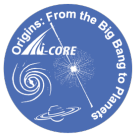


# Multimessenger Emission from GRBs

Tsvi Piran

Elly Leiderschneider, Ofek Birnholtz, Ehud Nakar, Evgeny Derishev  
Shotaro Yamasaki, Matteo Pais



CDY Sept 1 2021



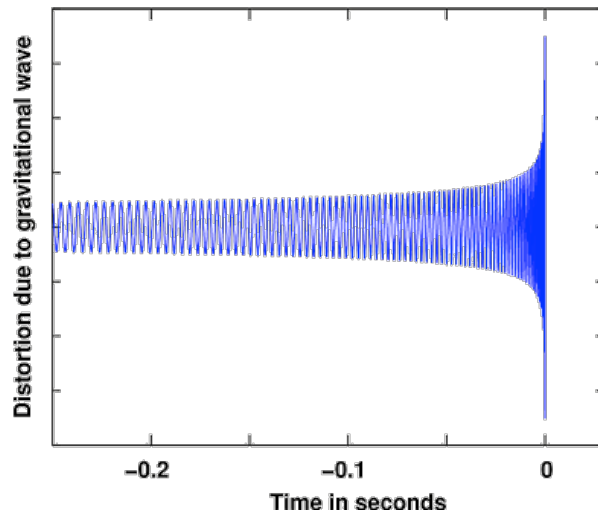
# Outline

- GW from GRBs
- Jet-GWs from Jets: Birnholtz & TP PRD, 2013; Leiderschneider & TP [arXiv:2107.12418](https://arxiv.org/abs/2107.12418)
- Hidden Jets: TP + [arXiv 1704.08298](https://arxiv.org/abs/1704.08298); [ApJ, 2019](https://doi.org/10.1086/70000)
- TeV emission from GRBs Afterglows : Derishev & TP [ApJL, 2019](https://doi.org/10.1086/70000), [arXiv: 2021.12035](https://arxiv.org/abs/2021.12035)

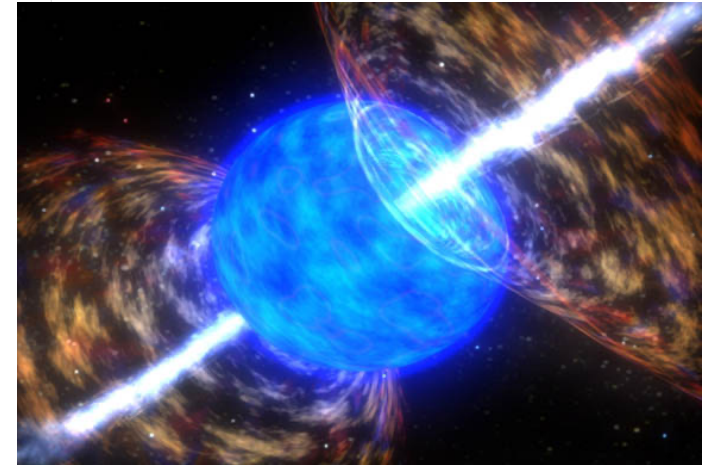
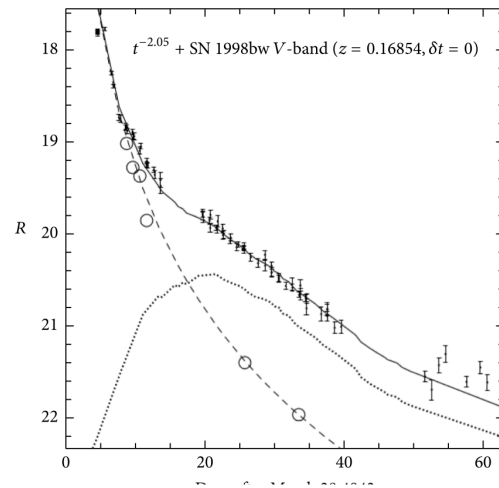
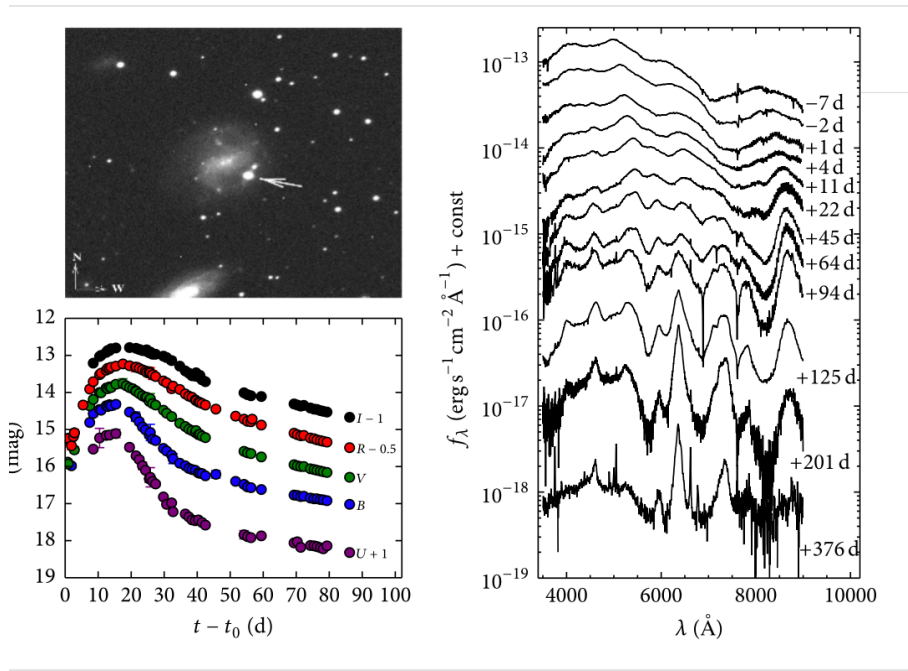
# Binary Neutron Stars Gravitational Waves

$$P = \frac{dE}{dt} = -\frac{32}{5} \frac{G^4}{c^5} \frac{(m_1 m_2)^2 (m_1 + m_2)}{r^5}$$

$$t_{merge} \approx 30 \text{ sec} (P / 30)^{-8/3} (M / M_{sun})^{-5/3}$$



# Supernova and Long GRBs



98bw GRB 980425

Collapsars

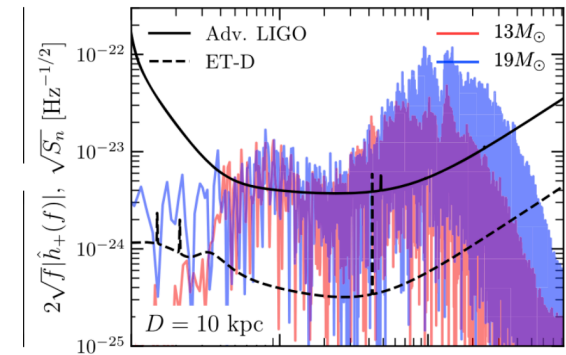
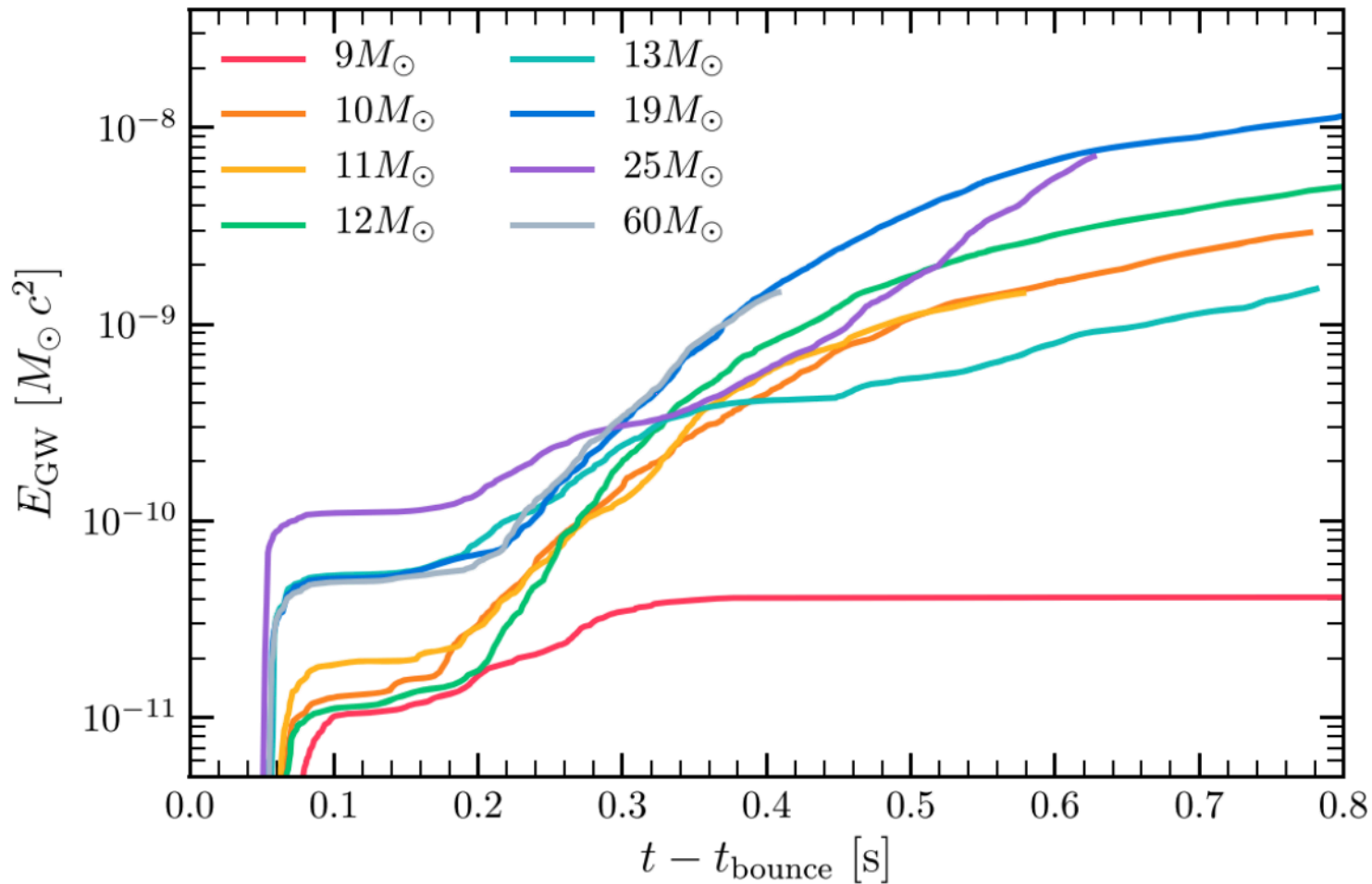
# GW from Supernova

$$h = \epsilon \frac{GM}{c^2 d} = \epsilon \frac{r_g}{d} \quad \nu = \sqrt{G\rho}$$

