

(High-mass) Gamma-ray binaries hosting a pulsar: Prospects at the highest energies

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The Extreme Non-Thermal Universe: CDY Initiative

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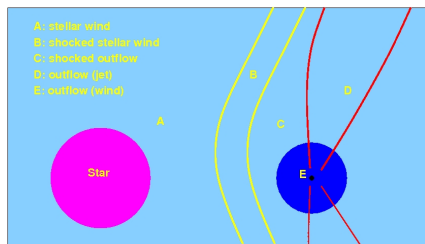
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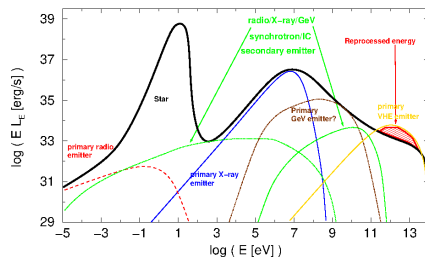
High-mass gamma-ray binaries (with CO)

- **The phenomenological term *Gamma-ray binary usually means star+CO and SED dominance of gamma rays (without the star).***
- High-mass gamma-ray binaries (HMGB) are among the most powerful galactic sources:
 $L \sim 10^{36-37}$ (MeV), 10^{34-37} (GeV)
and 10^{32-35} erg s⁻¹ (TeV).
- The great majority of the known HMGB are VHE emitters.

(Some reviews: Mirabel 2006; B-R & Khangulyan 2009; Dubus 2015; Paredes & Bordas 2019; Chernyakova & Malyshev 2020...)



Main elements of a HMGB.

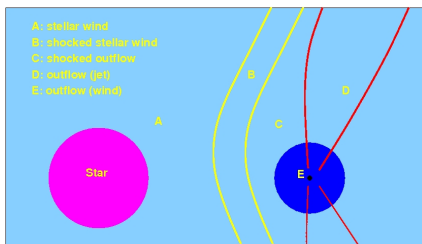


Typical SED of a HMGB.

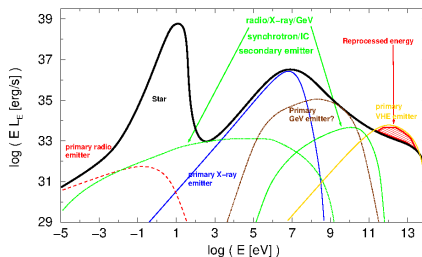
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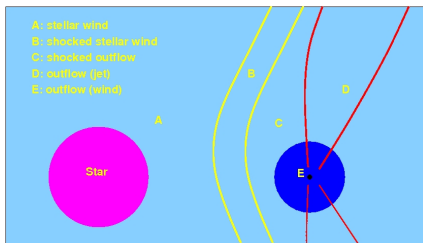


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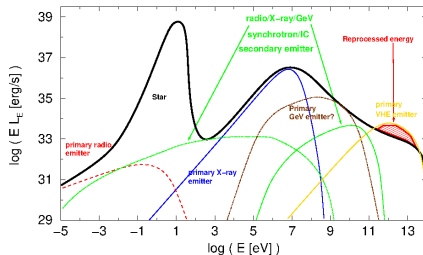
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Some history

- **In the 70s and 80s, a non-accreting pulsar plus a star were a common scenario to explain gamma-ray activity in binaries.**
- In the 90s, microquasars became also popular, often seen as scaled down versions of active galactic nuclei.
- In the 00s, pulsar and microquasar scenarios actively *competed* as the explanations behind the gamma-ray binary phenomenology.

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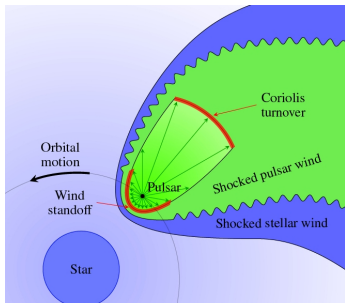
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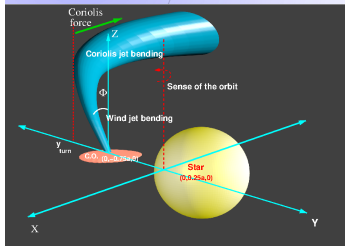
Non-accreting (pulsar) vs accreting (HMMQ) HMGB

- **A non-accreting HMGB consists of a young pulsar plus an OB star whose winds interact.**
- A HMMQ consists of a CO plus an OB star in which the wind is accreted and jets form, which interact with the wind.
- In both cases, outflows interacting along the orbit are complex and emit radio, X- and gamma rays, likely through synchrotron and IC, plus $\gamma\gamma$...

