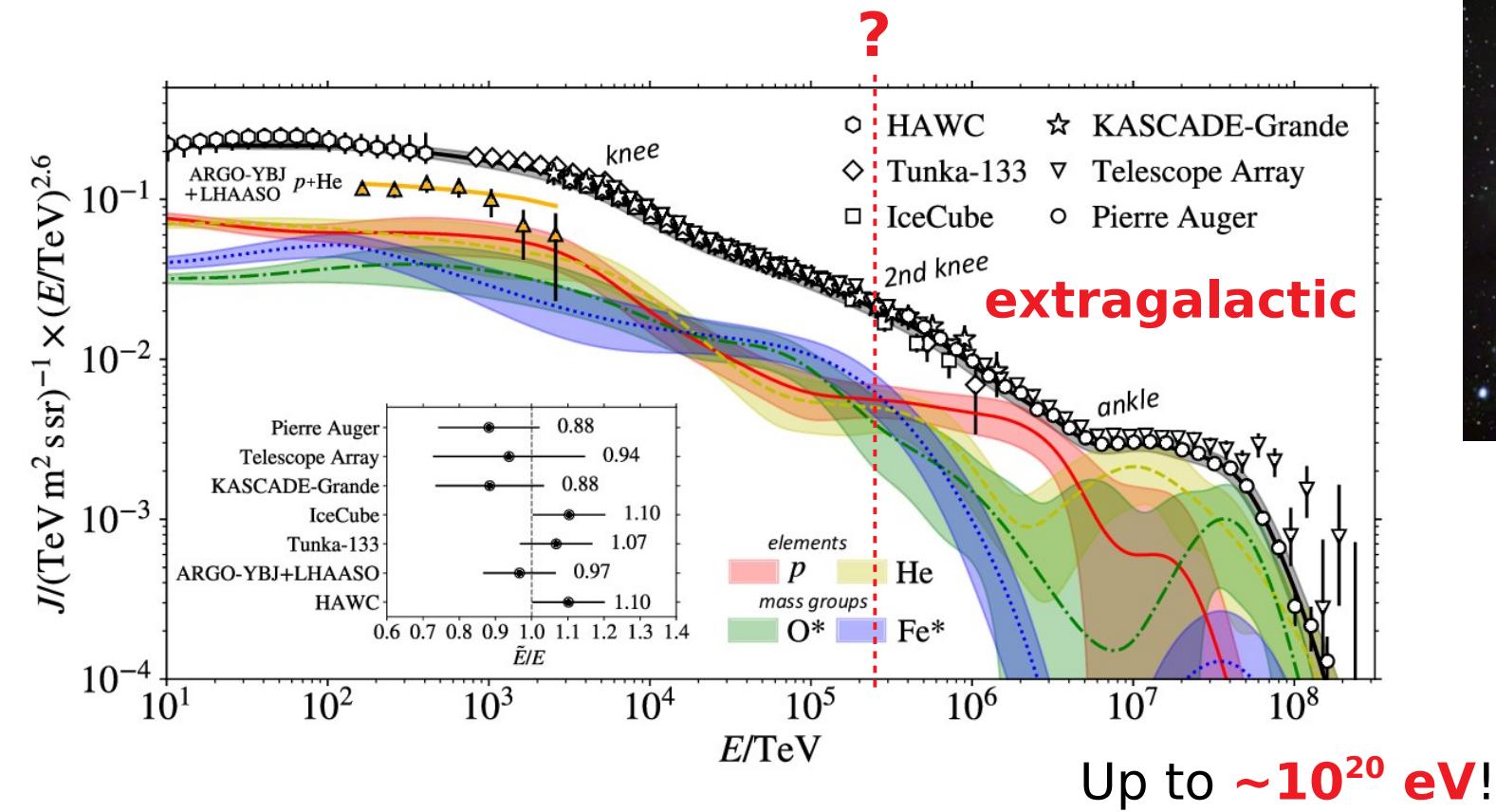
# Cosmic-Ray Spectrum



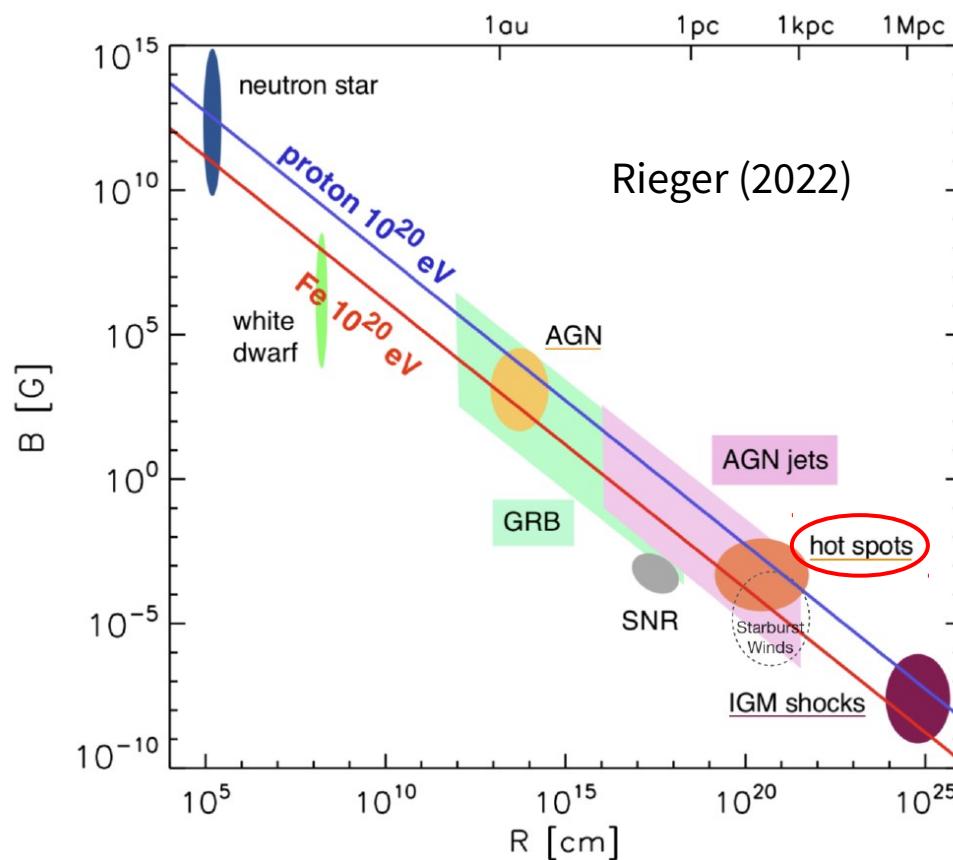
# Ultra-High Energy Cosmic-Rays



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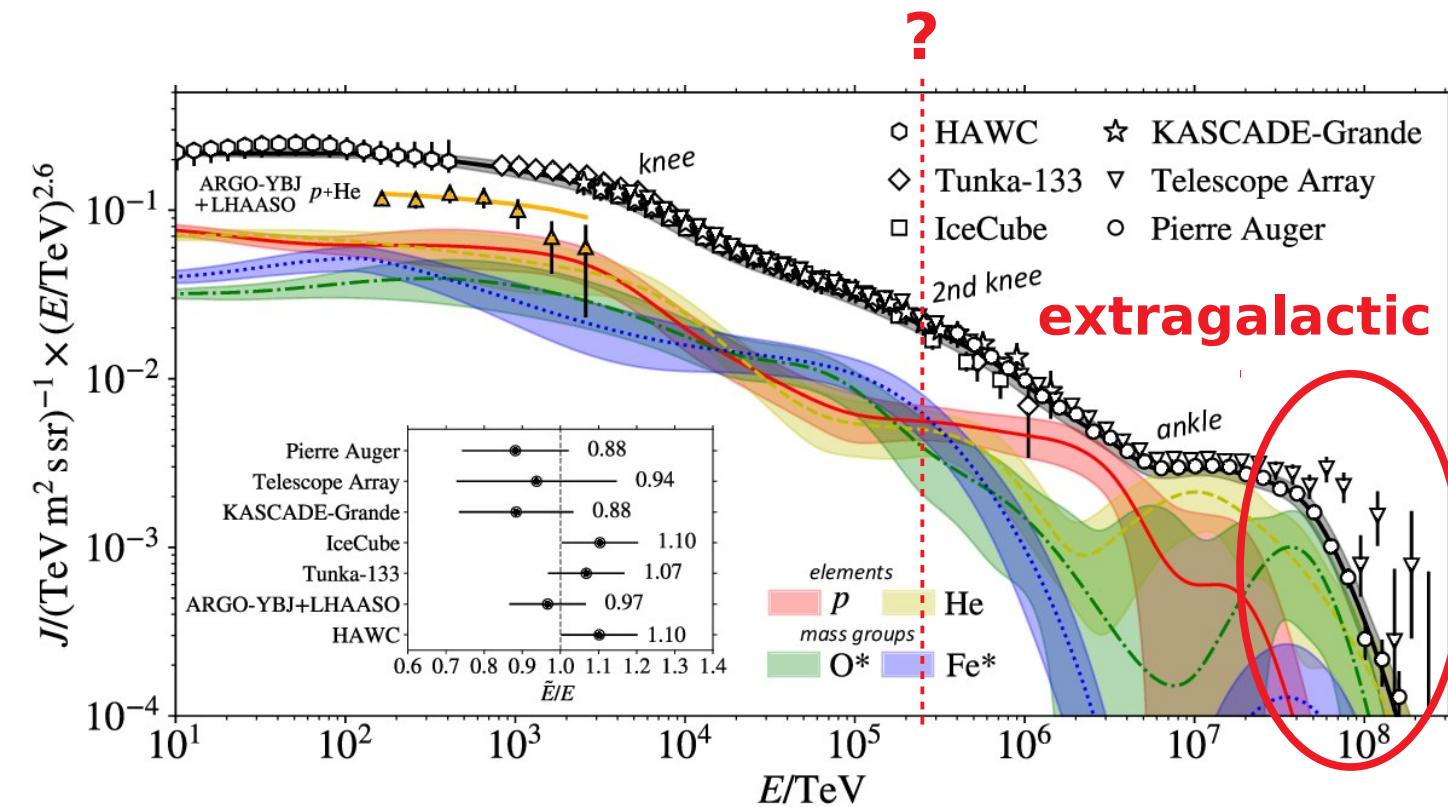
# Ultra-High Energy Cosmic-Rays



**Hillas criterion  
(Hillas 1984):**  
 $R_g < \sim \text{size source}$

$$E \leq 10^{20} Z \left( \frac{B}{10 \mu \text{G}} \right) \left( \frac{L}{10 \text{ kpc}} \right) \text{ eV}$$

# Ultra-High Energy Cosmic-Rays



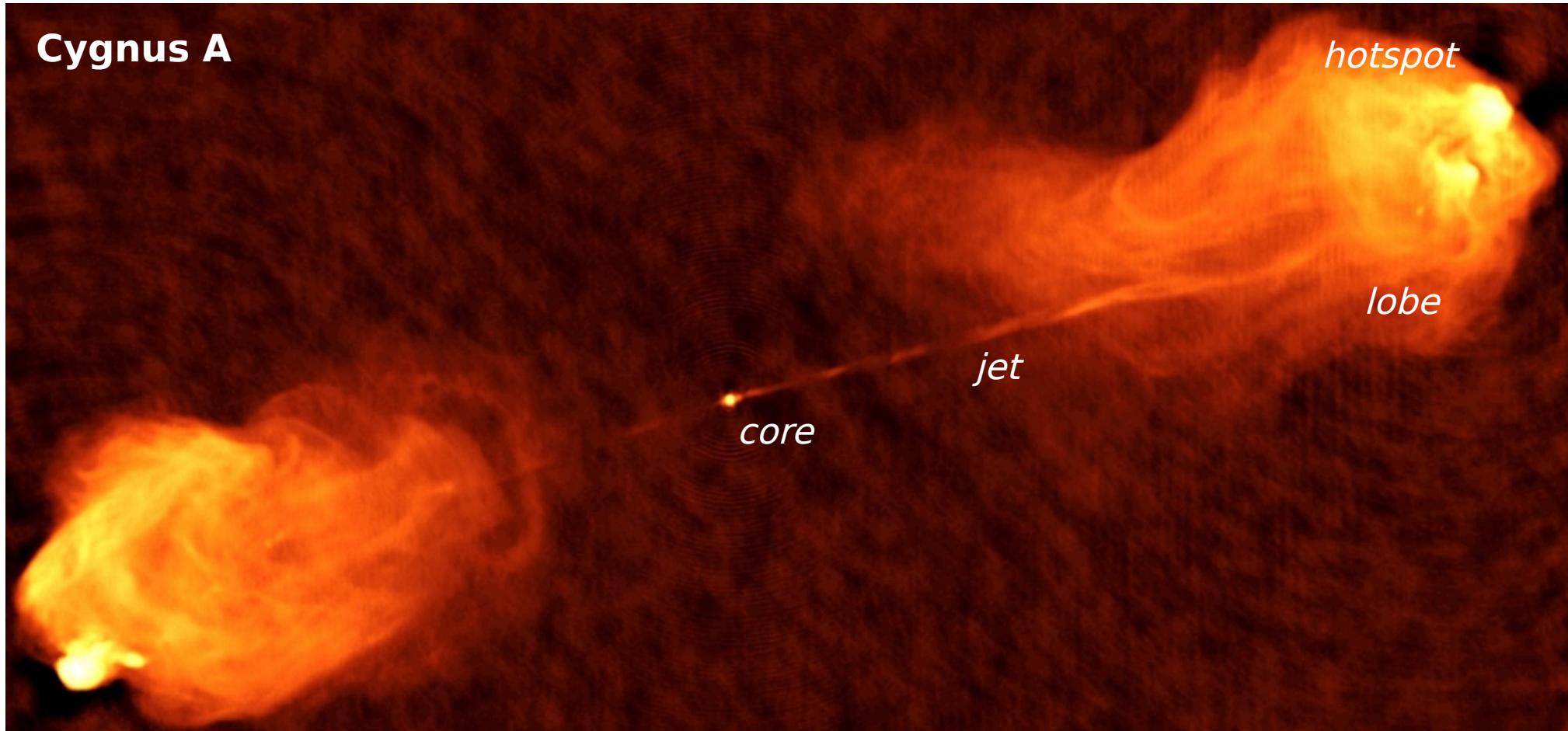
→ AGN jet hotspots/  
lobes could confine  
UHECRs.