**Once upon a time (late 70ies)**

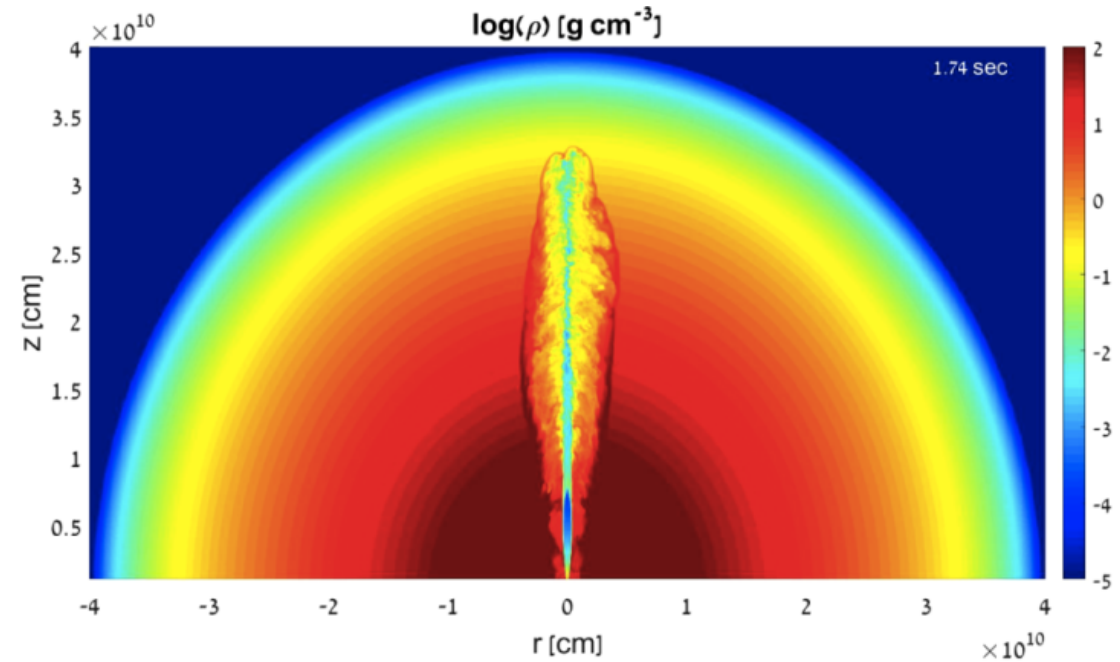
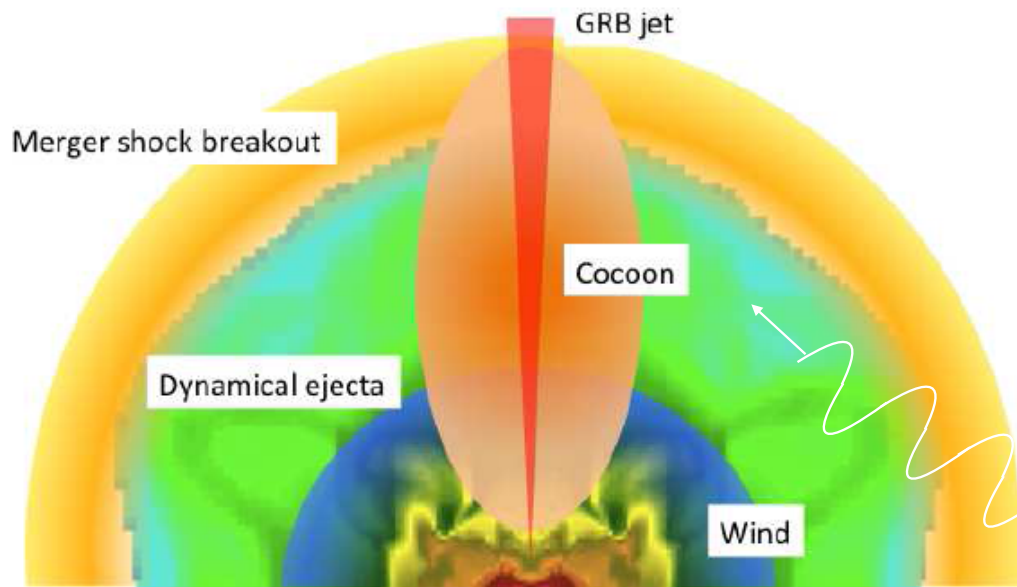
$$h=10^{-21} @ \text{ kHz} \Leftrightarrow E_{\text{gw}}=10^{51} \text{ erg} @ 10 \text{ Mpc}$$

# GW from Supernova



From  $\sim 10 \text{ kpc}$

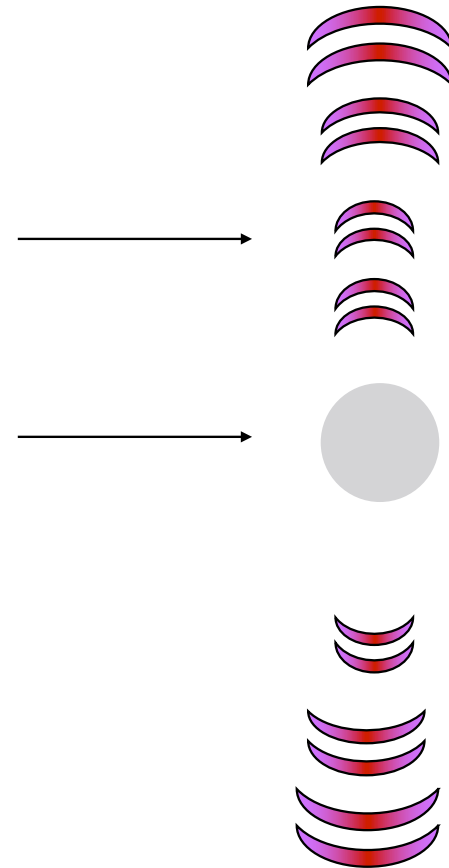
# More Gravitational Waves



Gravitational waves from the jet acceleration process  
(TP 2002; Birnholtz & TP 2014; Liedershneider & TP 2021  
Segalis & Ori 2001)

Jet  
 $\sim 10^{50}$  erg  
Poynting flux or particles?  
Opening angle  $\sim 5^\circ$

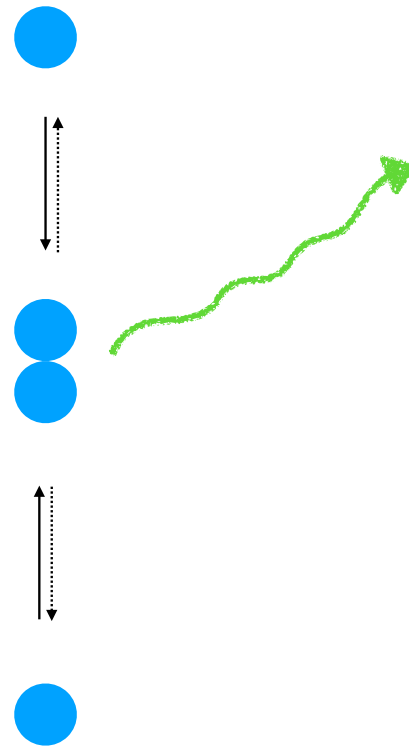
Central Engine  $10^6$  cm



**$10^{50}$  erg were standing still and then they suddenly they move at  $c$   
 $\Rightarrow$  This must produce gravitational waves**



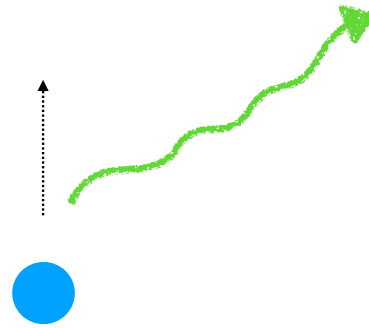
# Weinberg - GW from particle collisions



The sun emits  $10^{15}$  ergs/sec in GW (Weinberg)

The typical energy of the sun's gravitons is  $\sim$  keV with frequency of  $10^{18}$  Hz

Only 1 in  $\sim 10^{25}$  scatterings produces a graviton

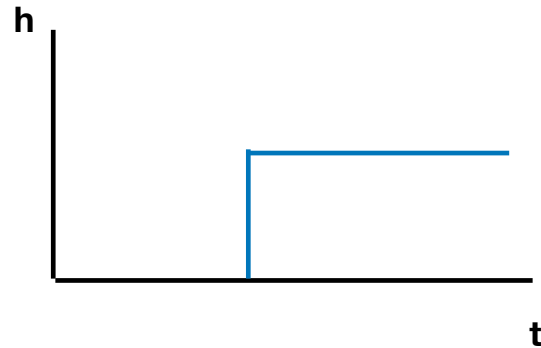


A “quarter” of the collision give us the results of an instantaneously accelerated particle.

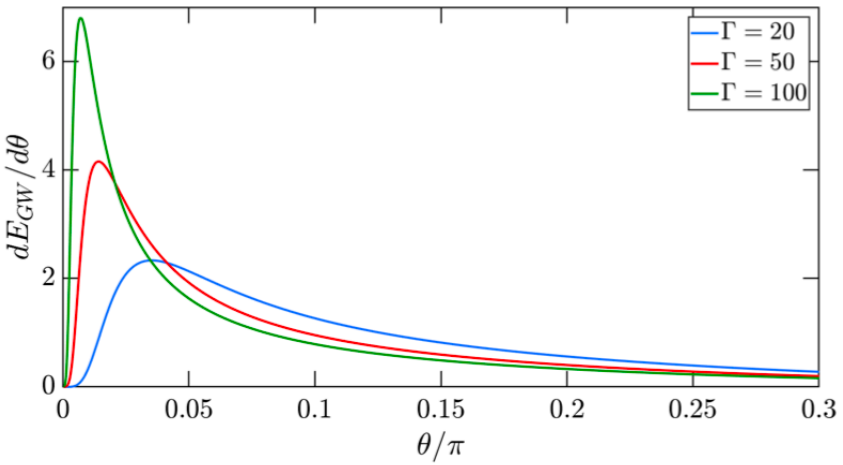
# The ZFL approximation

(Instantaneous acceleration)

$$h^{TT}(\theta_v) = h_+ + ih_x = \frac{2\mathcal{E}\beta^2}{r} \frac{\sin^2 \theta_v}{1 - \beta \cos \theta_v} e^{2i\phi}.$$



$$h_{\max} = \frac{4\mathcal{E}}{r}, \quad \text{at} \quad \theta_{\max} = \sqrt{2/\Gamma}.$$