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High-mass star+young psr (Zabalza et al. 2013)

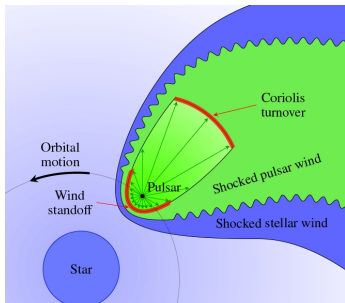


High-mass microquasar (Barkov & B-R 2021)

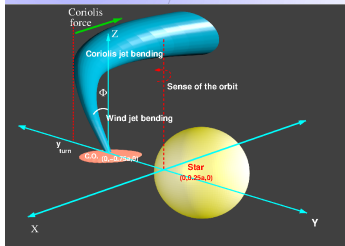
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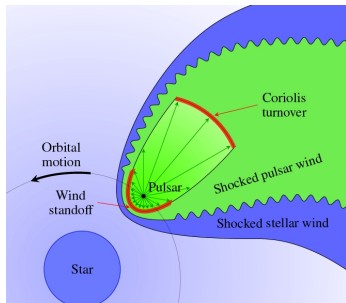


High-mass microquasar (Barkov & B-R 2021)

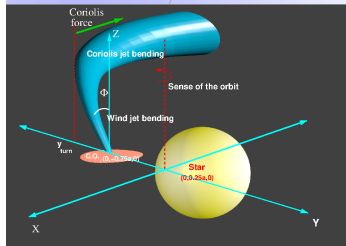
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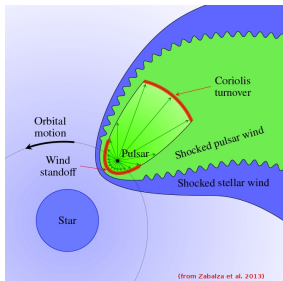


High-mass star+young psr (Zabalza et al. 2013)



High-mass microquasar (Barkov & B-R 2021)

Non-accreting HMGB



(Zabalza et al. 2013)

RHD simulations with PLUTO of 2-wind-orbit interactions (low e).

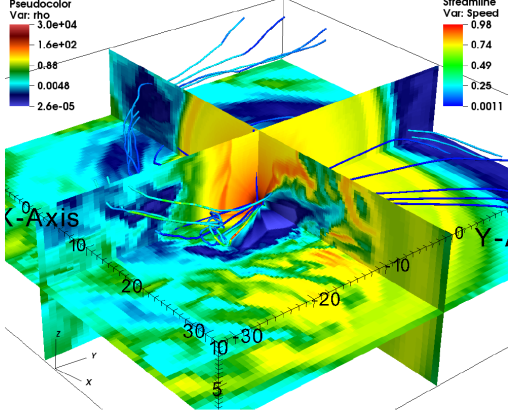
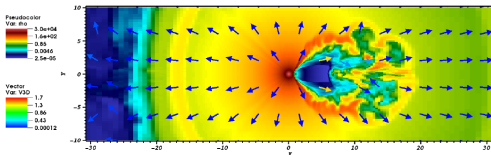
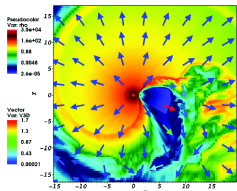