Hillas criterion  
(Hillas 1984):  
 $R_g < \sim \text{size source}$

$$E \leq 10^{20} Z \left( \frac{B}{10 \mu \text{G}} \right) \left( \frac{L}{10 \text{ kpc}} \right) \text{ eV}$$

→ ... But can they be accelerated there?

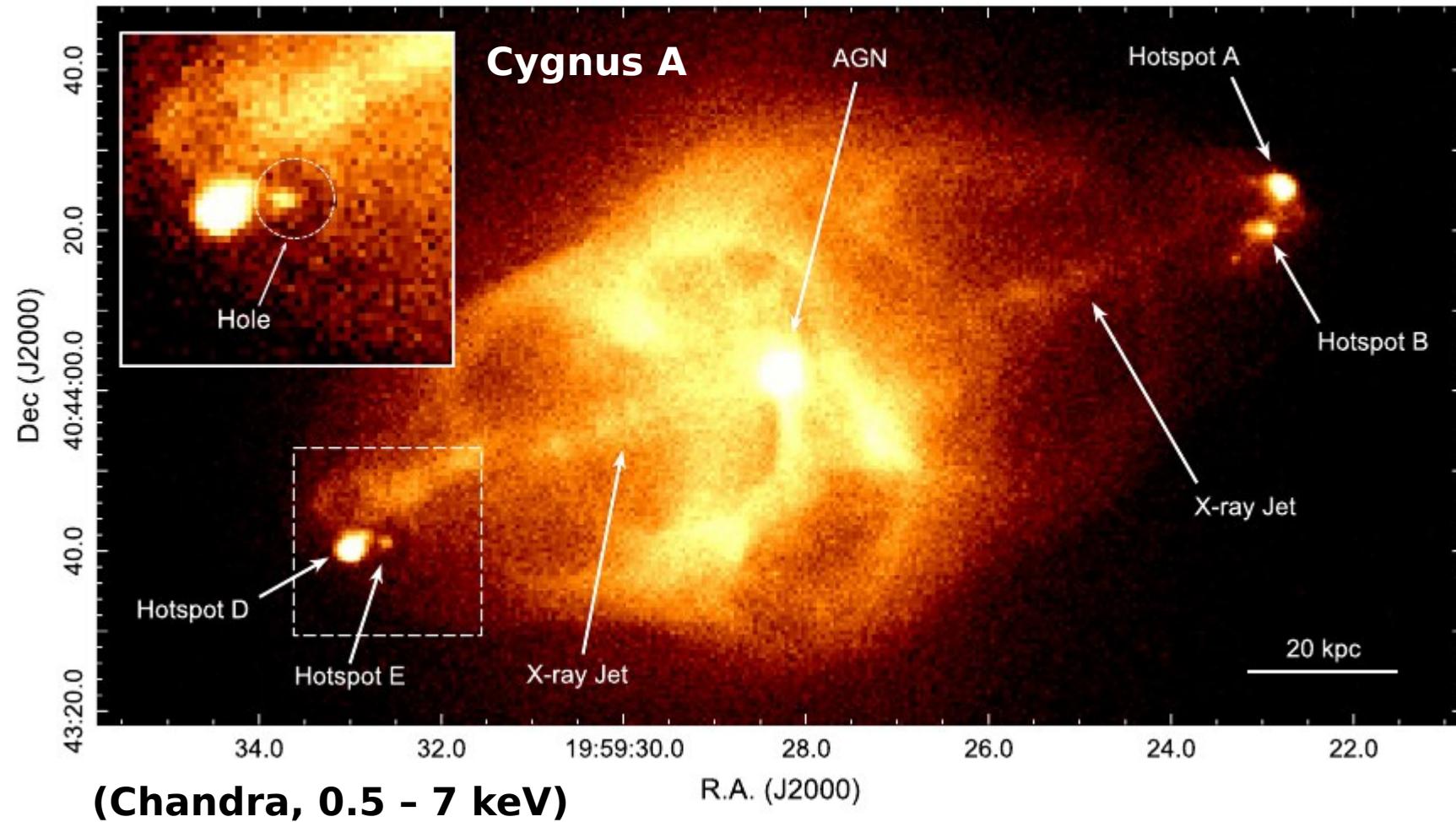
# FR II jets



# In-situ part. acceleration: Cygnus A hotspots

THE ASTROPHYSICAL JOURNAL, 891:173 (10pp), 2020 March 10

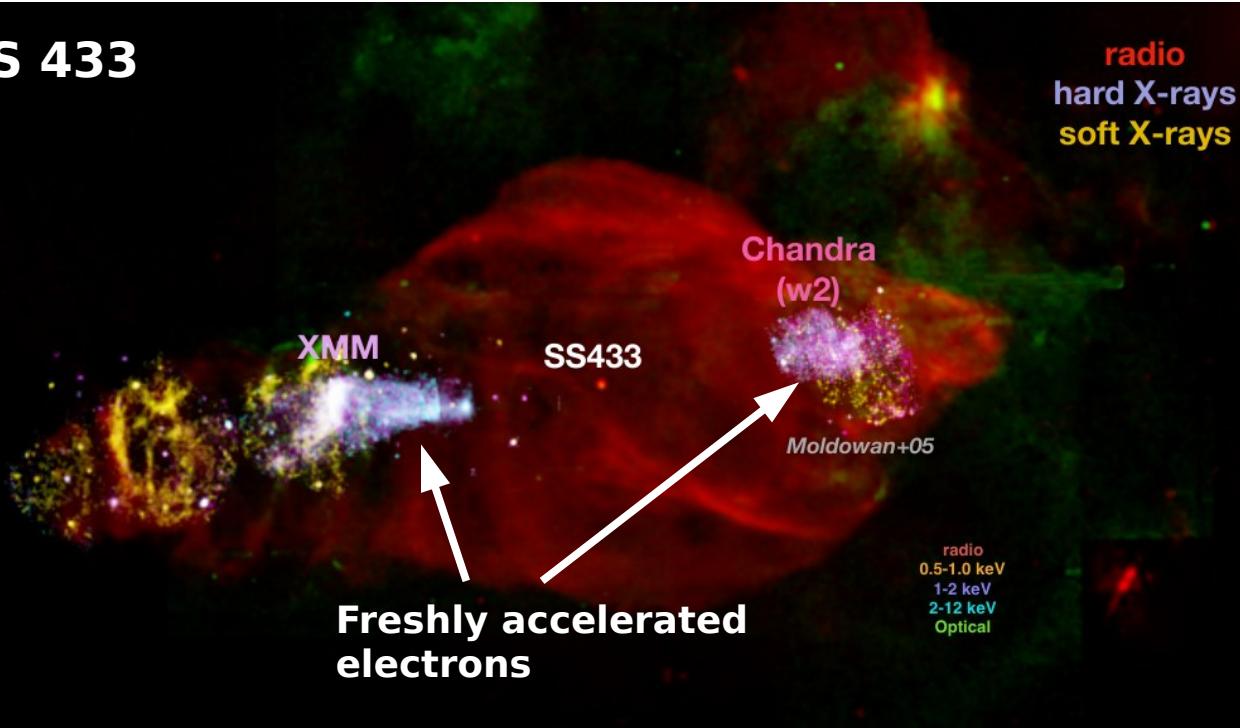
Snios et al.



# Same mechanism (for e<sup>-</sup>) in microquasars?

Safi-Harb et al. (2022)

SS 433



Detected by H.E.S.S.  
and HAWC.

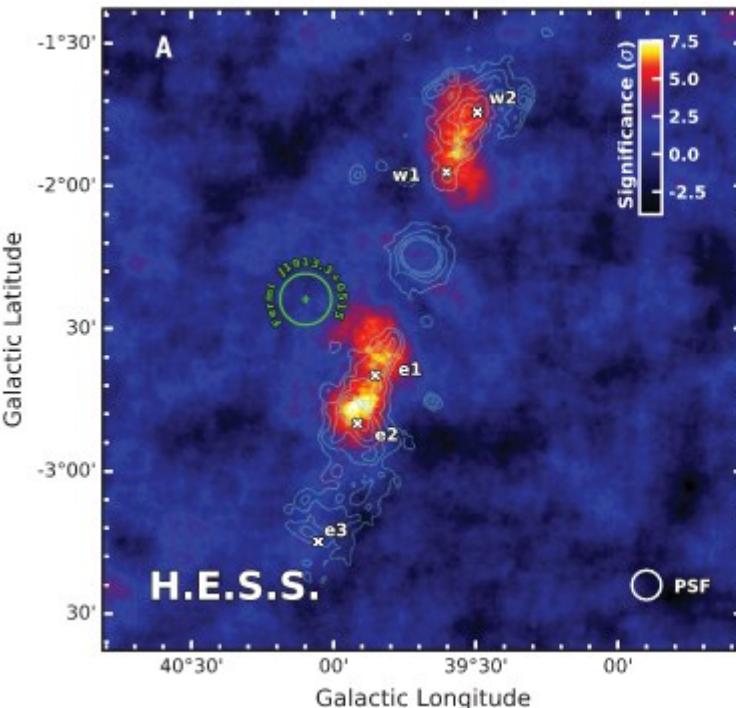
In-situ particle  
acceleration,  
to > 100 TeV.