Anti-beaming

**γ-rays**

**Acceleration  
to  $\Gamma \gg 1$**

**γ-rays**

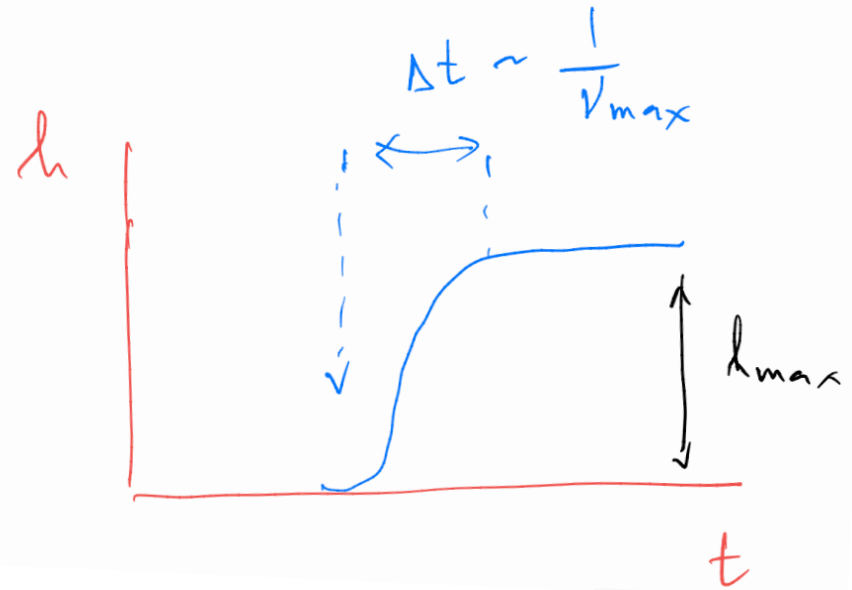
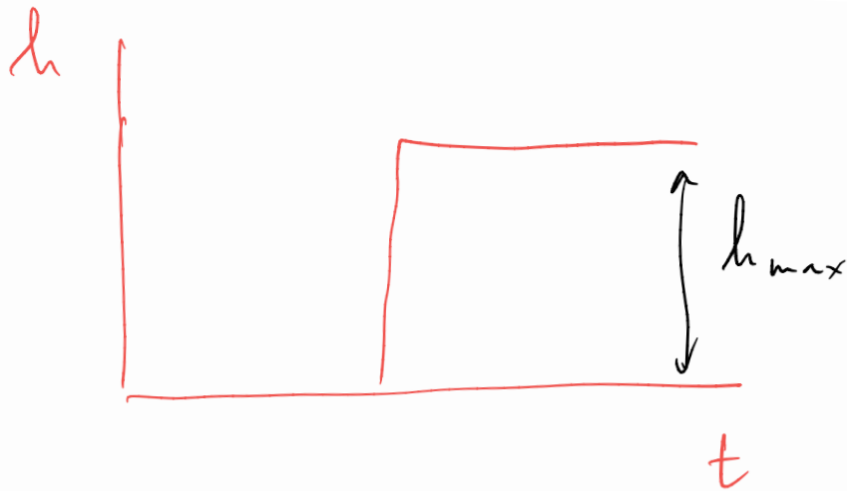
$$E_{GW} = \frac{2GE^2}{\pi c^5 \delta t} \approx 10^{48} \text{ ergs}$$

$$0 < \nu < \delta t^{-1} \approx 100 \text{ Hz}$$

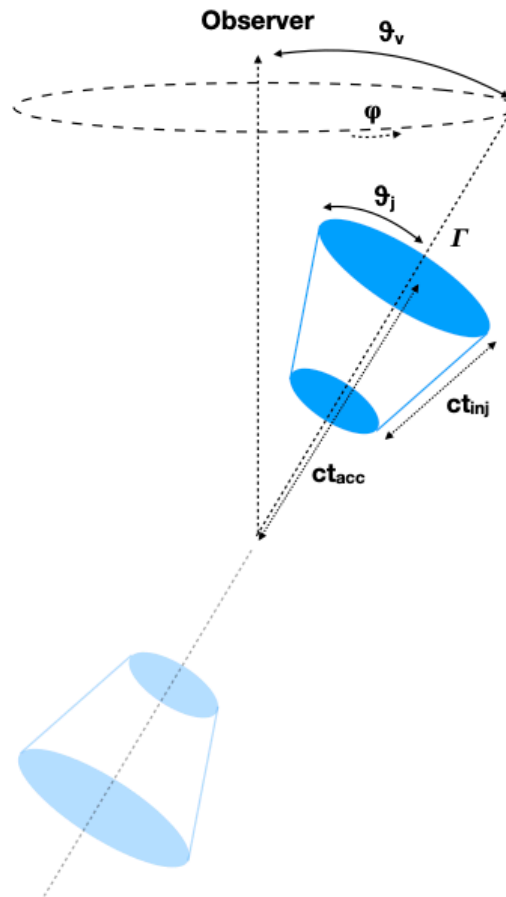
$$h \sim 10^{-25} \text{ for } D \sim 100 \text{ Mpc}$$

This slide is from my 2001 talk at GRG16

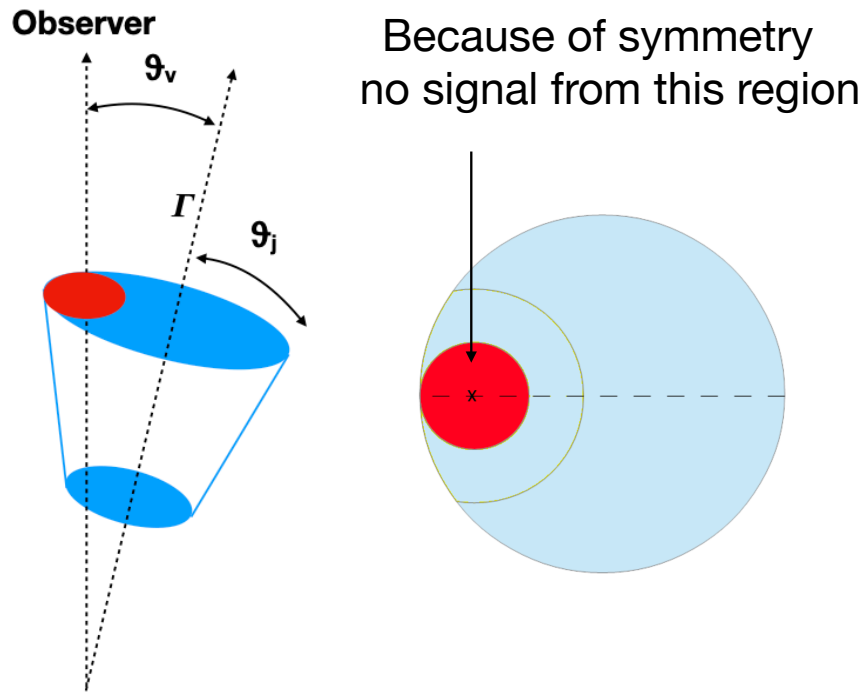
# Realistic light curve



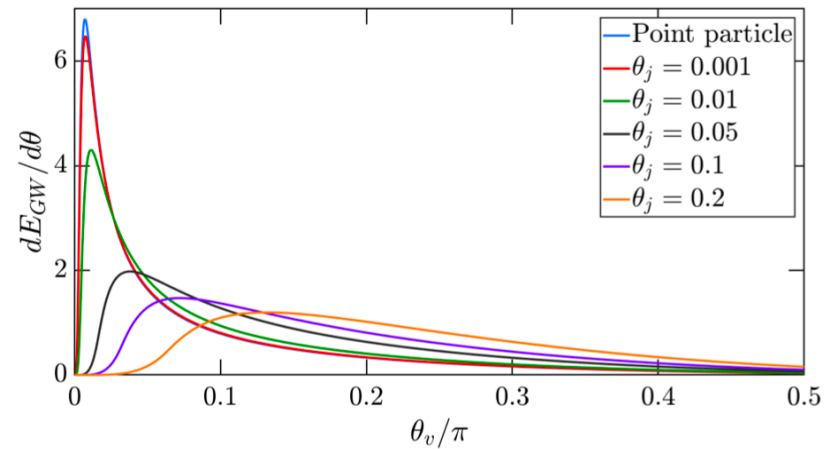
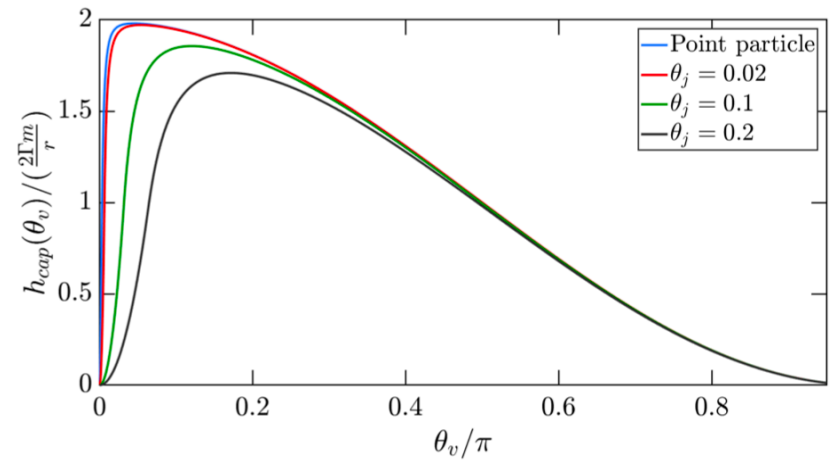
# A more realistic configuration



# A jet

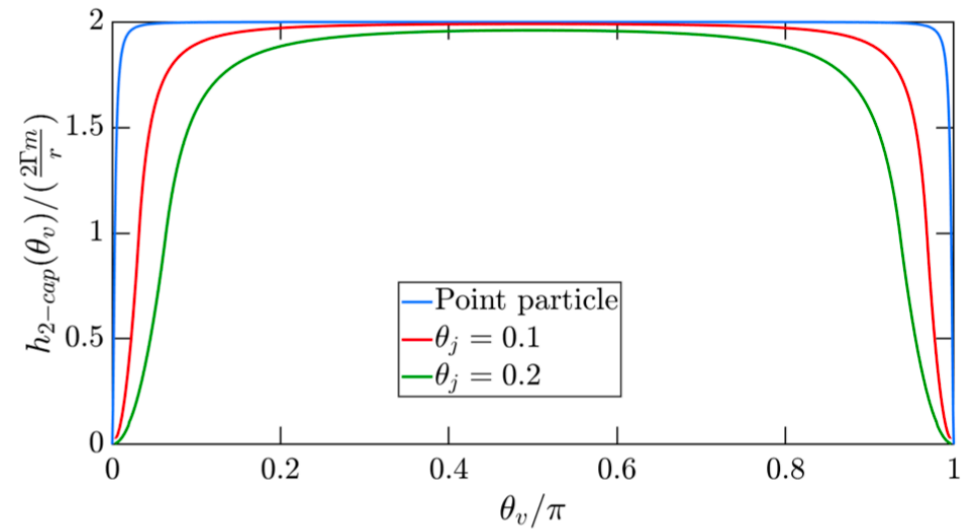
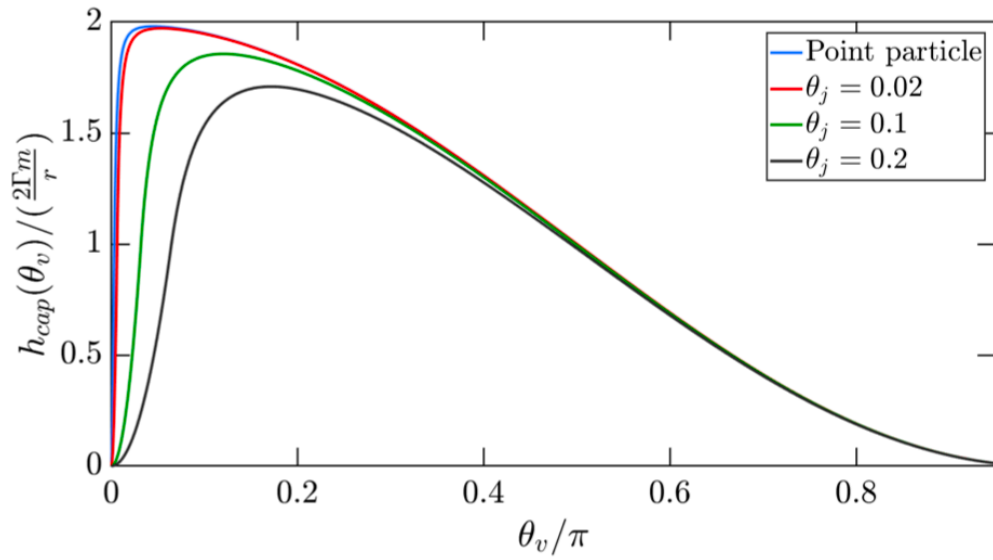