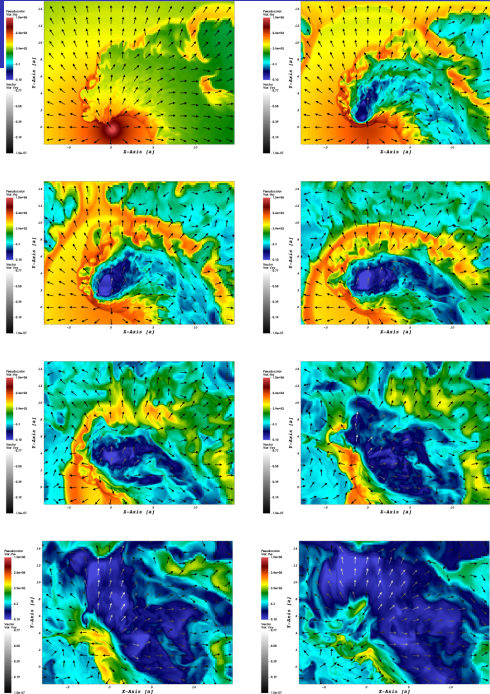
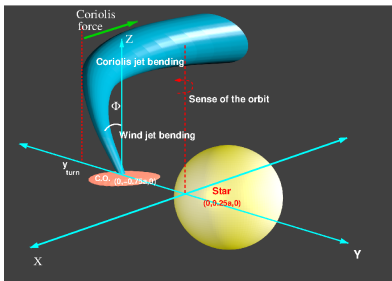
Fig. 2. Representation of the distribution of density in the XY -, XZ -, and YZ -planes for 3Df at $t = 3.9$ days (apastron). Streamlines are in 3D.



(LS 5039 at apastron; B-R, Barkov & Perucho 2015)

High-mass microquasar

RHD simulations with PLUTO of jet-wind-orbit interactions.



(HMMQ jet in a $e = 0$ orbit; Barkov & B-R 2021)

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- **Non-accreting pulsars are nowadays favored for those HMGB with unknown CO, (arguably) the main reason being the apparent lack of accretion features:**
 - LS 5039; LS I +61 303; HESS J0632+057; FGL J1018.6–5856; LMC P3; 4FGL J1405.1-6119; HESS J1832–093.
- The number of known accreting versus non-accreting sources are similar:
 - Cyg X-1, Cyg X-3, SS 433 versus PSR B1259-63, PSR J2032+4127.

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Pulsar or not (observation/phenomenology)

- **Accreting and non-accreting scenarios have been proposed for unknown CO sources from observations and phenomenology.**
- Interestingly, different pulsar regimes have been proposed in some cases; accreting pulsars, and even magnetars, are possible CO.
- No observational evidence for pulsations, accretion, or determining CO masses, exist beyond hints of a pulsar for some cases.
- Within the class of HMGB, different objects, and even orbital phases, may realize different scenarios.

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- **Non-thermal emission mostly comes from an interacting outflow (caveats: accretion or unshocked pulsar wind features).**
- The outflow structure, and thus all- λ spectra and variability, and morphology (particularly radio), are similarly affected by star+orbit.
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Eccentric, relatively compact Be+CO? binary (end of 70s)

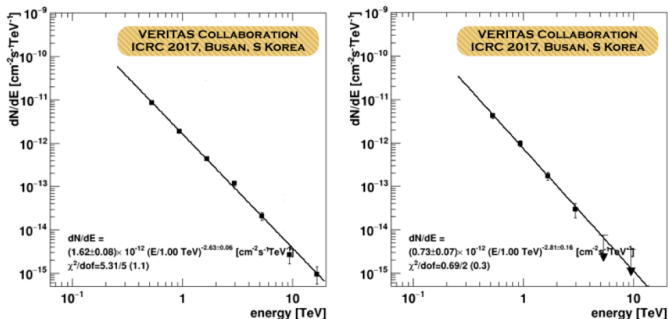
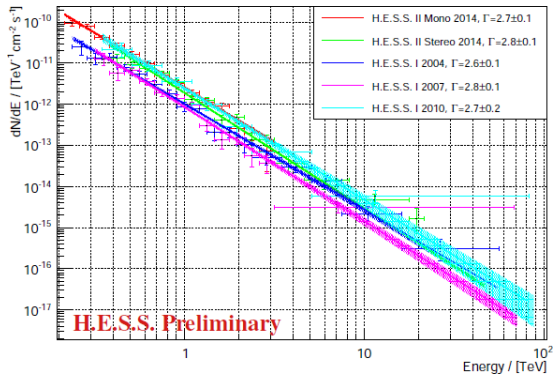


Figure 3: Spectral energy distribution (SED) for LS I +61°303 for two parts of the orbit (parts of the orbit shown on top panels). SED on the *left* is near apastron passage covering $\phi = 0.5 \rightarrow 0.8$ and SED on the *right* is for the rest of the orbit for $\phi = 0.8 \rightarrow 0.5$. The orbital parameters shown on top panel are used from [14]

(VERITAS: Kar et al. 2017)

VHE spectrum of PSR B1259–63

Highly eccentric, wide Be+pulsar binary (beginning of 90s)