# Same mechanism (for e<sup>-</sup>) in microquasars?

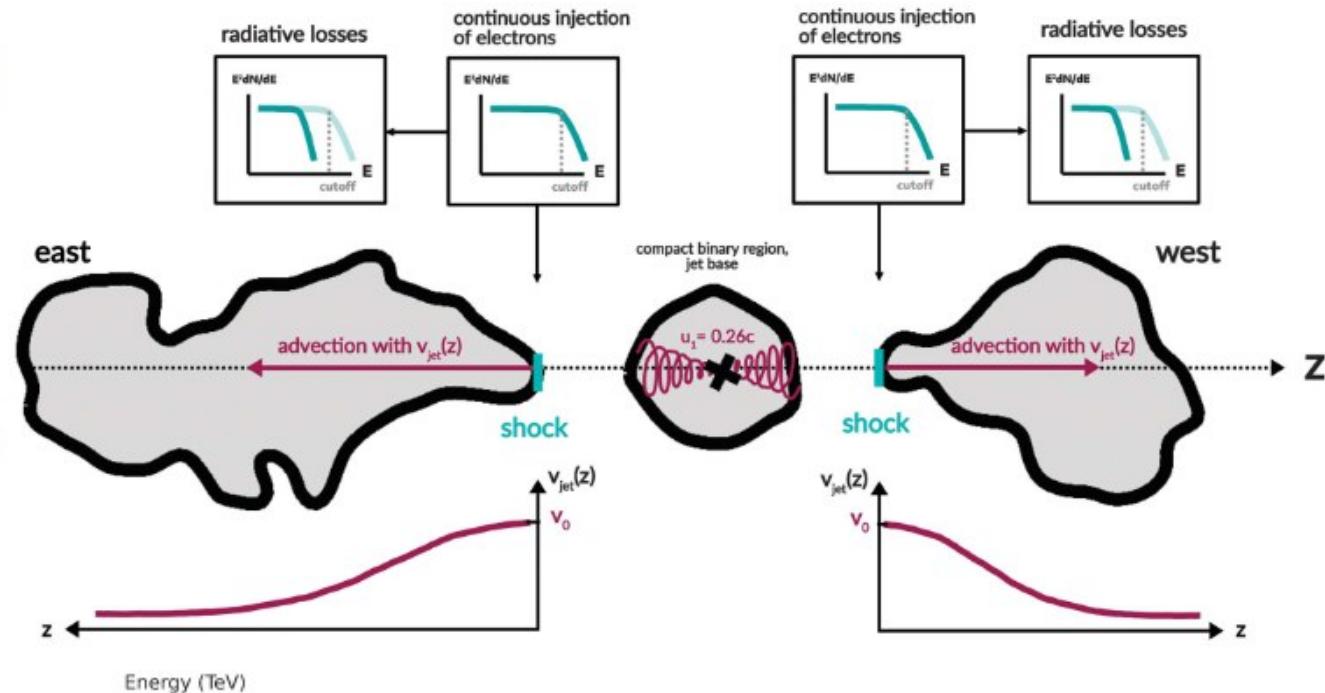
RESEARCH

GAMMA-RAY ASTRONOMY

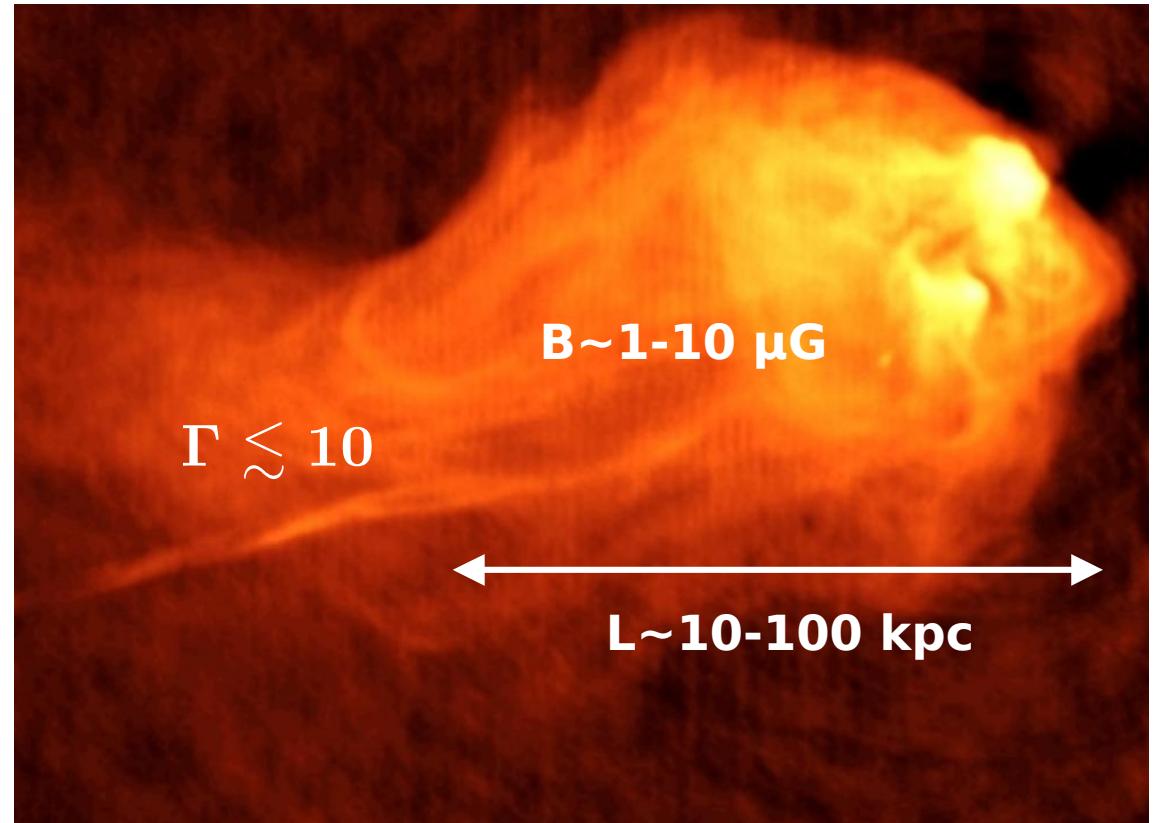
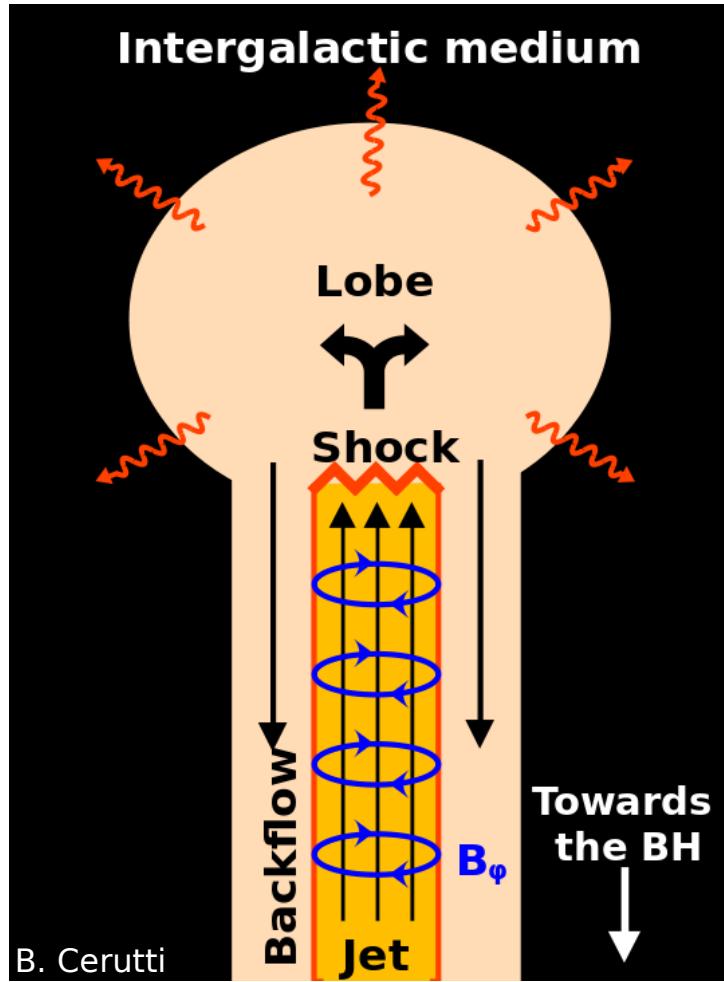
Acceleration and transport of relativistic electrons in the jets of the microquasar SS 433



H.E.S.S. Collaboration,  
Science 383, 402–406 (2024)

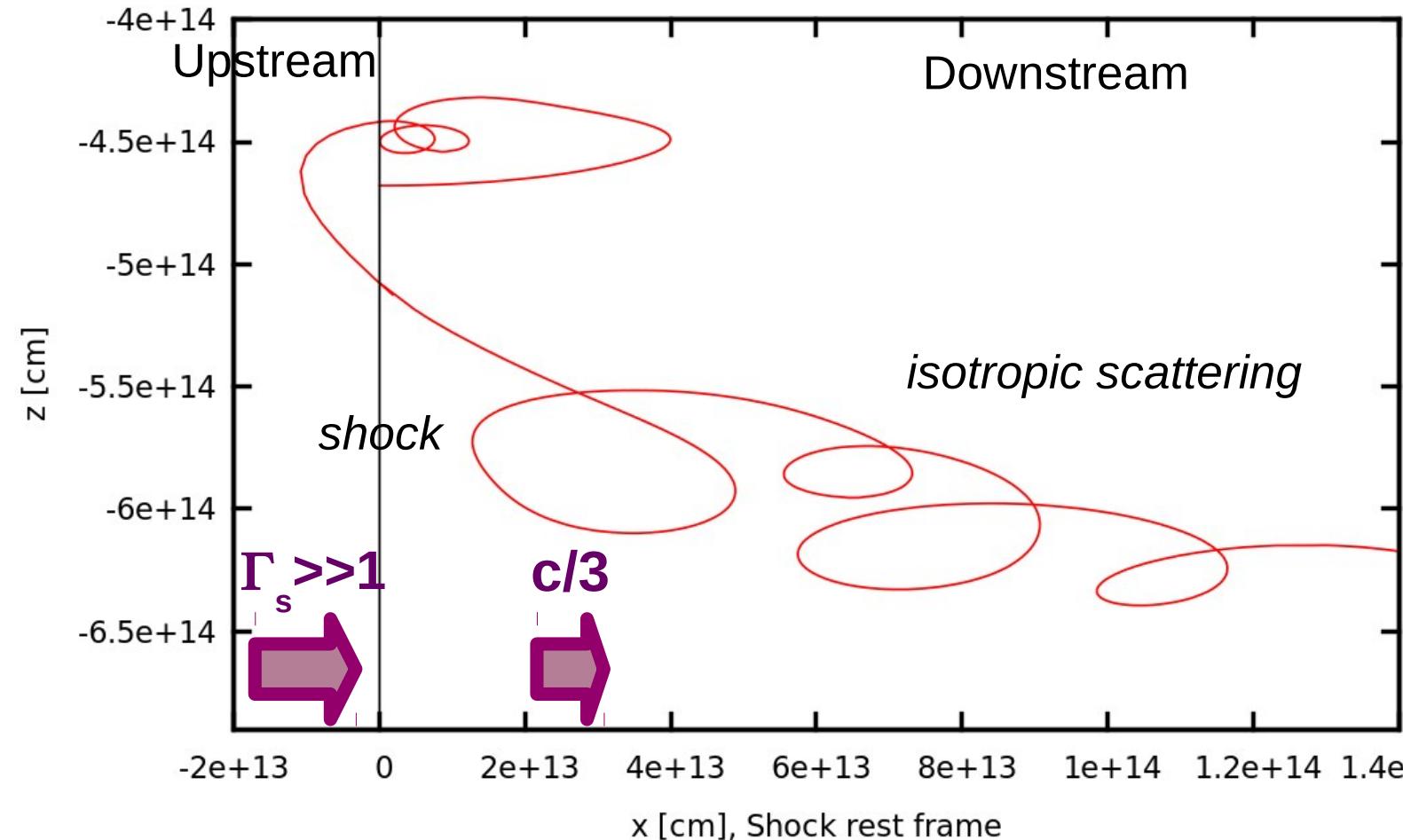


# Jet Termination Shock Region



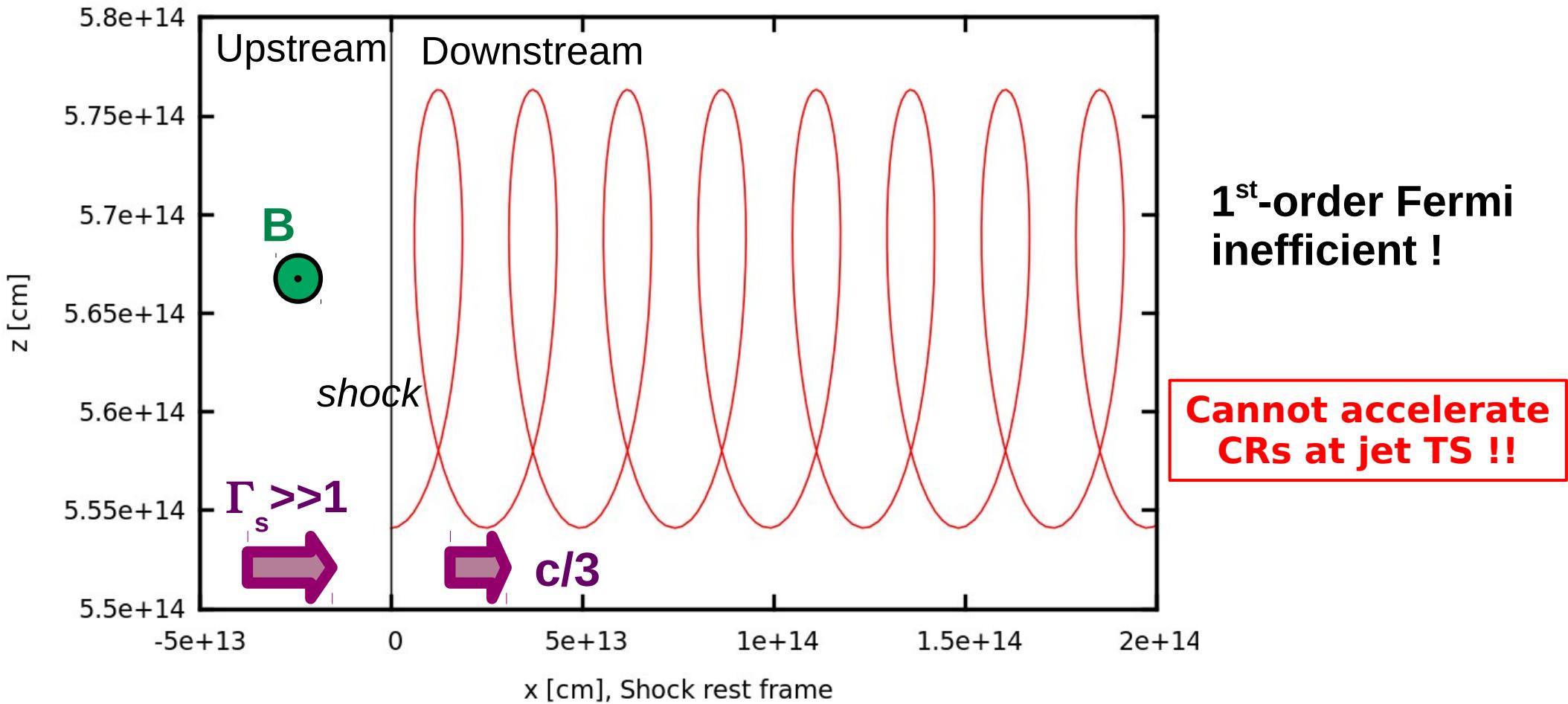
**Magnetization:**  $\sigma_0 = \frac{B_0^2}{8\pi\Gamma_0 n_0 m_p c^2} \sim 0.01 - 1.$

# Particle acceleration - relativistic shocks



$E^{-2.2}$  spectrum at ultra-relativistic shocks.