Antibeaming is defined by  $\vartheta_j$

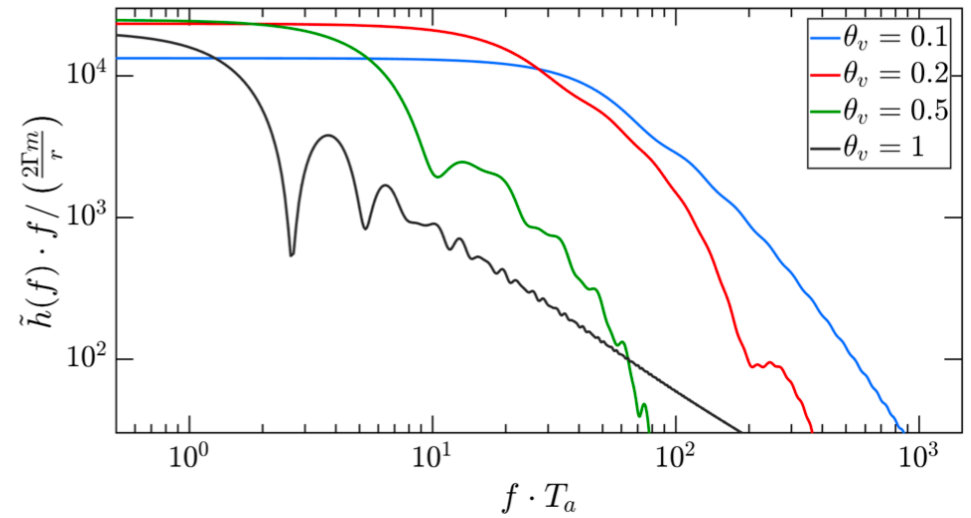
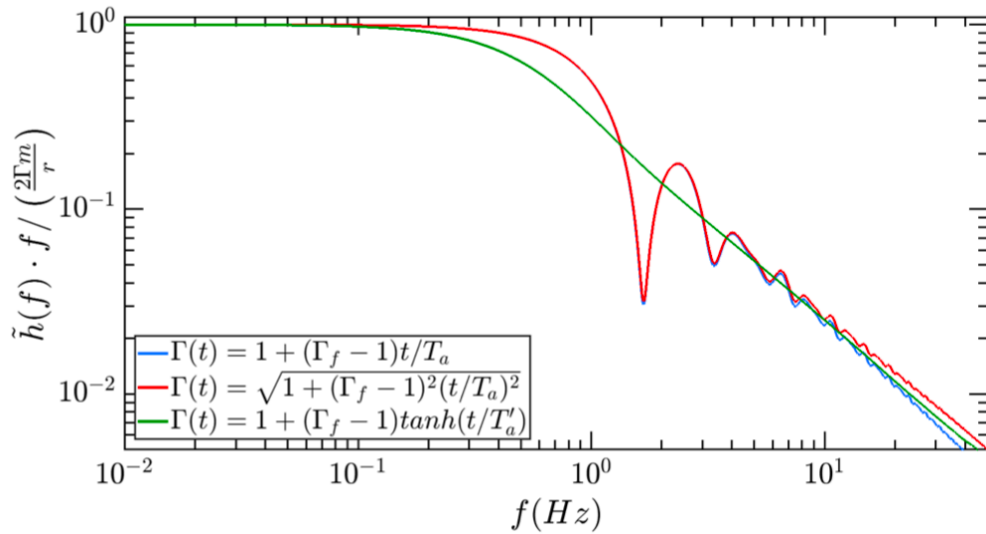




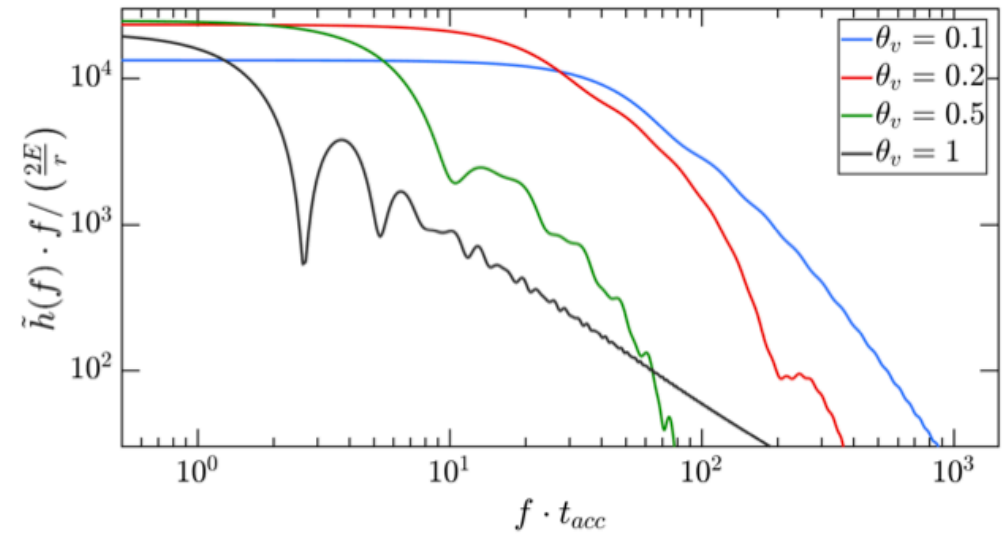
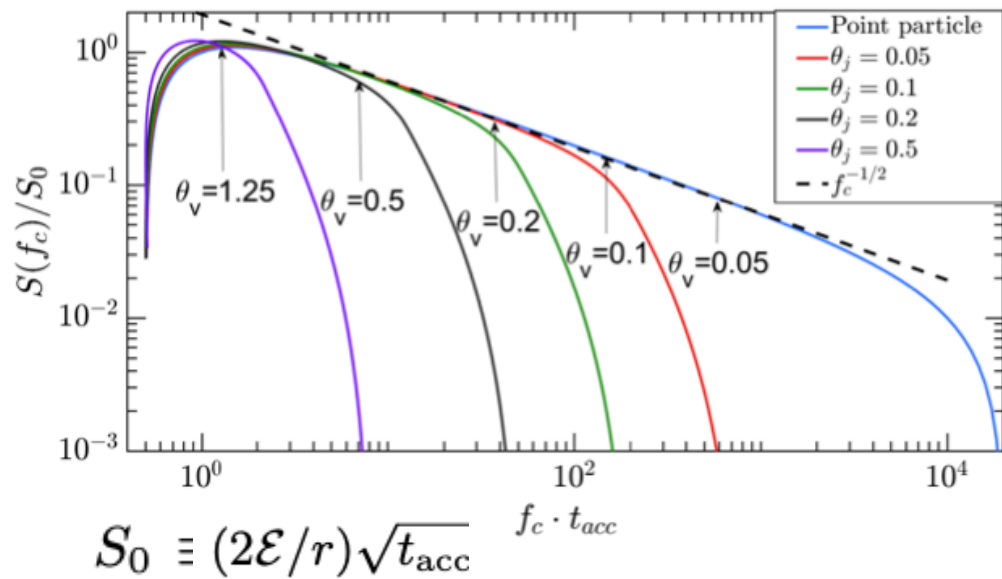
# Two sided jets



# Different acceleration models

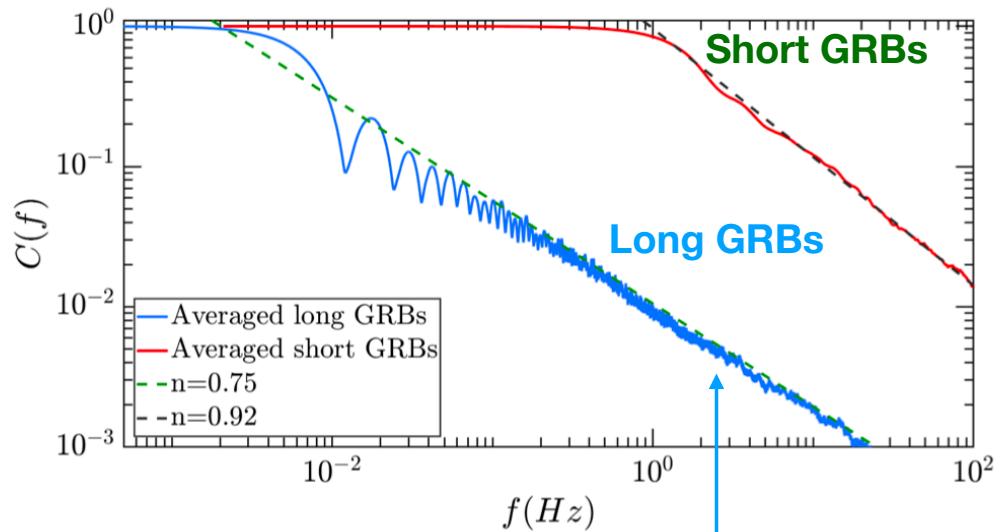


# The crossover frequency

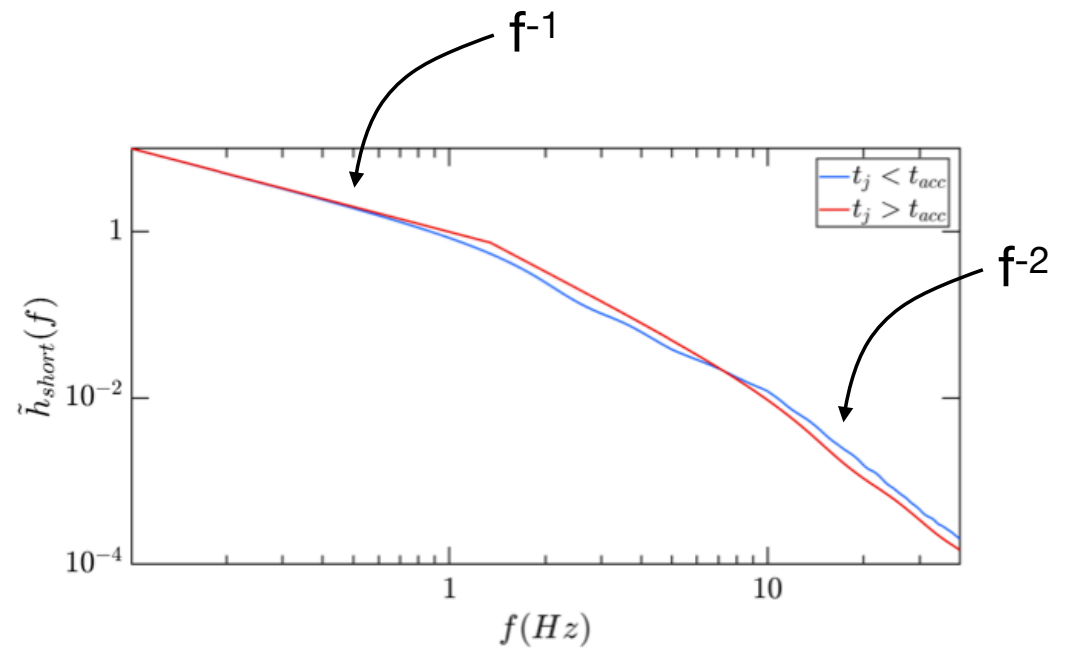


A transition frequency from  $f^{-1}$  at low frequency to  $f^{-\alpha}$  with  $\alpha > 3/2$  at higher frequencies