(HESS: Bordas et al. 2015)

Moderately eccentric, compact O+CO? binary (90s)

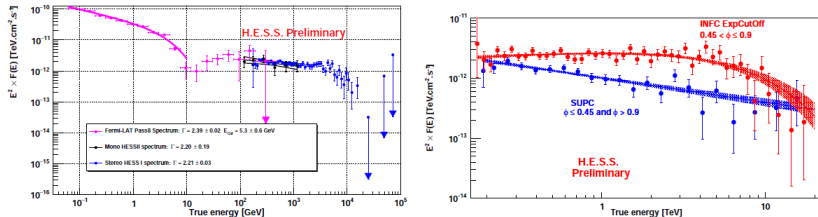
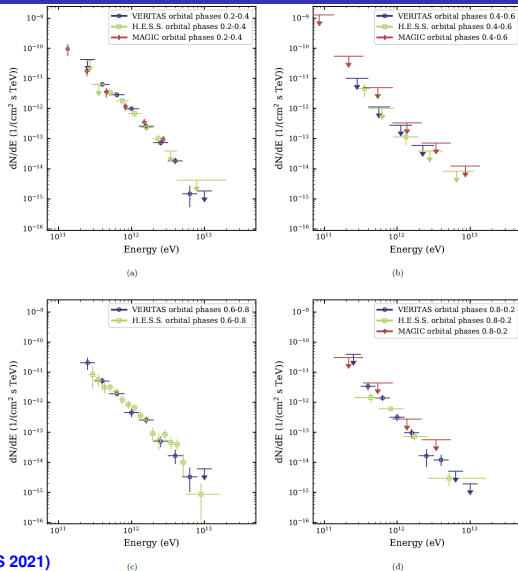


FIGURE 4. *Left:* SEDs obtained from monoscopic and a stereoscopic analyses of the H.E.S.S.-II and H.E.S.S.-I data sets, respectively. Results of fits with power-law functions are given in the inset. Also an SED obtained from a re-analysis of Fermi-LAT data is shown. *Right:* SEDs resulting from H.E.S.S.-I analyses for parts of the orbit corresponding to the inferior or superior conjunction. The corresponding orbital phase ranges are given for reference. Fit results are given in the main text.

(HESS: Bordas et al. 2015)

VHE spectrum of HESS J0632+057

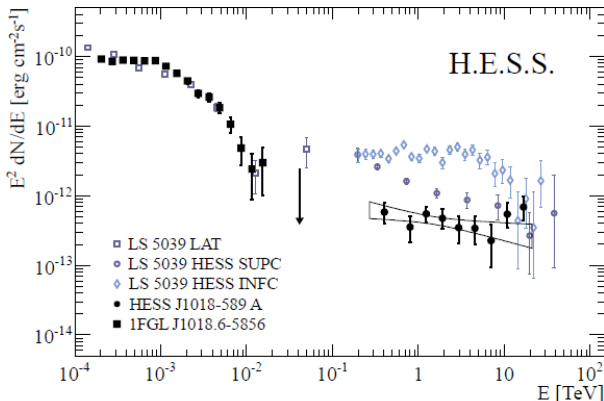
**Eccentric,
rather wide
Be+CO?
binary (00s)**



(HESS, MAGIC & VERITAS 2021)

Figure 7. Differential energy spectra of photons above 200 GeV obtained by H.E.S.S., MAGIC and VERITAS averaged over all available orbits. The figure shows the results for four different orbital phase bins: (a) orbital phases 0.2-0.4; (b) orbital phases 0.4-0.6; (c) orbital phases 0.6-0.8; (d) orbital phases 0.8-0.2. Vertical error bars show 1σ uncertainties; downwards pointing arrows indicate upper limits at the 95% confidence level.