# At relativistic perpendicular shocks...



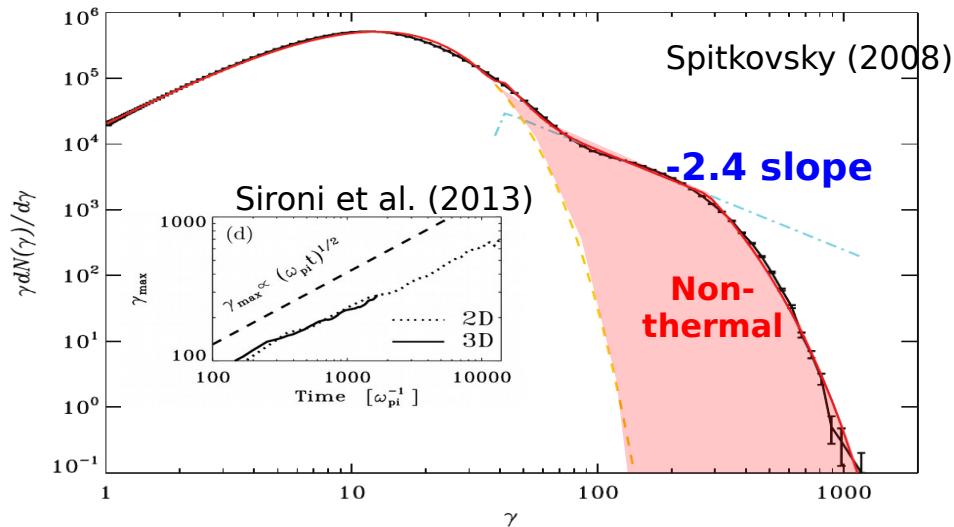
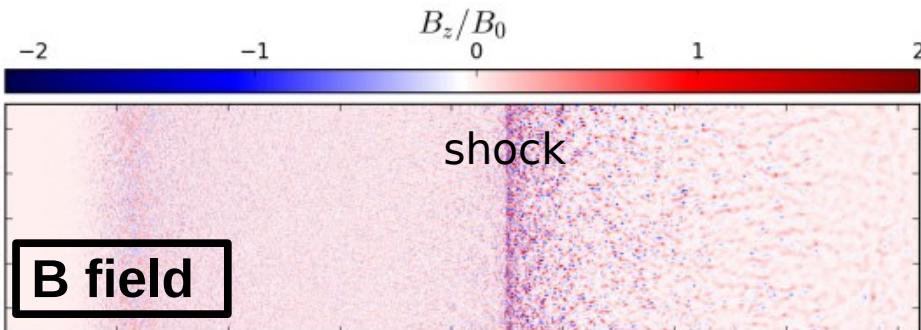
# Particle-In-Cell (PIC) simulations

→ Unmagnetized case ( $\sigma=0$ ):

*Spitkovsky (2008), Sironi + (2013),  
Plotnikov+ (2018), Lemoine+ (2019)*

**Good but slow** accelerators.

**Maximum energy grows as  $t^{1/2}$**   
*(Reville & Kirk 2010, Plotnikov et al. 2013)*



Weibel-dominated shock: Fermi-acceleration  
on small-scale plasma turbulence

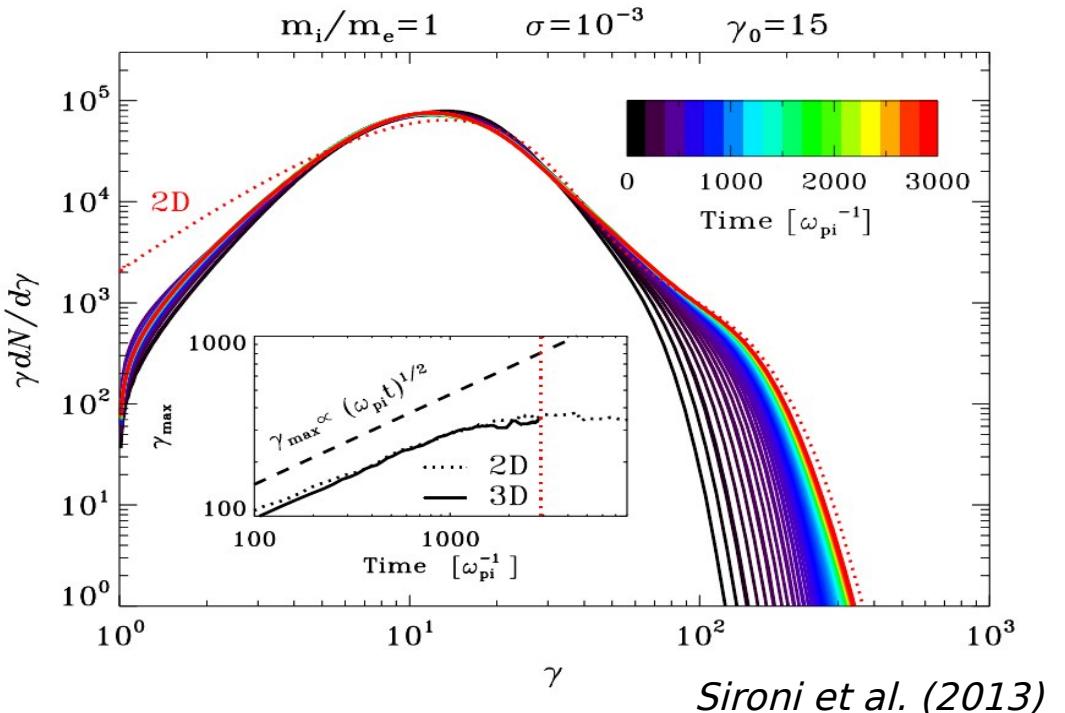
# Particle-In-Cell (PIC) simulations

→ Magnetized case ( $\sigma > 10^{-3}$ ):

Even weak magnetization levels stop particle acceleration.

$E_{\max}$  quickly saturates.

**Cannot accelerate CRs to UHE at jet TS !!**

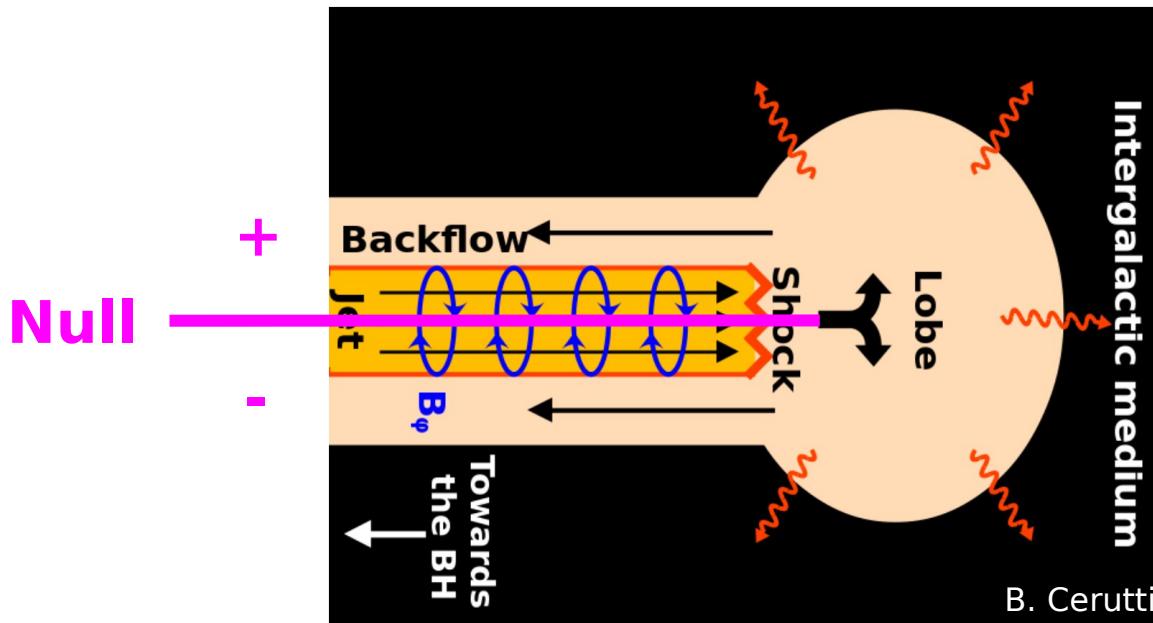


Sironi et al. (2013)

# Solution: Global B field geometry

This was for **plane-parallel, homogeneous** shocks...

**LARGE-SCALE GEOMETRY OF THE MAGNETIC FIELD MAY SOLVE THE PROBLEM!**

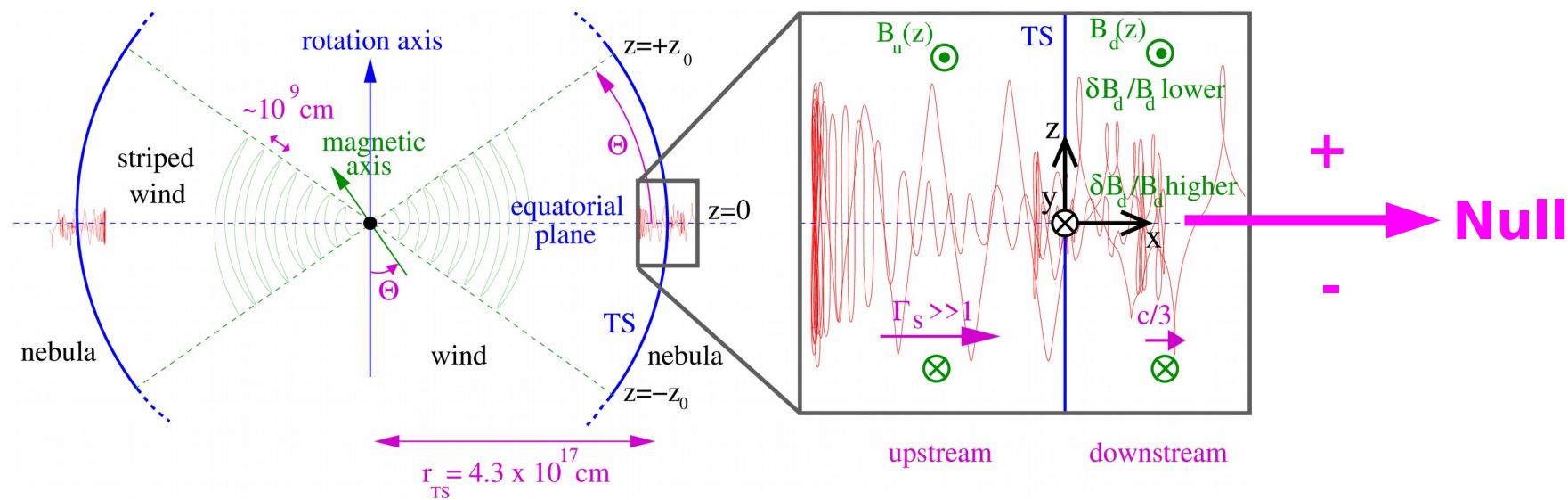


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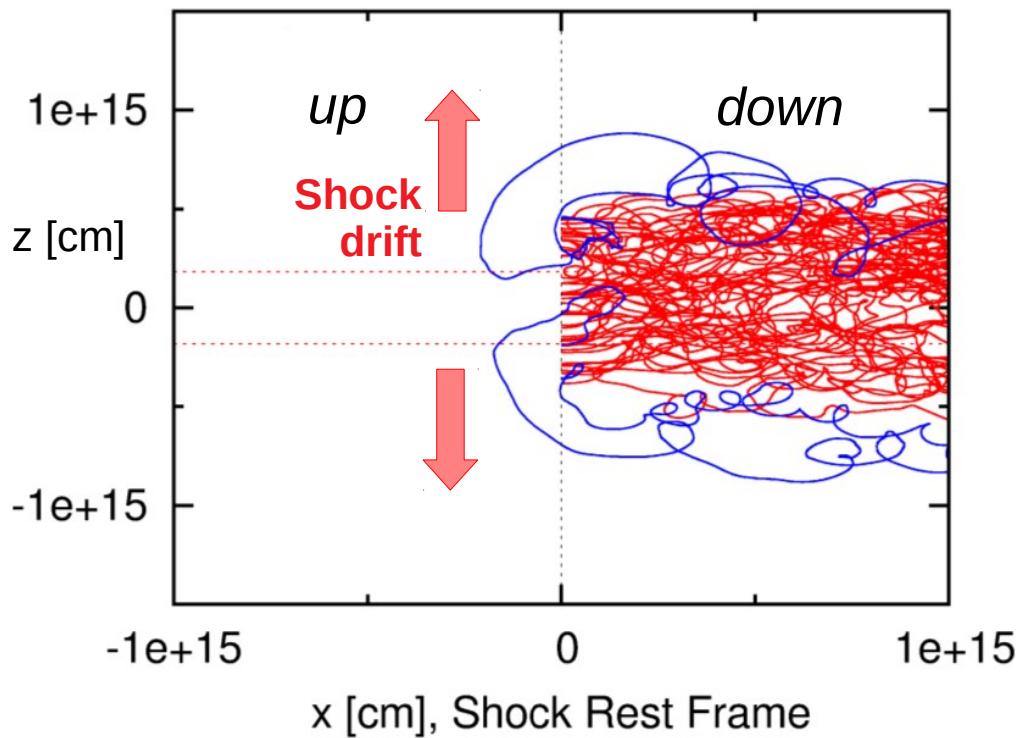
**LARGE-SCALE GEOMETRY OF THE MAGNETIC FIELD MAY SOLVE THE PROBLEM!**