# GW from GRB jets



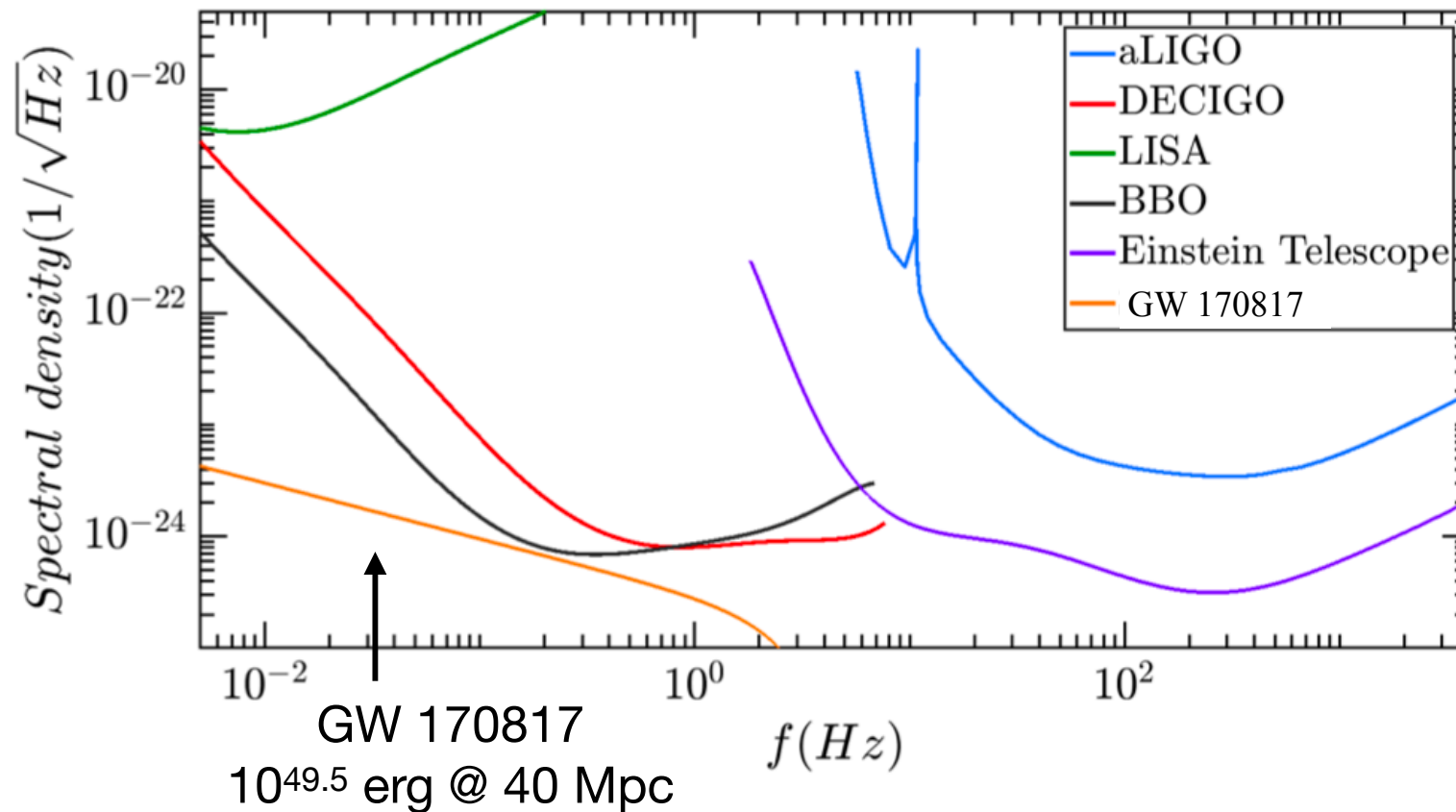
Beloborodov 2000



H for short GRBs

Assuming that the observed GRB power spectrum reflects the jet output

# GW from the jet of GRB 170817



Low frequency  $\lesssim 1$  Hz and weak  $< 10^{-24}$  but possibly detectable with BBO, and DECIGO

**Are there better sources?**

# SN 1997ef

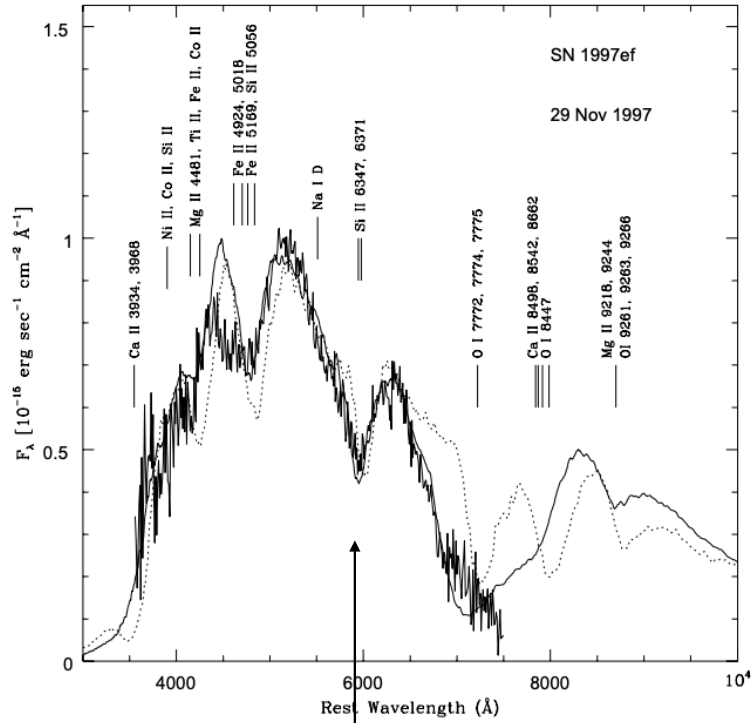


FIG. 2.— The observed, smoothed spectrum of SN 1997ef on Nov 29, 1997 (thick line), compared to two synthetic spectra computed for  $t = 9$  days. The dashed line is a spectrum computed with the original model CO100, while the fully drawn thin line is a spectrum computed with the modified outer density described in the text.

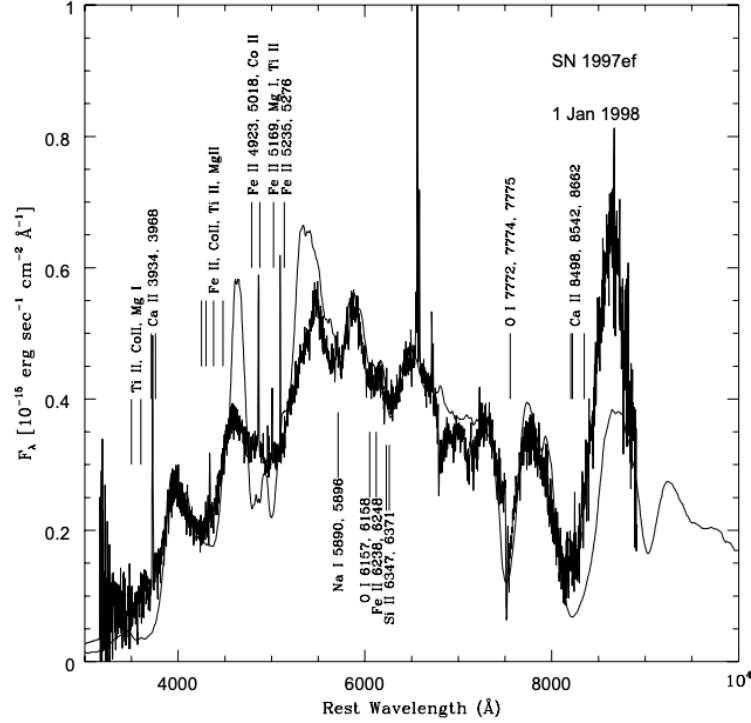


FIG. 6.— The observed spectrum of SN 1997ef on Jan 1, 1998 (thick line), compared to a synthetic spectra computed for  $t = 42$  days using the modified density described in the text (thin line).

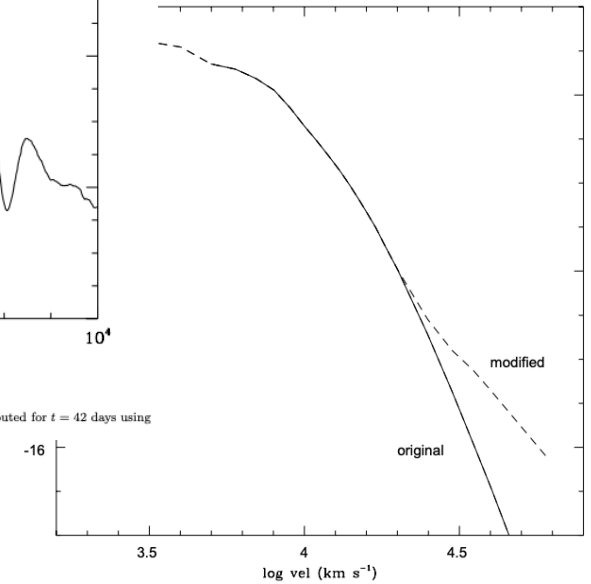
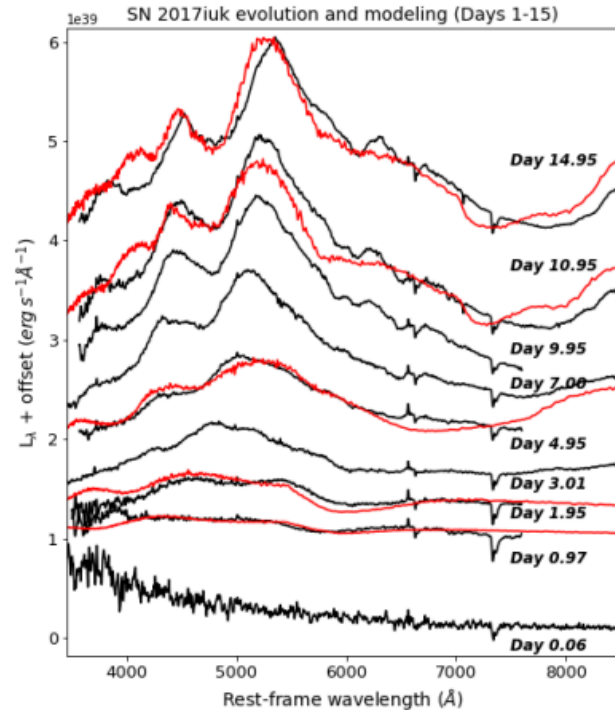


FIG. 8.— The original density structure of the hydrodynamical model CO100 (continuous line) and the modifications introduced to improve the spectral fits (dashed line). The outer part of the modified density structure has a power law index  $n = -4$ , while the inner extension has  $n = 1$  between  $v = 3000$  and  $v = 5000$  km s $^{-1}$  and  $n = -1$  below  $v = 3000$  km s $^{-1}$ .