Moderately eccentric?, relatively compact O+CO? binary (10s)

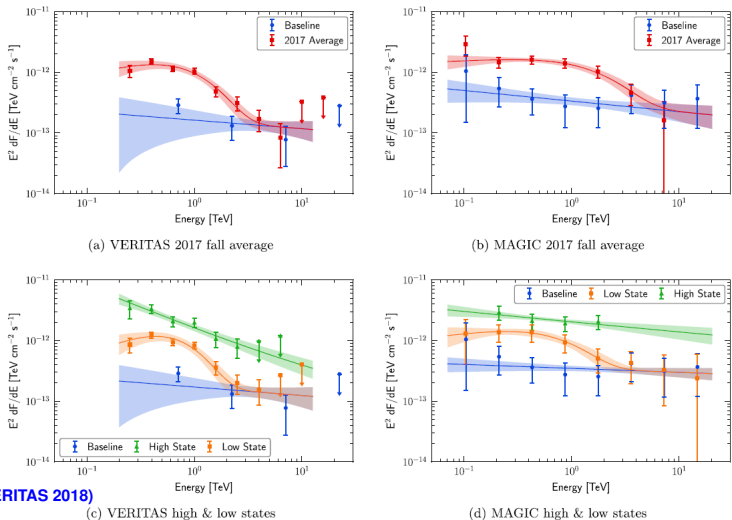


(HESS 2015)

Fig. 1. SED of HESS J1018–589 A/1FGL J1018.6–5856 is shown in black (filled squares and circles for the LAT and HESS detection). For comparison, the SEDs of LS 5039 during superior (SUPC) and inferior conjunction (INFC) are also included (blue points from [Hadasch et al. 2012](#); [Aharonian et al. 2005a](#)).

VHE spectrum of PSR J2032+4127

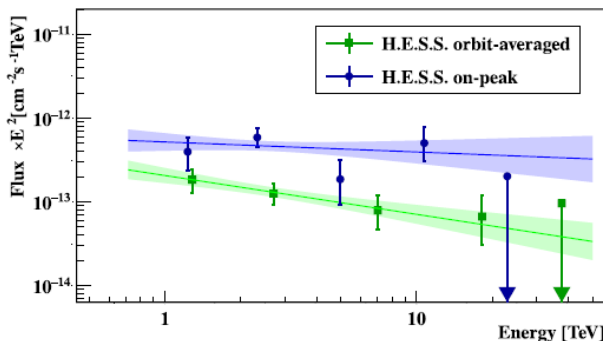
Extremely eccentric, very wide Be+pulsar binary (00s or 10s?)



(MAGIC, VERITAS 2018)

Figure 3. Spectral energy distributions for PSR J2032+4127/MT91 213 and TeV J2032+4130 from VERITAS (left) and MAGIC (right). The blue butterflies are the spectral fits to TeV J2032+4130. The red butterflies in the upper plots are fits to the 2017 fall data: the sum of a power-law fit to TeV J2032+4130 and a cutoff power-law fit to PSR J2032+4127/MT91 213. In the bottom plots, orange is the fit to the low-state data (PSR J2032+4127/MT91 213 is fit with a cutoff), while green represents the high-state data (PSR J2032+4127/MT91 213 is fit with a power law). The fit parameters are given in Table 1 and the time periods are defined in the text.

Moderately eccentric, compact O+CO? binary (10s)



(HESS 2018)

Fig. 3. Spectral energy distribution averaged over the full orbit (green, squares) and for the on-peak orbital phase range (orbital phase from 0.2 to 0.4: blue, circles). The data points have 1σ statistical error bars, upper limits are for a 95% confidence level. The best fit and its uncertainty are represented by the solid lines and shaded areas, respectively.

- **The source is rather similar to LS 5039 and 1FGL J1018.6–5856.**
- It is not known so far if this source emits VHE.
- The HE spectral information does not allow to extrapolate.

(see Corbet et al. 2019)

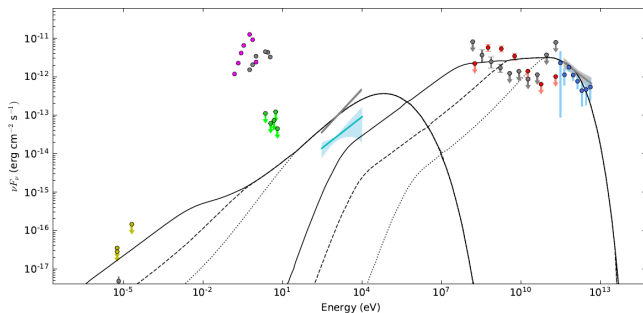
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Moderately wide O?+CO? binary? (10s)



(Martí-Devesa & Reimer 2020 ↑, Tam et al. 2019, and ref. therein; HESS 2015, Eggers et al. 2016)

- 1 Introduction (personal view)
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Non-accreting HMGB: acceleration sites

- **HMGB have plenty of regions where acceleration can occur (via Fermi I, II and shear; B -reconnection; converter mechanism...).**