See **Giacinti & Kirk (2018)** for Pulsar Wind Nebulae :



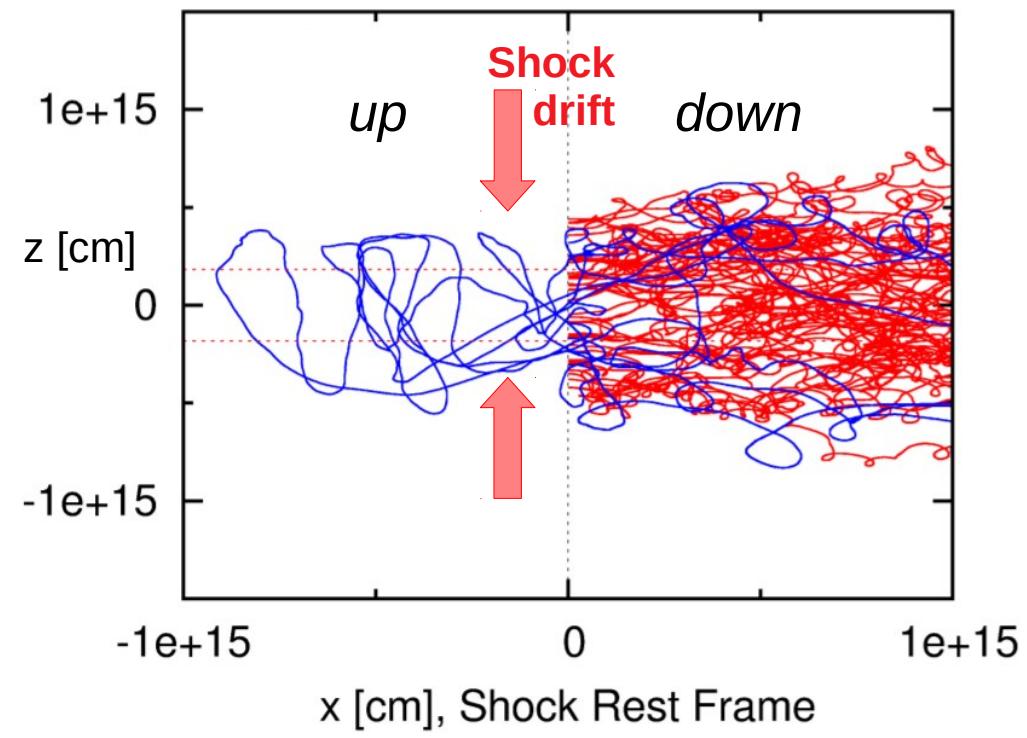
See **Giacinti & Kirk (2018)** for Pulsar Wind Nebulae :

**Positrons :**



No/little acceleration

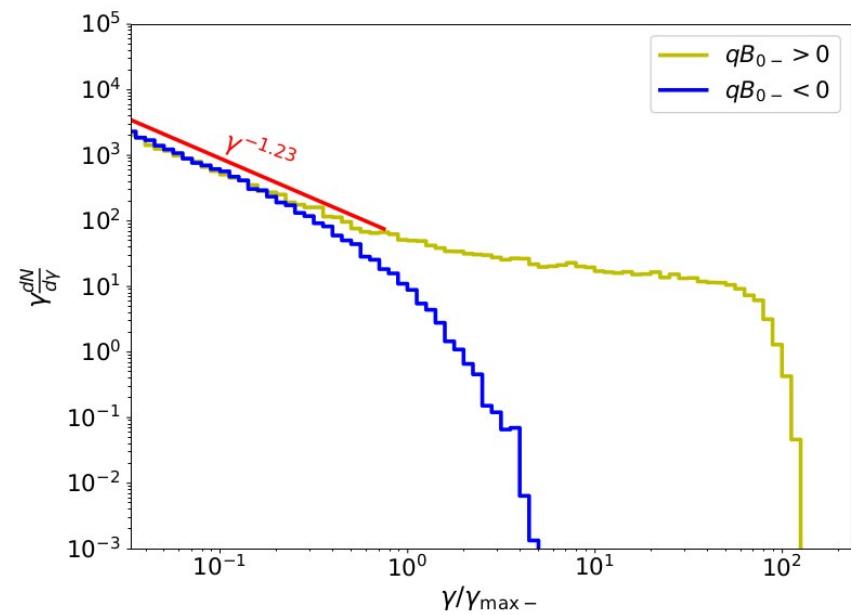
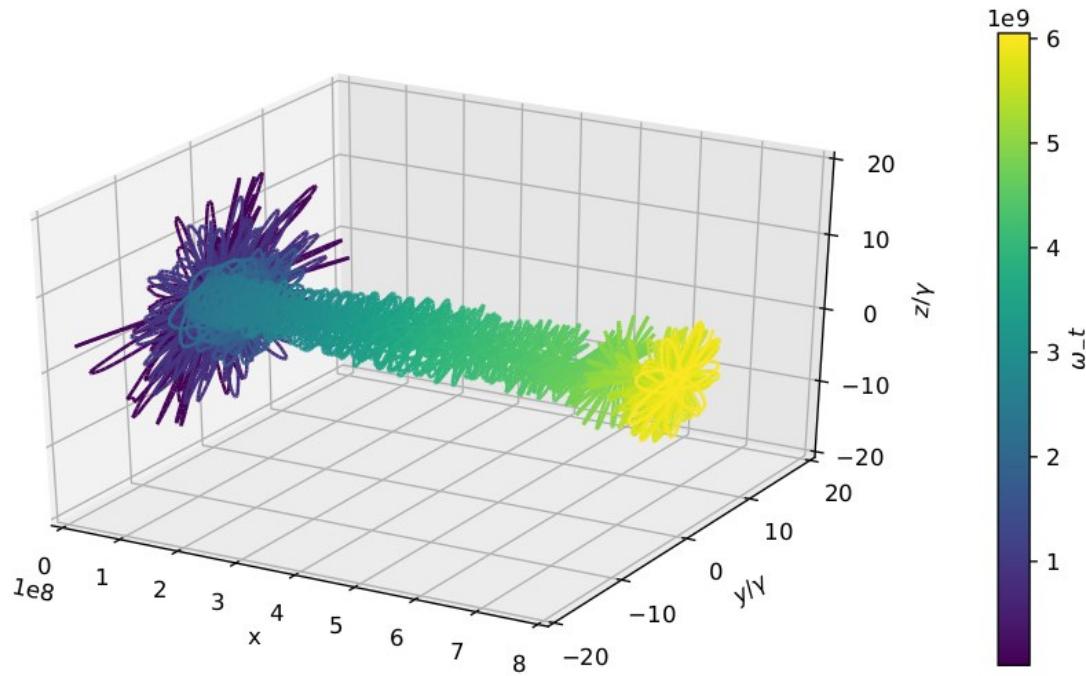
**Electrons :**



Acceleration

# Particle acceleration mechanism

→ Though shock acceleration if CR pressure is not too large (i.e. in test-particle limit):  
**Huang, Reville, Kirk, GG, MNRAS 522, 4955 (2023)**



**Key point: Particles (w/ correct sign of charge) remain around the null point**

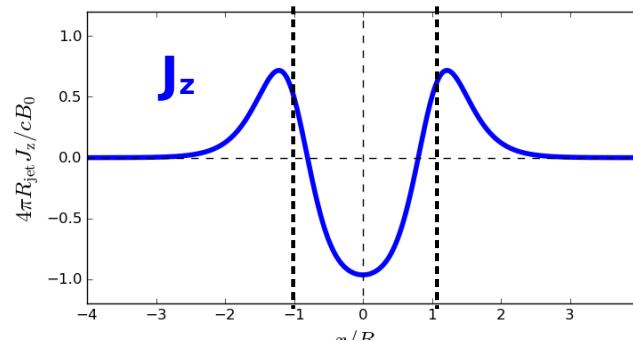
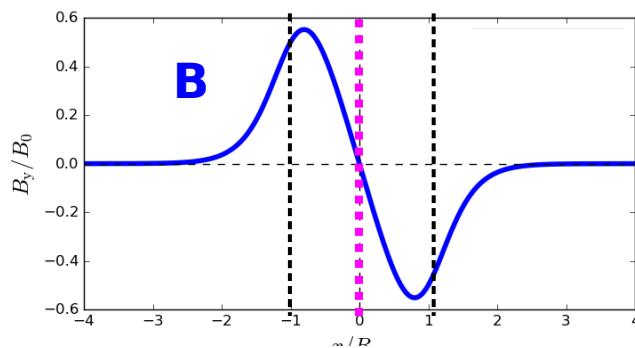
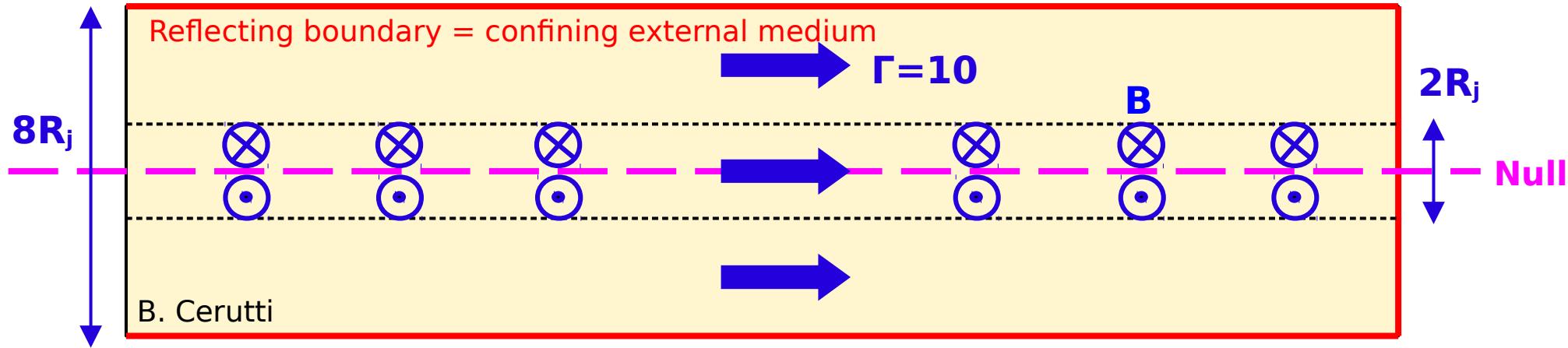
# Particle-In-Cell (PIC) setup

2D Cartesian box (xz-plane), **262,144×16,384 cells, or 6554x410 d<sub>i</sub> (ion skin depth)**