Mazzali et al., 2000

Very broad absorption lines disappear at later spectrum

# SN 2017iuk (GRB 171205A)



**Figure 2.** The spectral evolution of SN 2017iuk during the first 15 days after the GRB. All spectra are shown as black curves, and they have been de-reddened for Galactic extinction, with the GRB afterglow contribution being subtracted. The simulated emission (red curves) obtained from our synthesis model for some selected spectra are shown as red curves. For the spectral simulation at Day 0.957 an arbitrary constant has been considered, due to the uncertainty in the afterglow component continuum, to match the observed data.

epoch spectroscopic observations of SN 2017iuk, associated with GRB 171205A which display features at extremely high expansion velocities of  $\sim 100,000 \text{ km s}^{-1}$  within the first day after the burst<sup>4,5</sup>. These high-velocity components are characterized by chemical abundances different from those observed in the ejecta of SN 2017iuk at later times. Using spectral synthesis models

Very broad absorption lines disappear at later spectrum



From density distribution  
to energy distribution =>  $v^5$

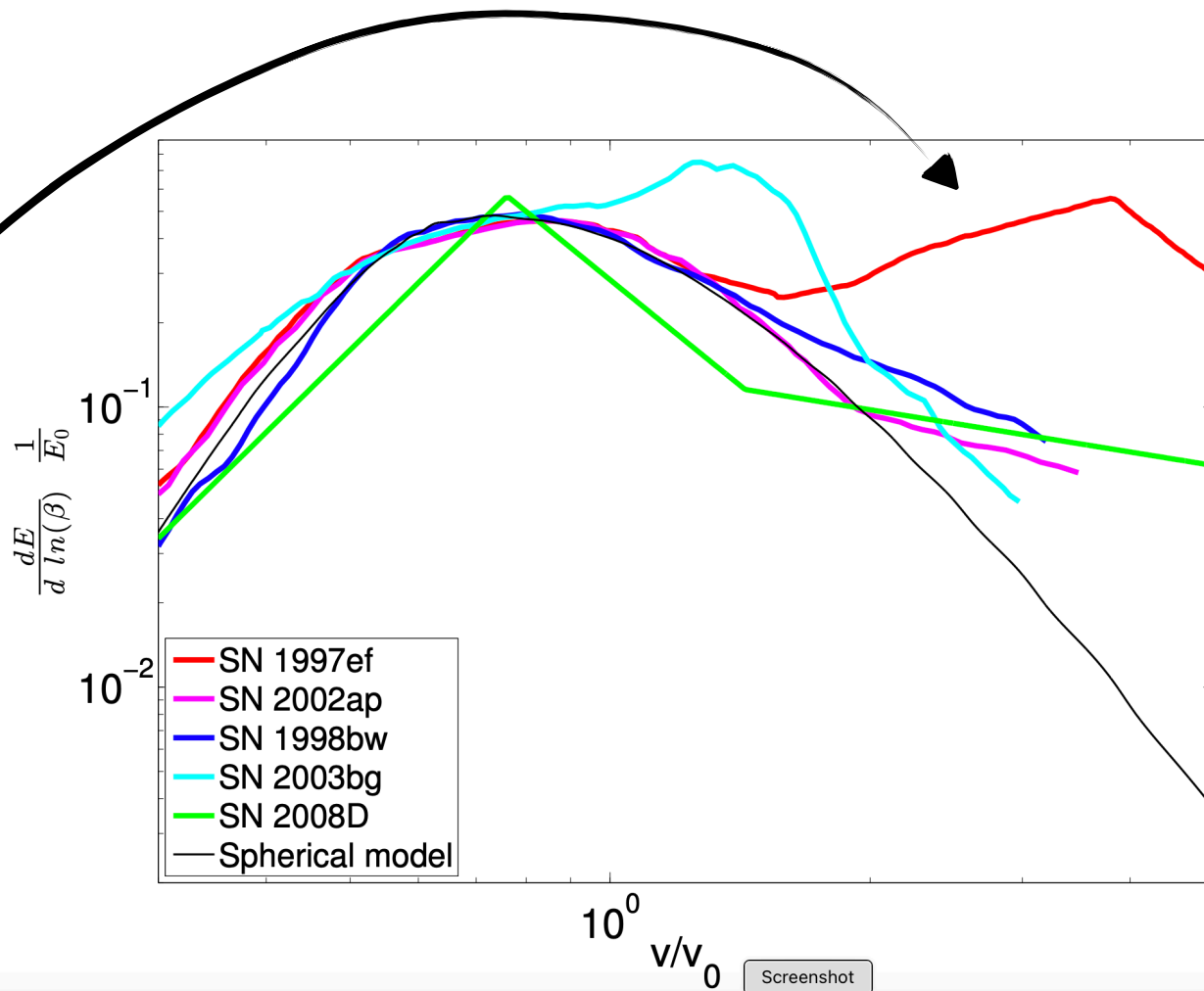
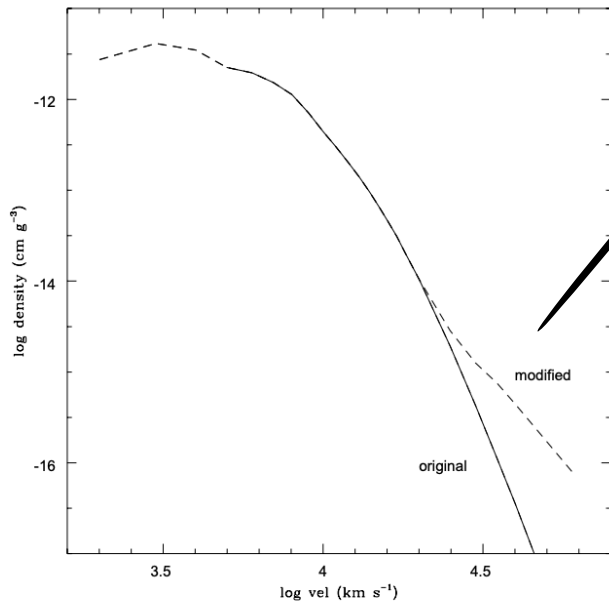
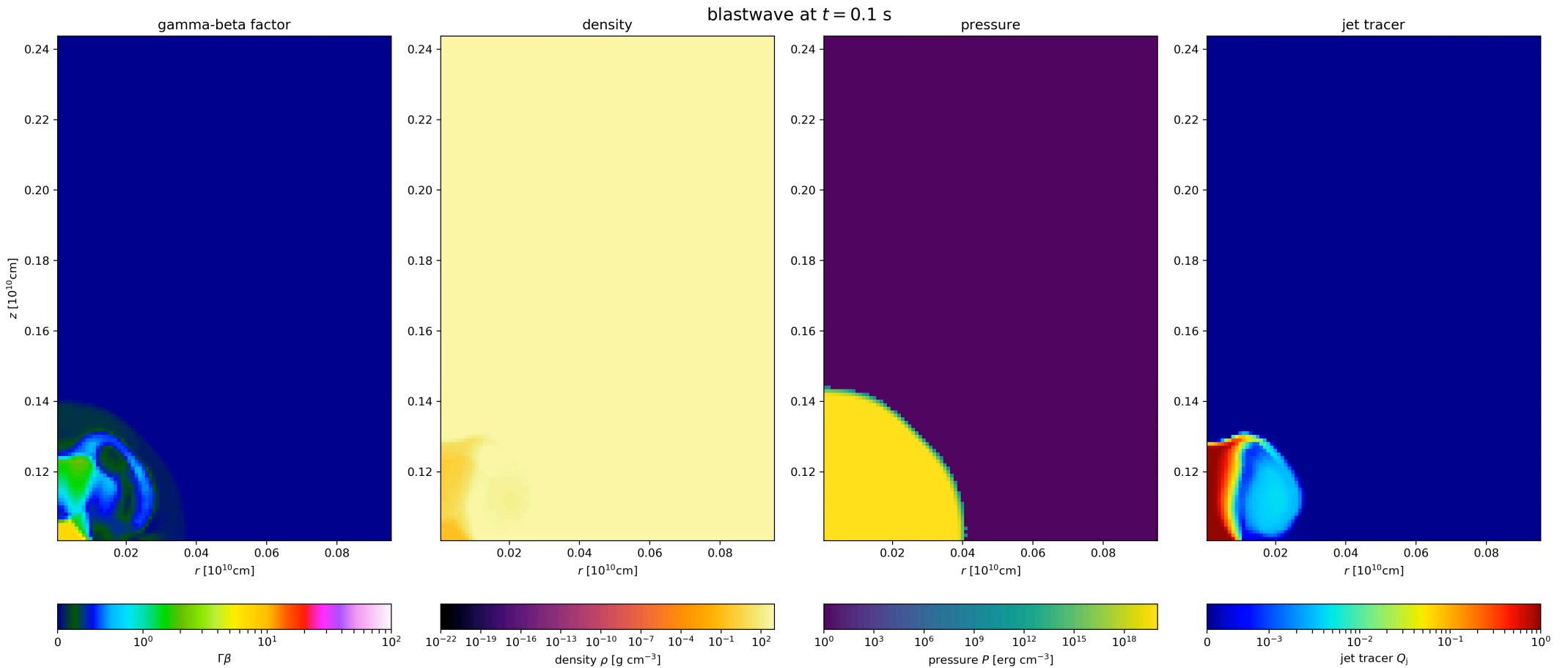


FIG. 8.—The original density structure of the hydrodynamical model CO100 (continuous line) and the modifications introduced to improve the spectral fits (dashed line). The outer part of the modified density structure has a power law index  $n = -4$ , while the inner extension has  $n = 1$  between  $v = 3000$  and  $v = 5000 \text{ km s}^{-1}$  and  $n = -1$  below  $v = 3000 \text{ km s}^{-1}$ .