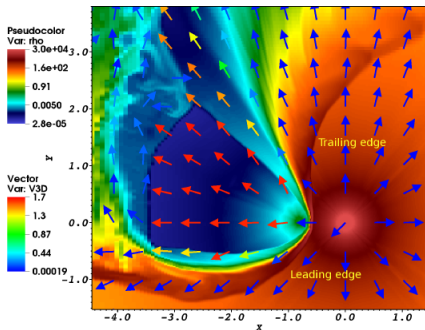
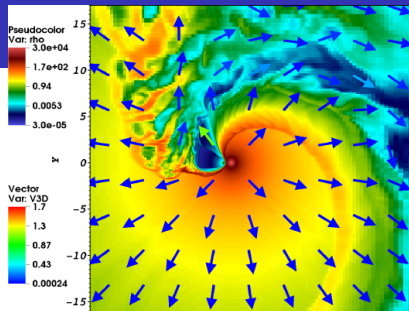
(e.g., Rieger et al. 2007; Khangulyan et al. 2008; B-R & Khangulyan 2009; Takahashi et al. 2009 B-R 2012; B-R & Rieger 2012, Derishev & Aharonian 2012)

Acceleration sites

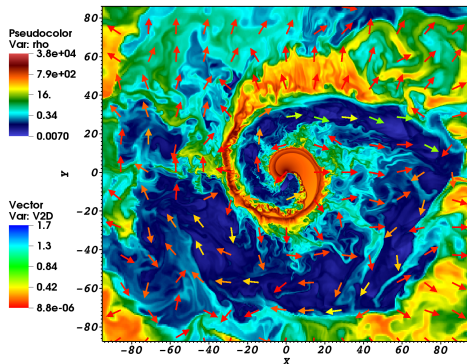
Important elements:

- **Ultrarelativistic weak- B' flow.**
- Perpendicular/oblique shocks of different speeds and strength (B' ?).
- Flow reacceleration and further shocks, shear layers, turbulence and mass-loading... (B' ?)

The next step is to include B .



(B-R et al. 2012 -2D-, 2015 -3D-; low e)

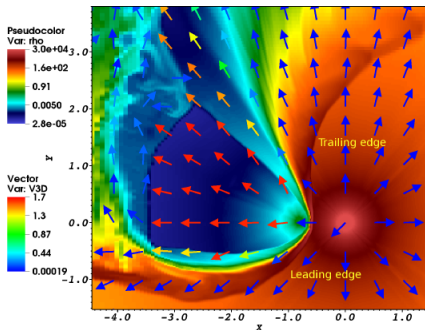
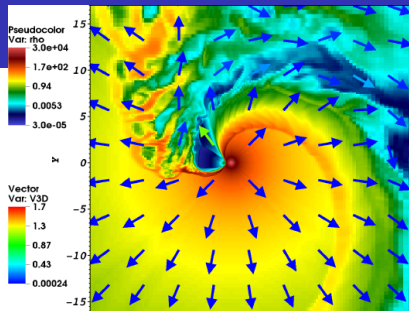


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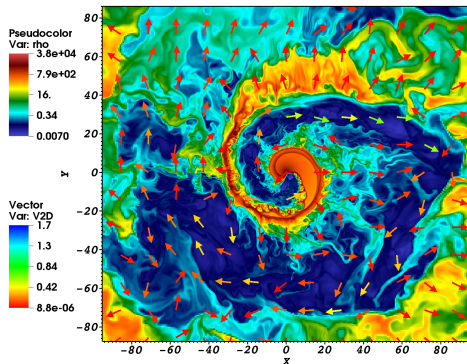
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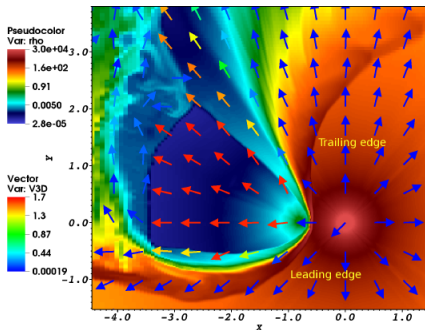
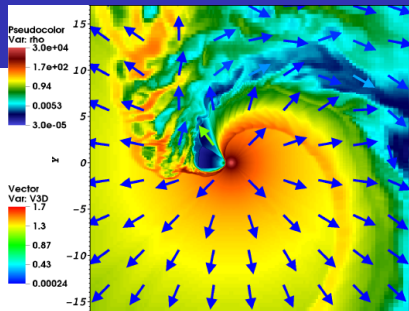