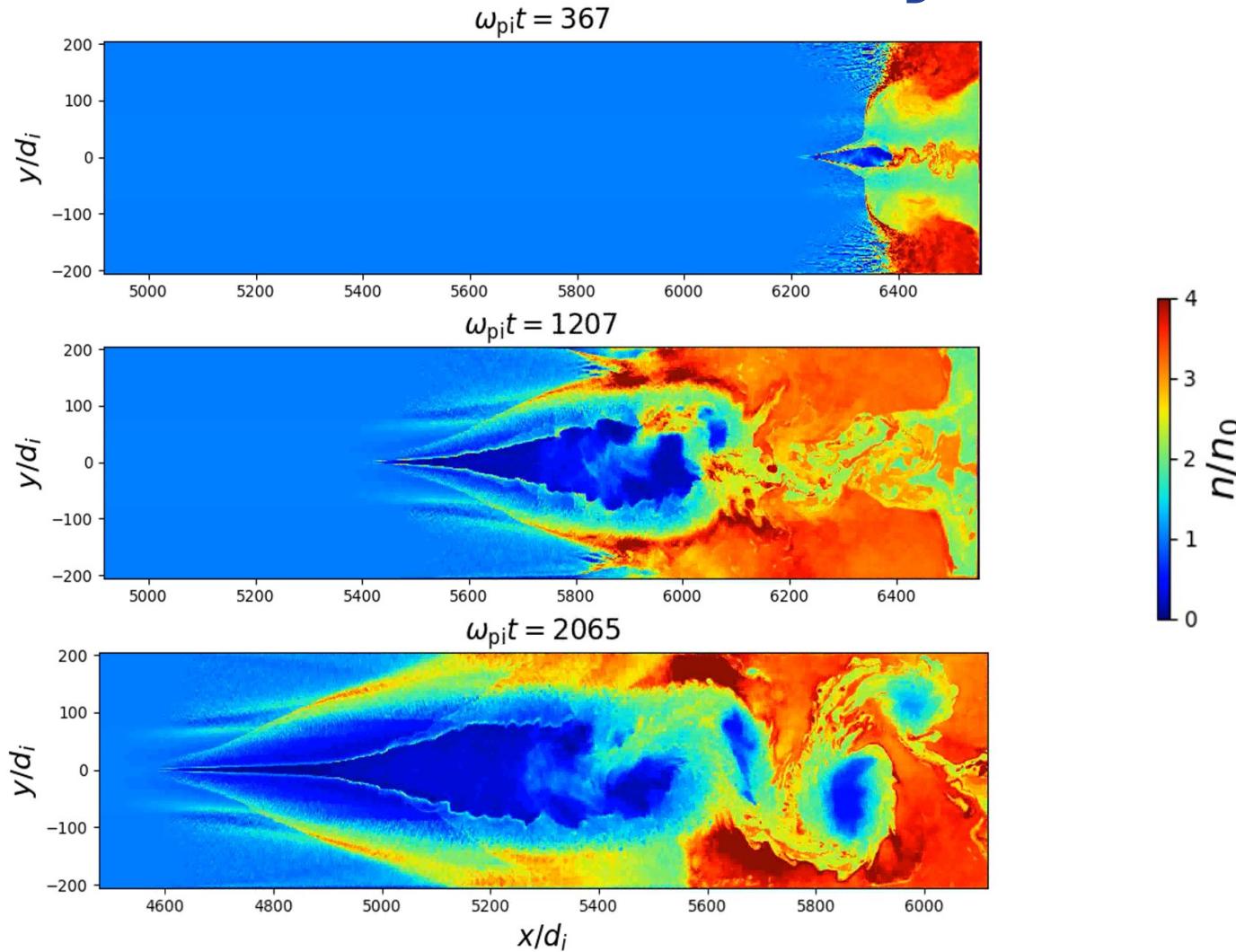
Electron-ion plasma with  $m_i/m_e=25$

Magnetization :  $\sigma=0.1, 1$

Reflecting boundary  
= contact discontinuity

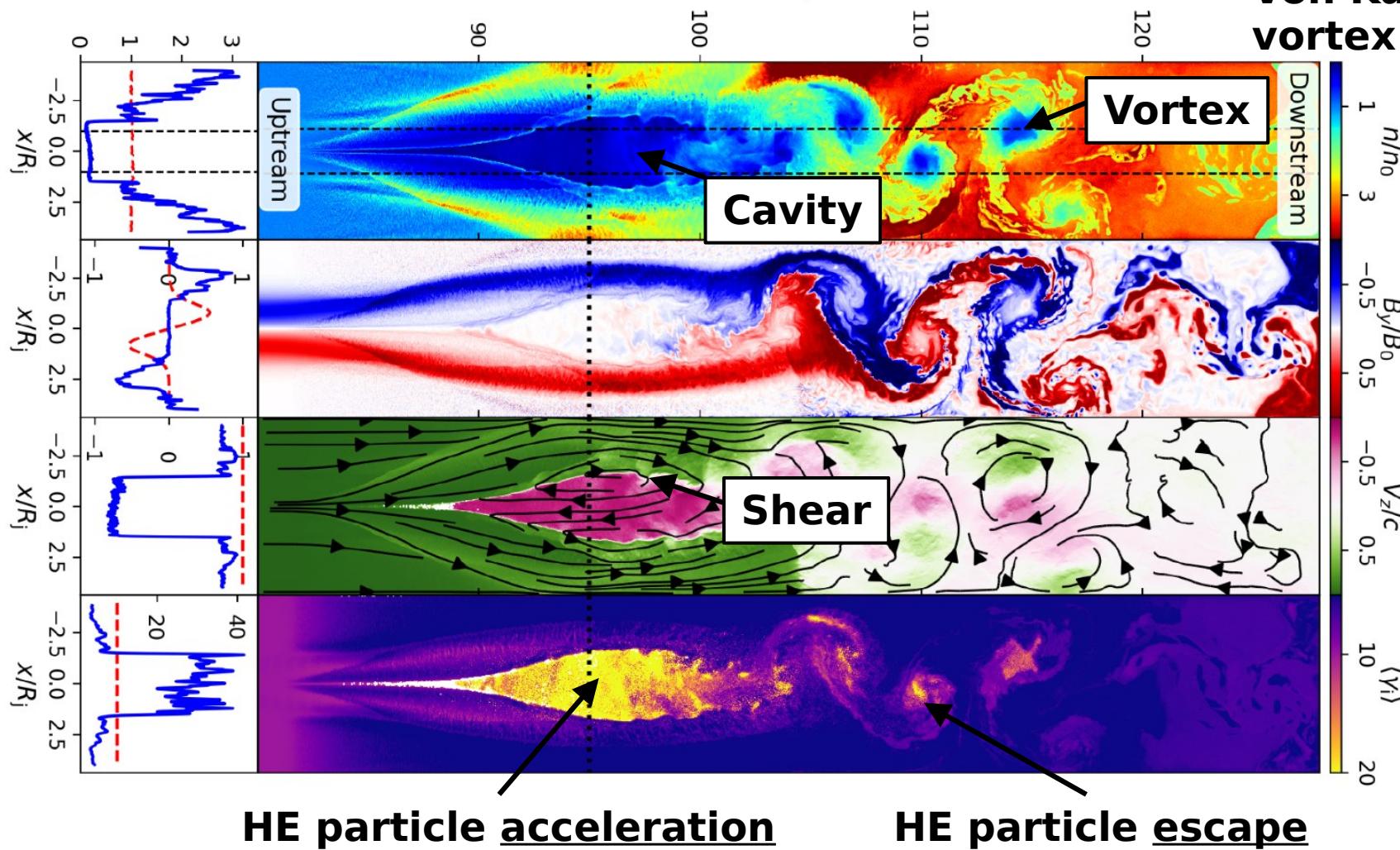


# Results PIC Sim.: Density evolution

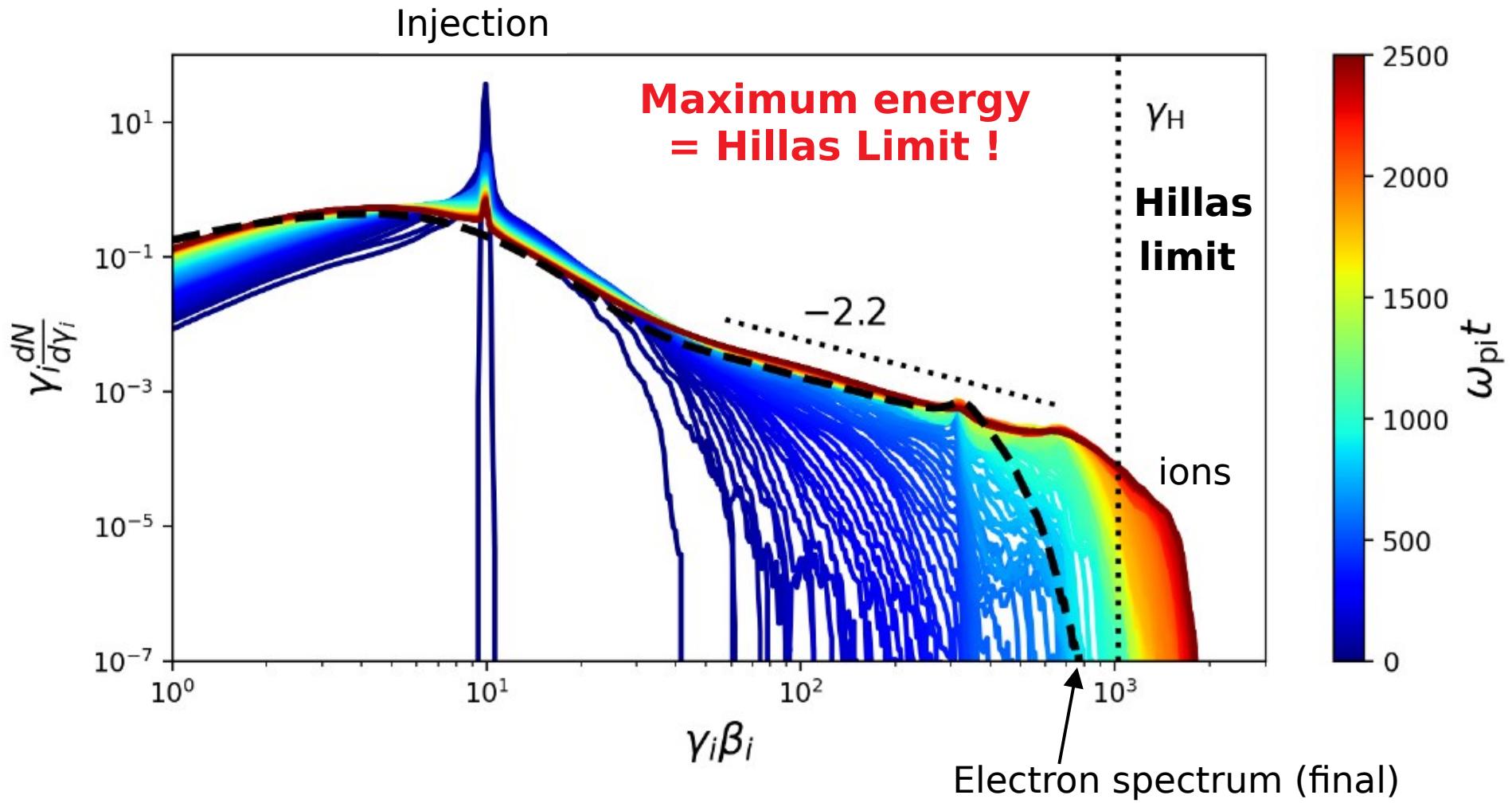


# Results - PIC Simulations

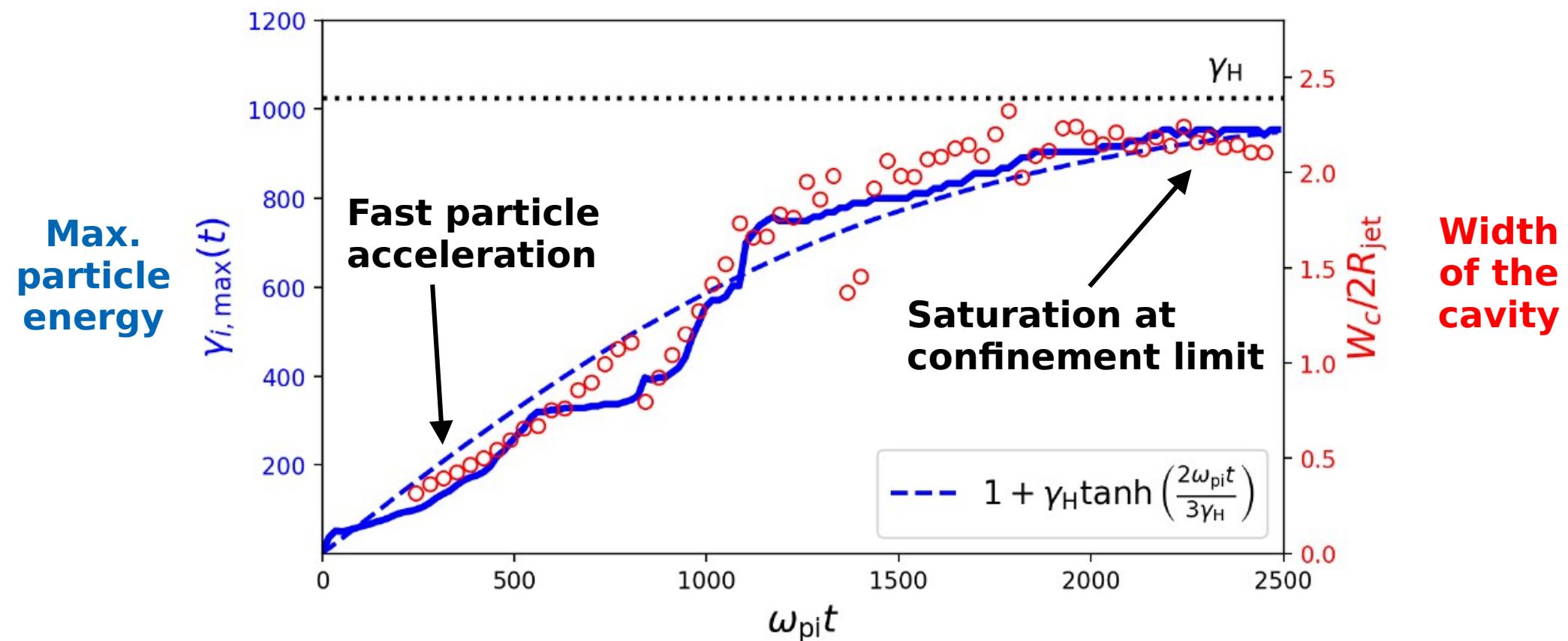
Density



# Ion spectrum: Time Evolution & $E_{\max}$



# $E_{\max}$ ions & Cavity size: Time Evolution



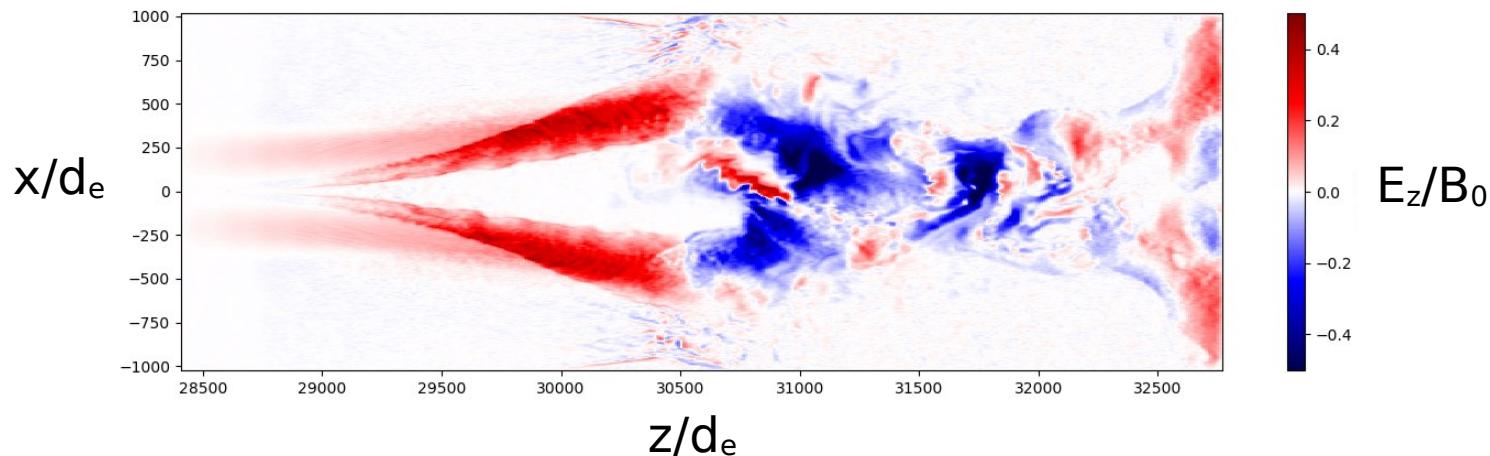
Maximum particle energy grows as cavity width.  