Mazzali et al., 2000

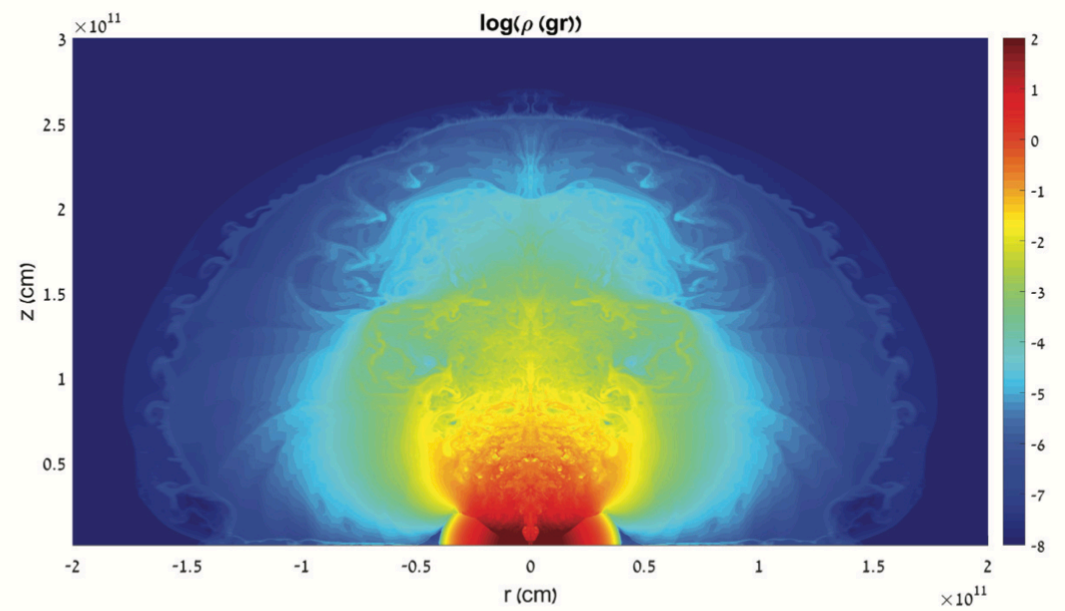
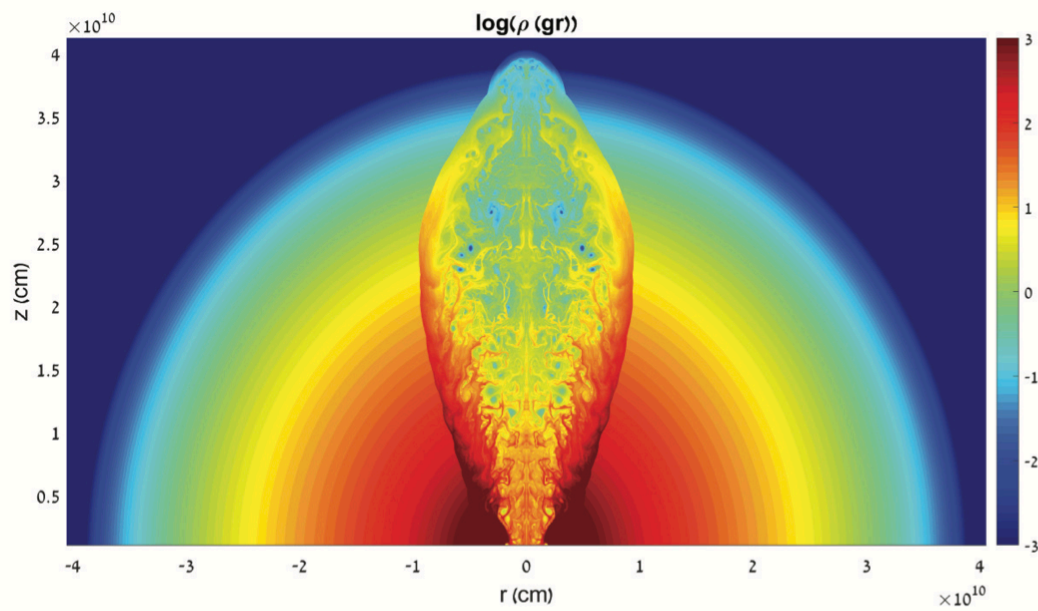
TP, Nakar, Mazzali & TP 2017, 2019

# Choked jet within a Star

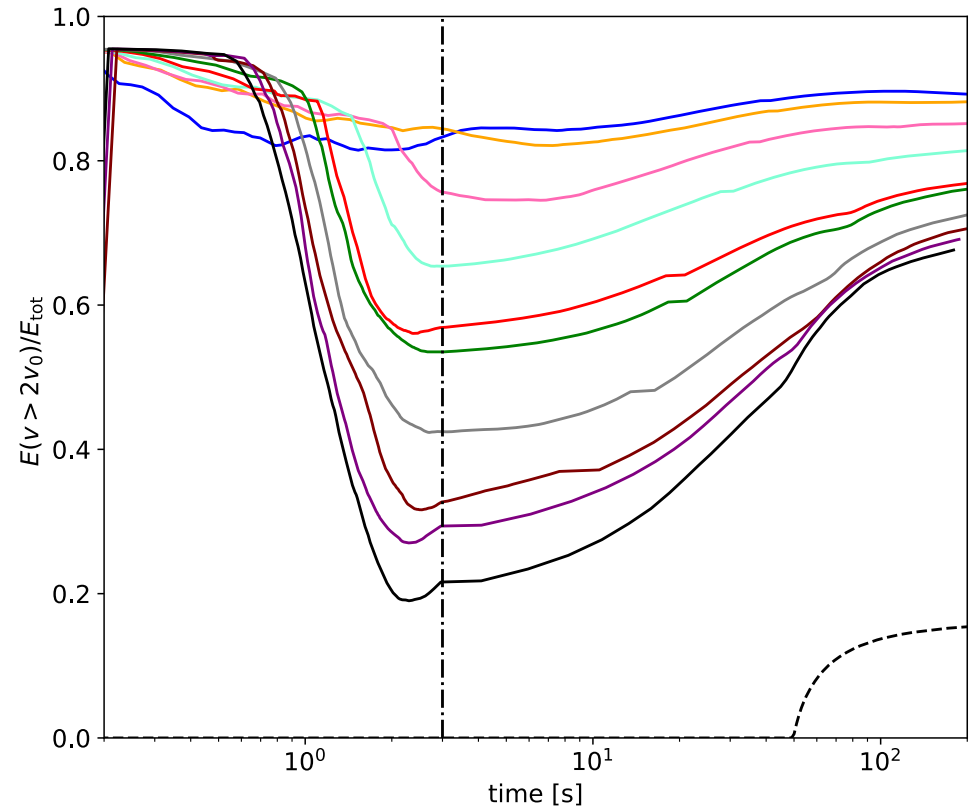
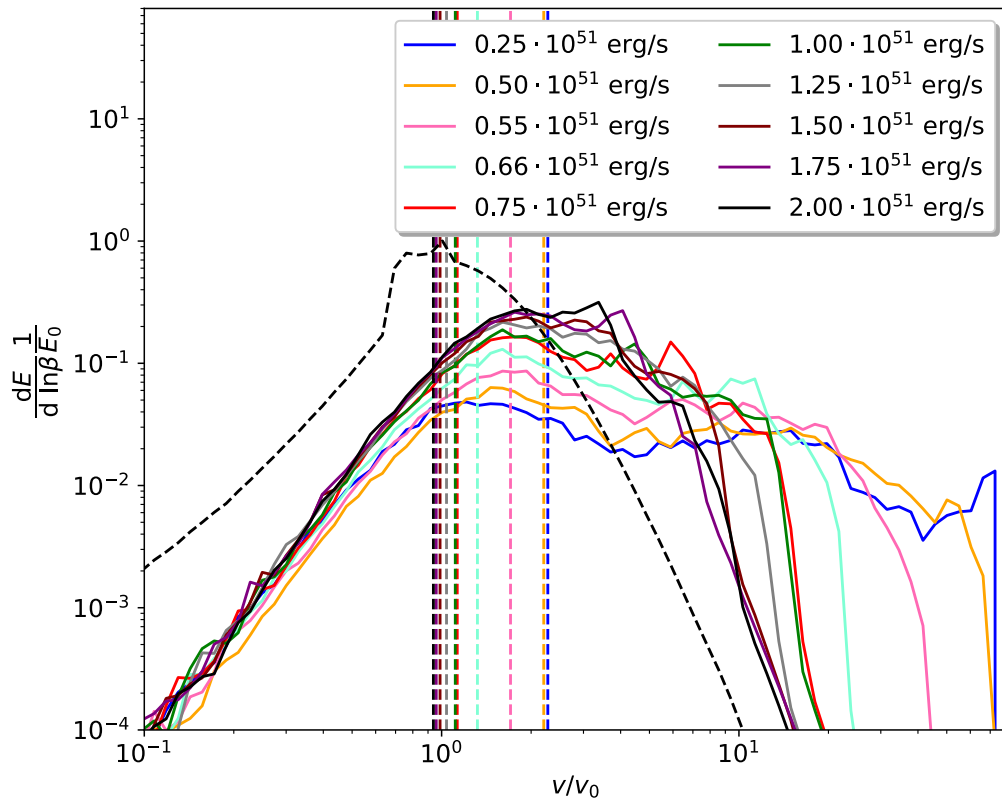