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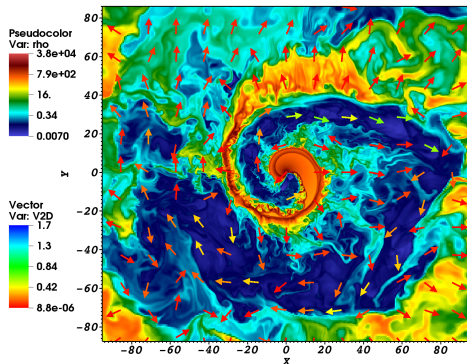
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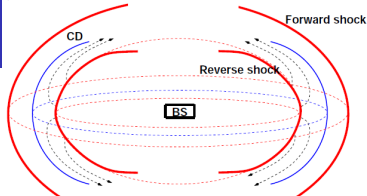
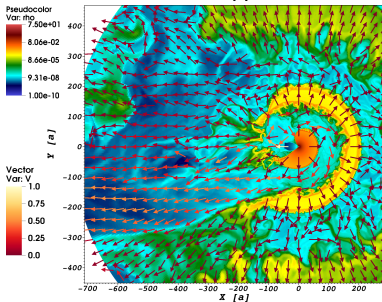
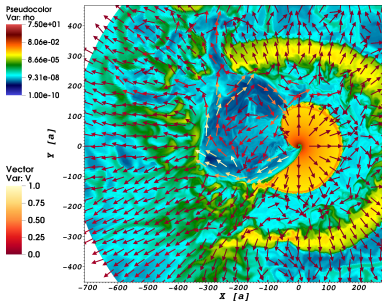
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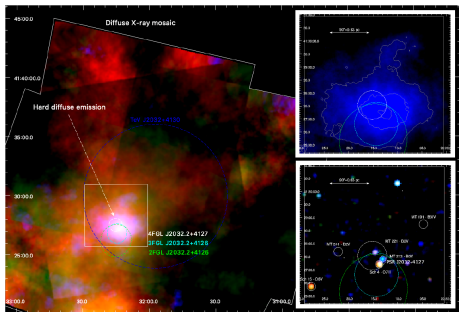


Flow termination.(B-R & Barkov 2011)
Mainly X-ray evidence.

(Paredes+2007 -LSI; Durant+2011 -LS-; Pavlov+2015

-PSRB-, Williams+2015 -1FGL-; Kargaltsev+2021 -HESS-;

Albacete-Colombo+2020 -PSRJ- ↓)



(Barkov & B-R 2018, high e)

Phenomenological E_{\max} in HMGB

- HMGB have plenty of regions where acceleration can occur (via Fermi I, II and shear; B -reconnection; converter mechanism...).

- E_{\max} for the most relevant processes ($t_{\text{acc}} \sim \eta E/qBc$; $D \sim \chi D_{\text{Bohm}}$; $RB \sim ct$):

- Hillas limit (e^\pm, p):

$$E_{\max}^{\text{H}} \sim 300 R_{12} B_0 \text{ TeV}$$

- Escape/adiabatic cooling (e^\pm, p):

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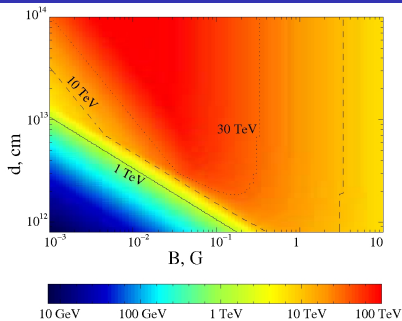
- Diffusion (e^\pm, p):