Cavity stops growing at  $\sim$  width jet  $\Rightarrow$  **Hillas criterion**.

# Particle acceleration mechanism

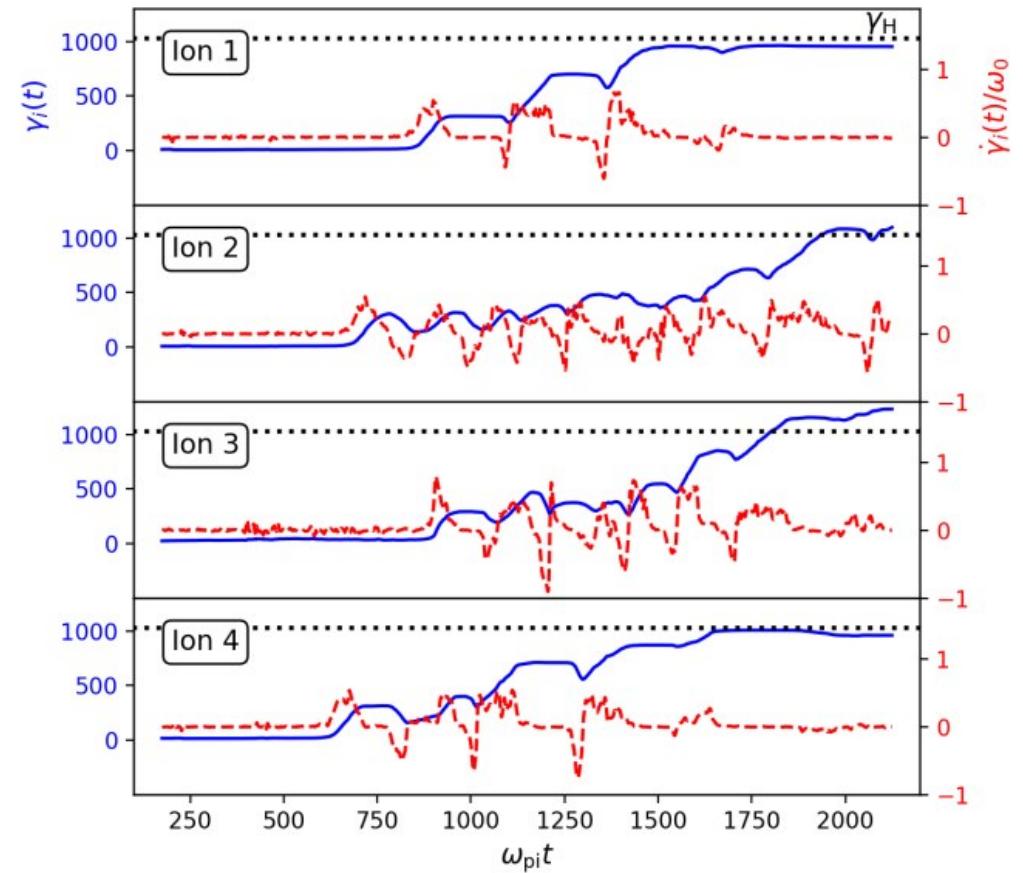
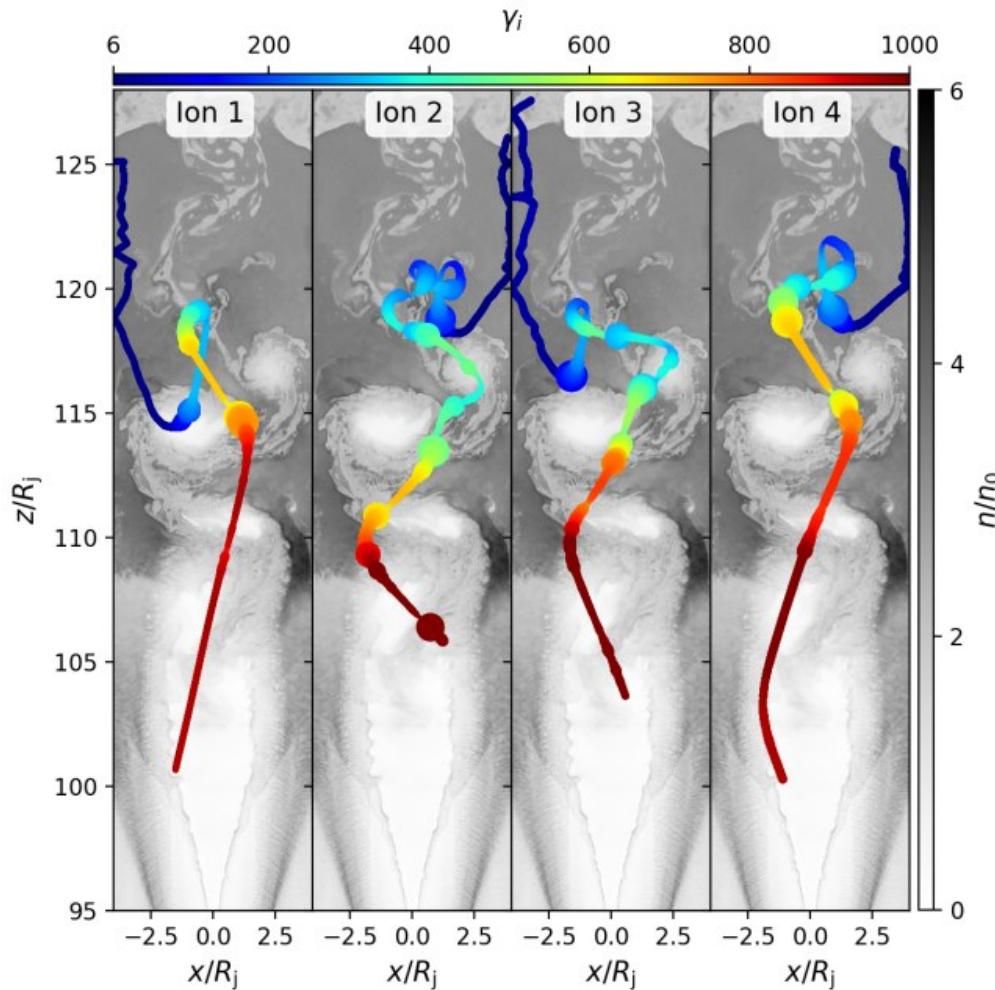
- Not standard shock acceleration mechanism here...
- **Shear-flow acceleration at the edges of the cavity instead**



Ideal E field in the lab frame:  $E = -\frac{\mathbf{V} \times \mathbf{B}}{c}$

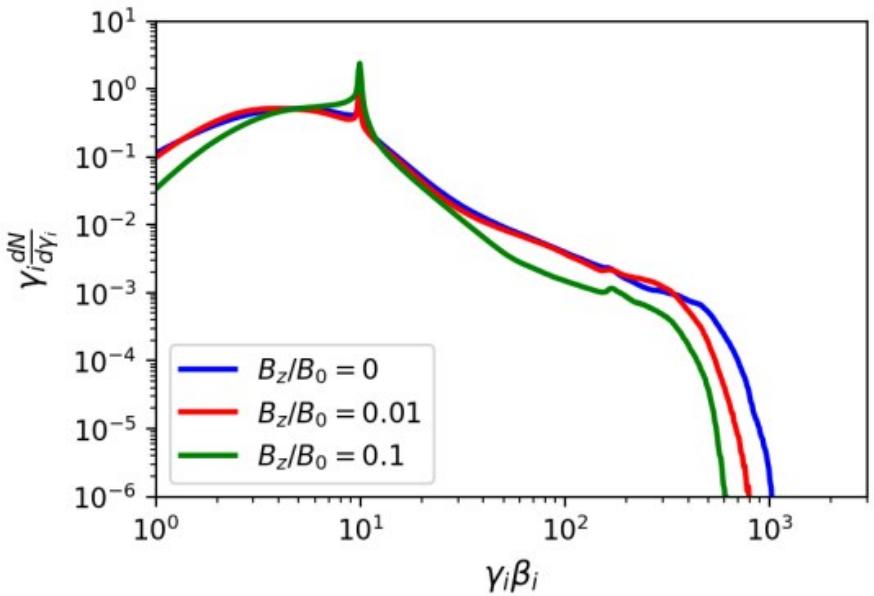
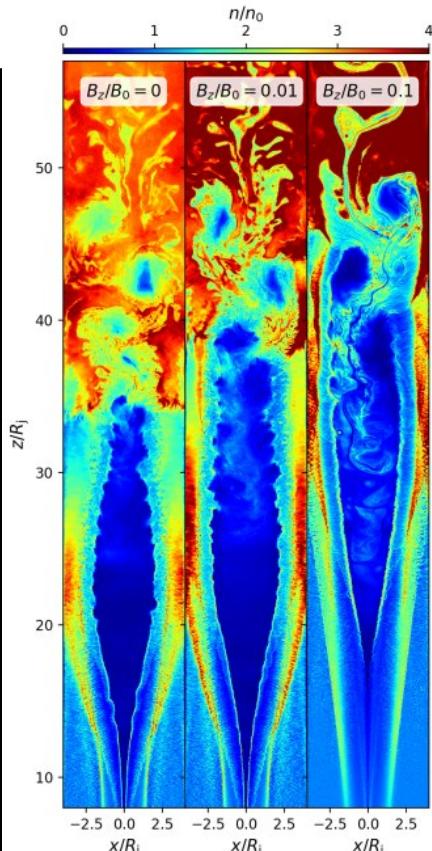
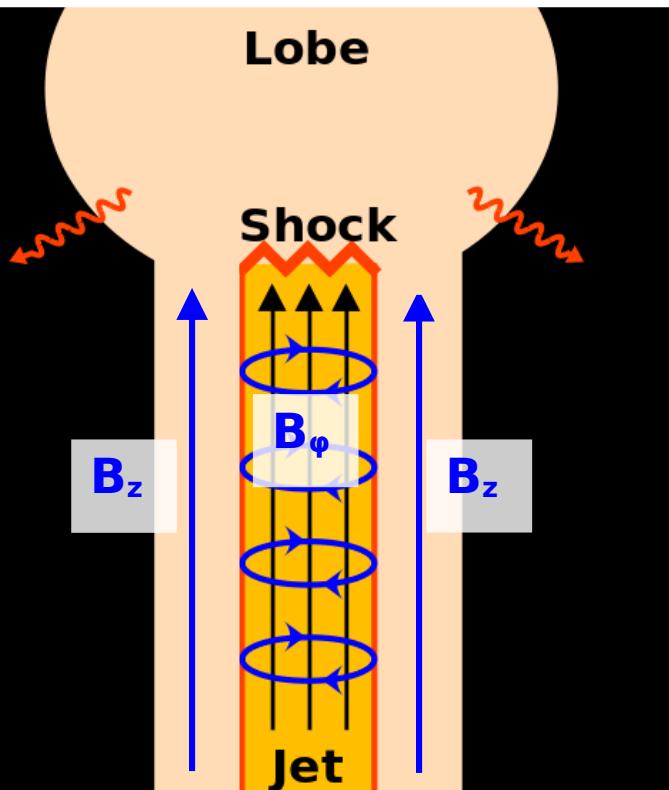


# Particle acceleration mechanism

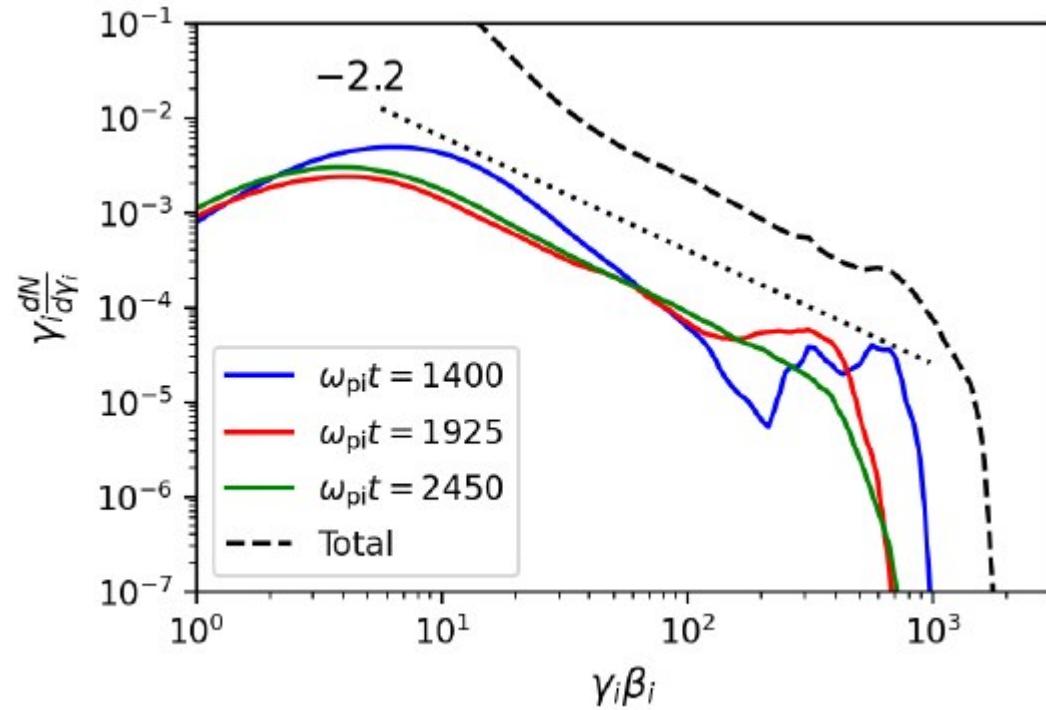
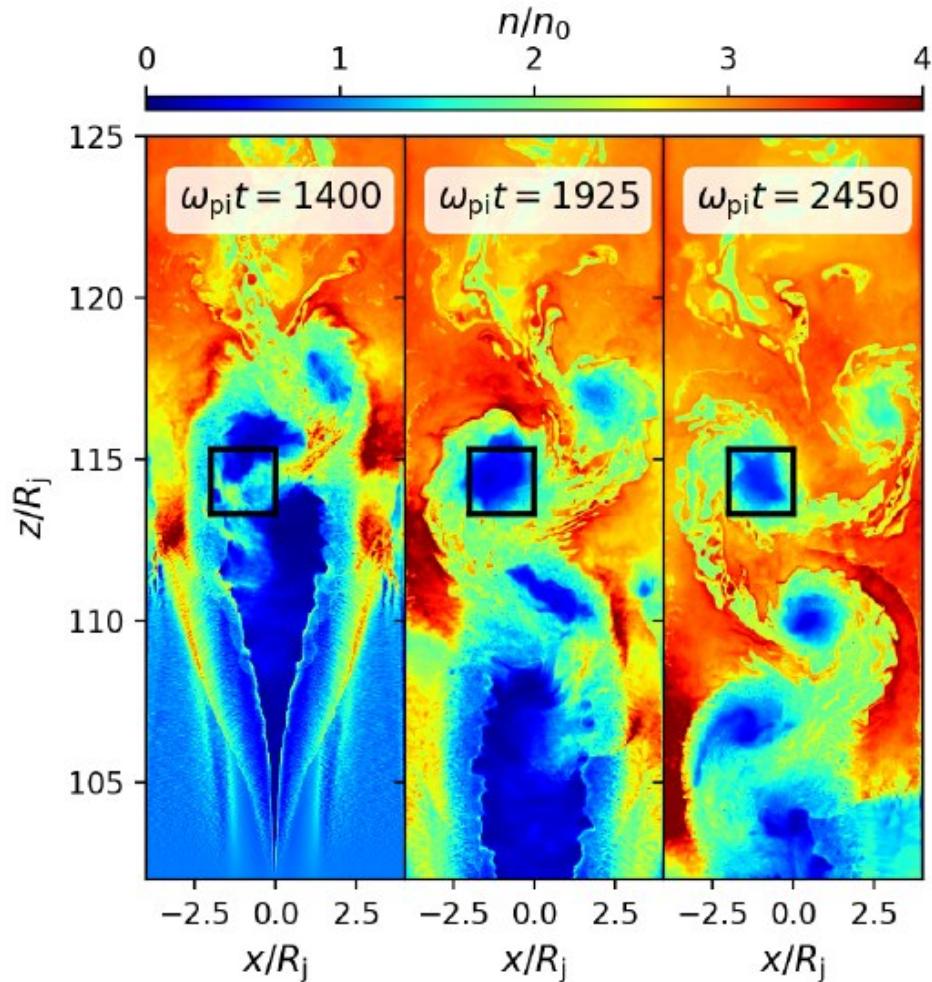


# Effect of a poloidal B field component

If  $B_z < B_\phi$  (expectation in jet TS region), particle acceleration remains efficient



# Mechanism for VHE particle escape



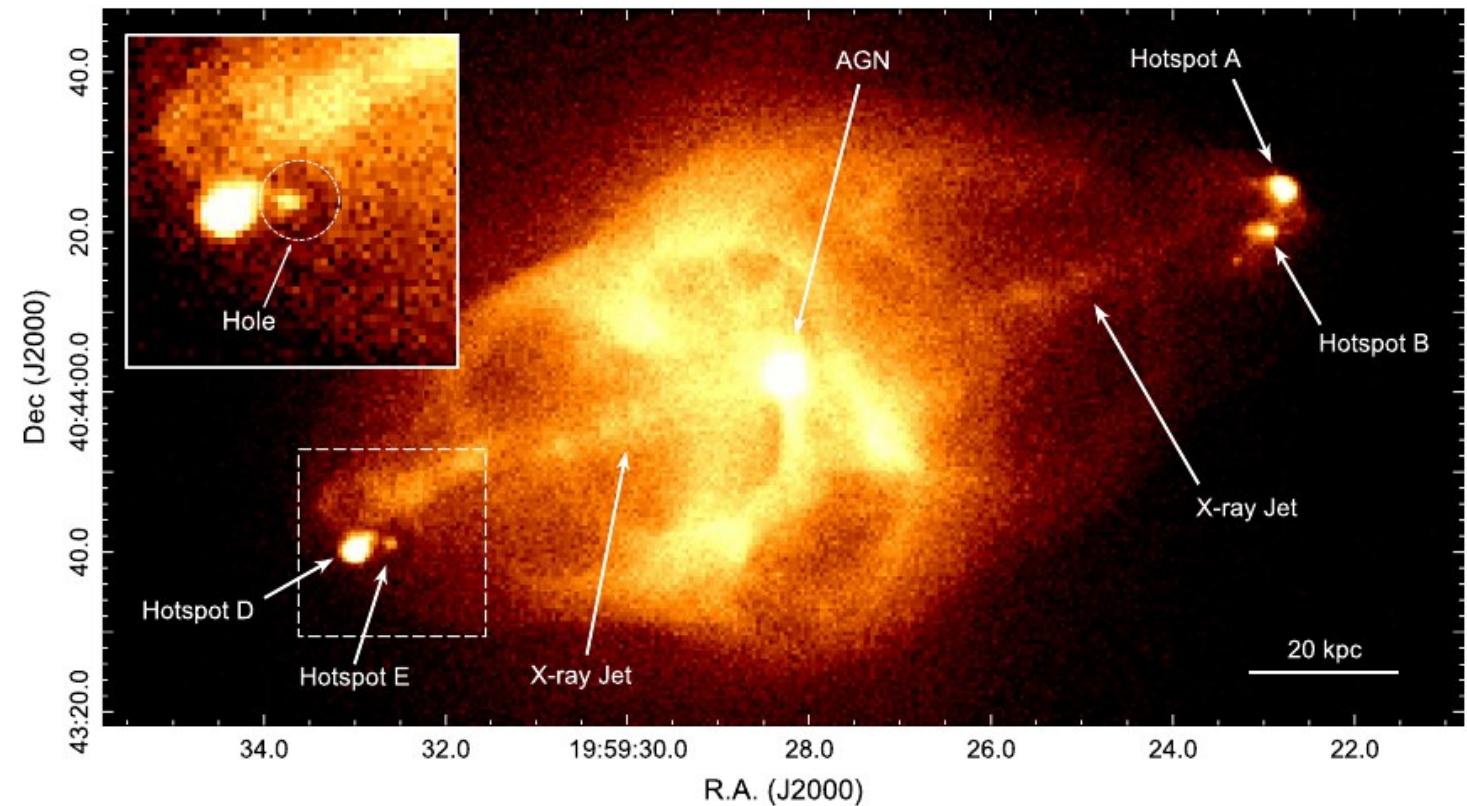
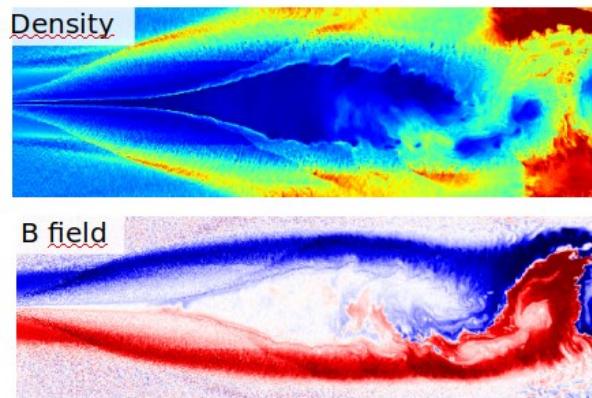
→ Particles escape in von Kármán vortices.  
→ **No E losses.**

# Observational test / evidence ?

THE ASTROPHYSICAL JOURNAL, 891:173 (10pp), 2020 March 10

Snios et al.

Cavity → under-luminous holes?



# Conclusions & Perspectives

- Large-scale structure of the B fields key for having particle acceleration,  
→ Very generic: May apply for GRBs, microquasars, PWNe, ...
- CR-dominated cavity at the shock front around the B field null point,  
→ Search for cavities.
- Particles accelerated at the shear flows around the cavity,
- Particles accelerated to the Hillas Limit (at least in the simulations)  
This mechanism could accelerate hadrons to UHEs at AGN jet TSs.  
... and to PeV in stellar-mass BH jets, e.g. in SS433,
- CRs escape in the downstream inside von Kármán vortices.
- Next step: Study this problem with PIC-MHD simulations.