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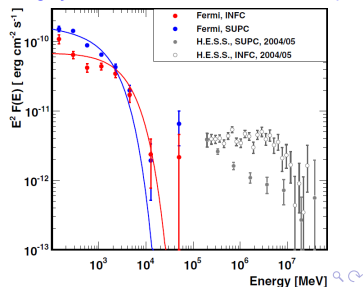
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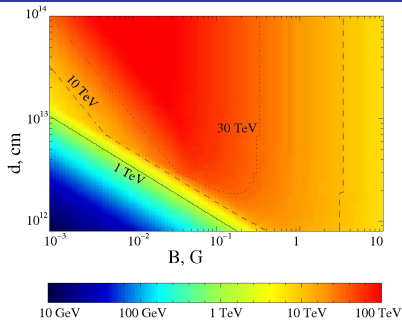
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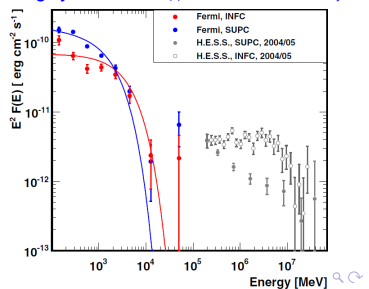
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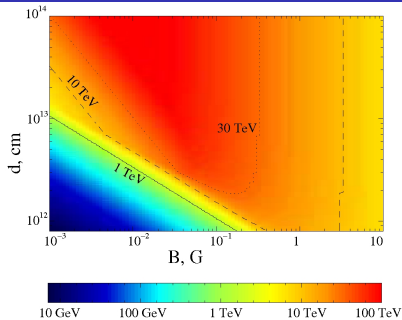
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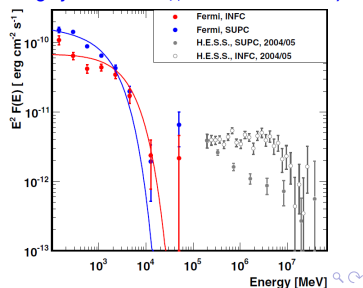
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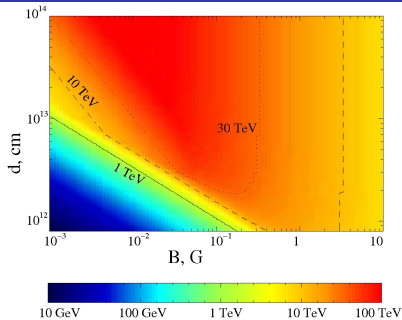
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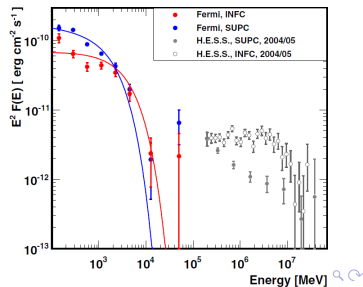
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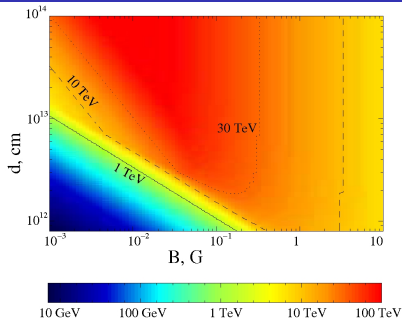
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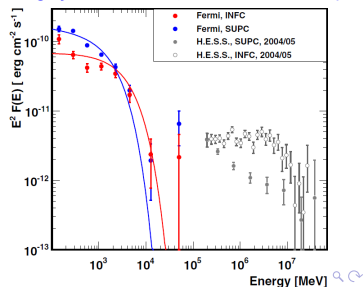
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Unshocked pulsar wind acceleration in a HMGB

- **Derishev & Aharonian (2012) showed that the converter mechanism (Derishev et al. 2003; Stern 2003) can operate in compact HMGB via e^\pm -creation in the unshocked pulsar wind if $\Gamma \gtrsim 10^4$.**
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- Pairs accelerate close to $t_{\text{acc}} = E/qBc$, with $\gamma_{\text{peak}} \sim \Gamma^2 \sim 10^8$.
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(Derishev & B-R, in prep.)

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- **Non-accreting pulsars are attractive but conclusive pulsation detection or mass characterization are needed, as radiatively inefficient BH accretion is still a possibility.**
- HMGB are perfect sites for multi-TeV particle acceleration and radiation:
 - However, leptonic and hadronic CR injection from the system may be inefficient due to mass-loading plus adiabatic losses.
 - On the other hand, large-scale outflow-medium interactions might be a suitable site for PeV CR production.
 - Finally, as LS 5039, LS I 61 040, 1FGL J1018.6–5856, 4FGL J1405.1-6119... are known within $\sim 1/4$ of the disk, the total power of the HMGB population may be $\gtrsim 10^{38} \text{ erg}^{-1}$.

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