Credit: Matteo Pais

# SNe harbor energetic jets



# The energy distribution



Credit: Matteo Pais

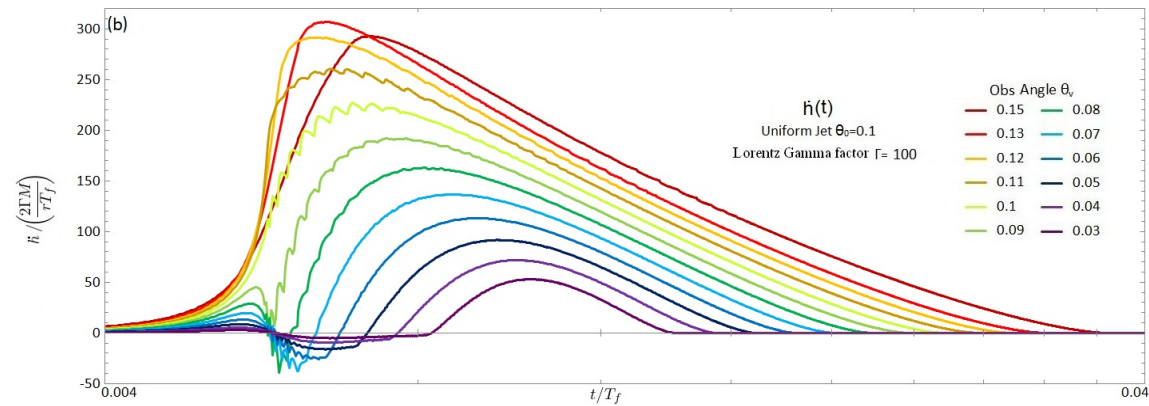
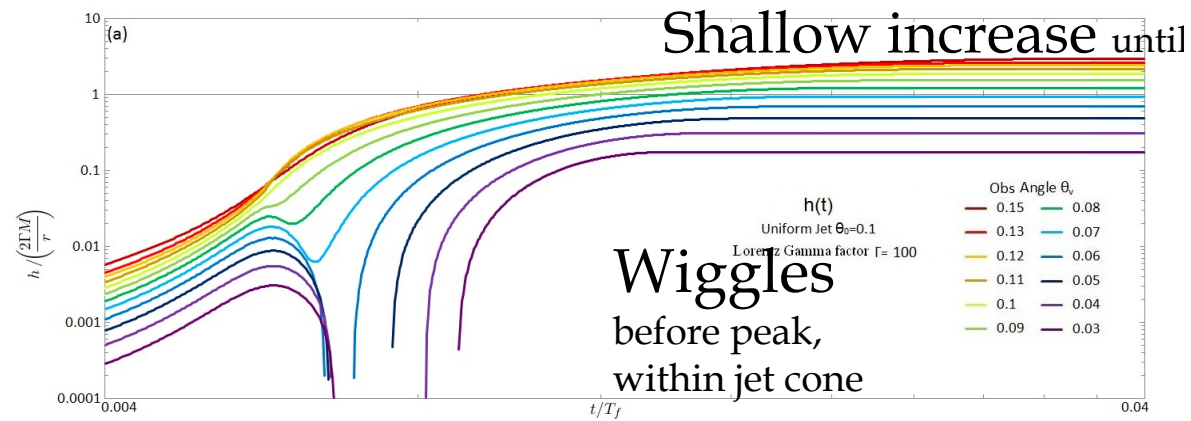
# Some SNe and their Jets

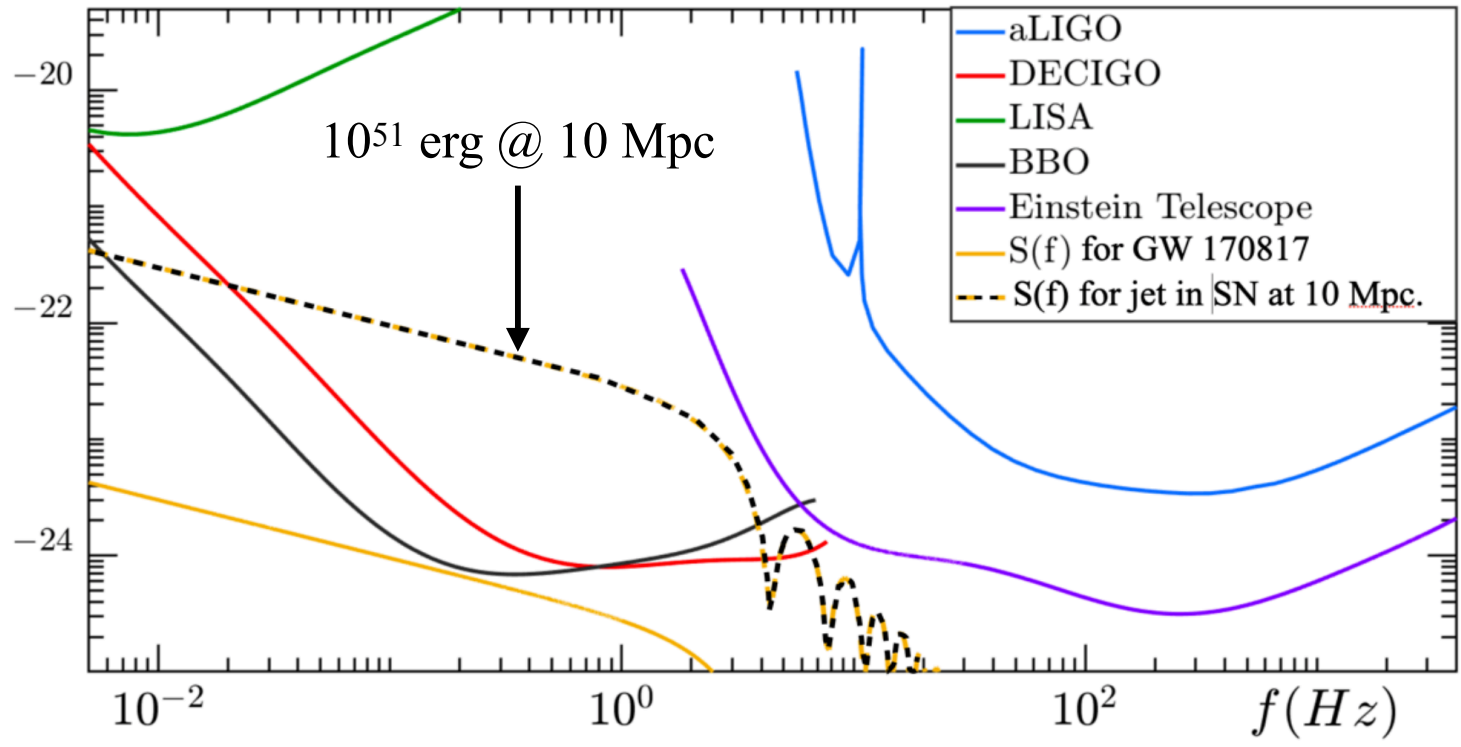
SN	Type	$E_{tot}$ [ $10^{51}$ erg]	$M_{ej}$ [ $M_{\odot}$ ]	$E_j$ [ $10^{51}$ erg]	$M_c$ [ $M_{\odot}$ ]	$\theta_c$ [deg]	Comments	ref.
1997ef	Ic-BL	20	8	9	0.4	$20^{\circ}$	No associated GRB	[16]
1998bw	Ic-BL	50	11	$\gtrsim 2$	-	-	Associated with a low luminosity GRB 980425	[17]
2002ap	Ic-BL	4	2.5	0.3	-	-	No associated GRB. No outflow faster than 0.3c.	[18]
2003bg	I Ib	5	4.5	1	0.2	$20^{\circ}$		[19]
2008D	Ib	6	7	1.4	-	-	Associated with a faint x-ray burst	[20]
2016jca	Ic-BL	50	10	$\gtrsim 2$	-	-	Associated with a long GRB 161219b	[21]

All the SNe are stripped, some associated with //GRBs

# GW from the jet

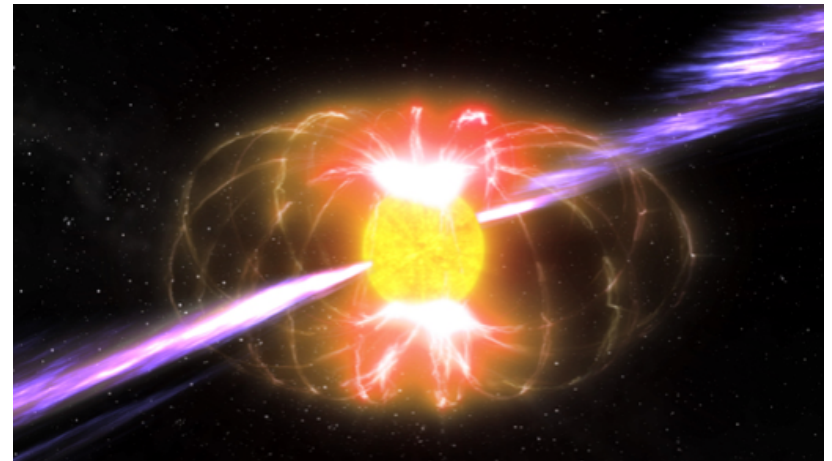
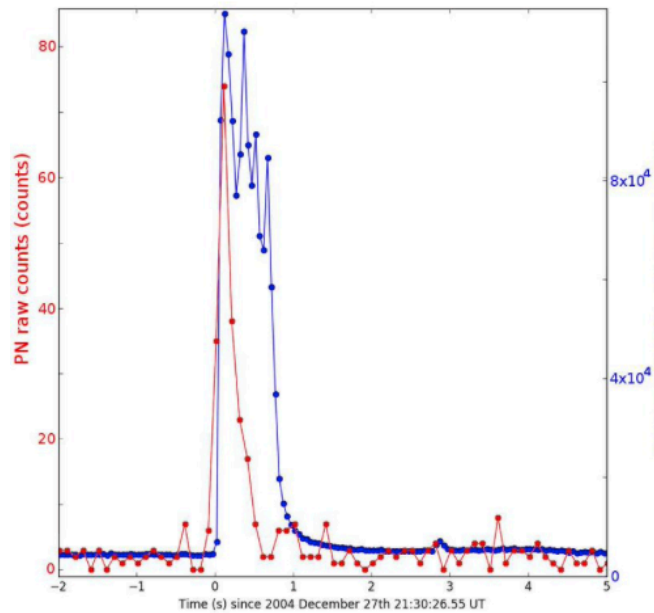
(Birnholtz & TP, 14)





Low frequency  $\approx 1$  Hz and mild  $< 10^{-23}$  but detectable with BBO, DECIGO and marginally Einstein Telescope

# SGR giant flares?



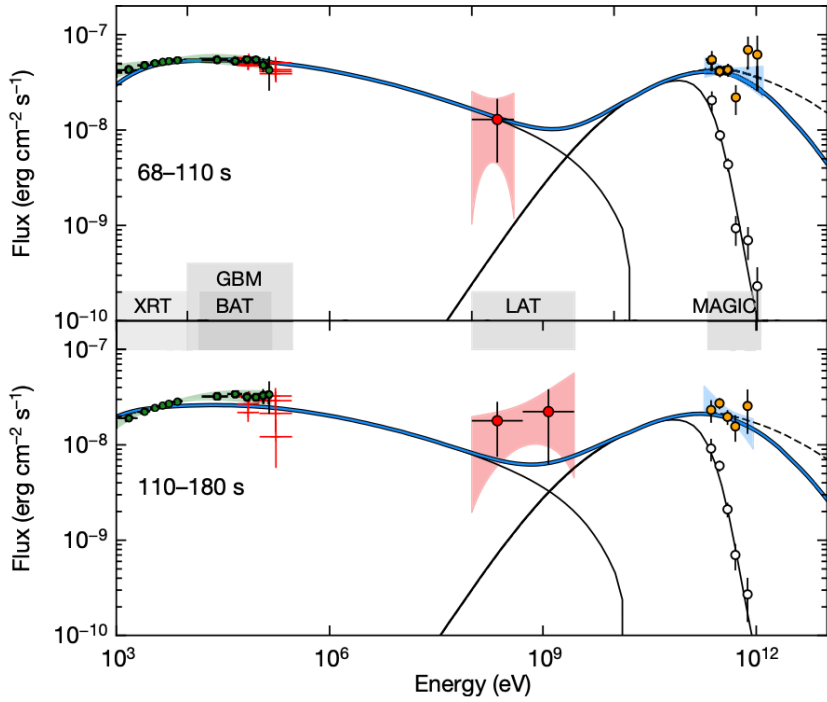
$10^{47}$  erg in a few millisecond. Is this good enough ?



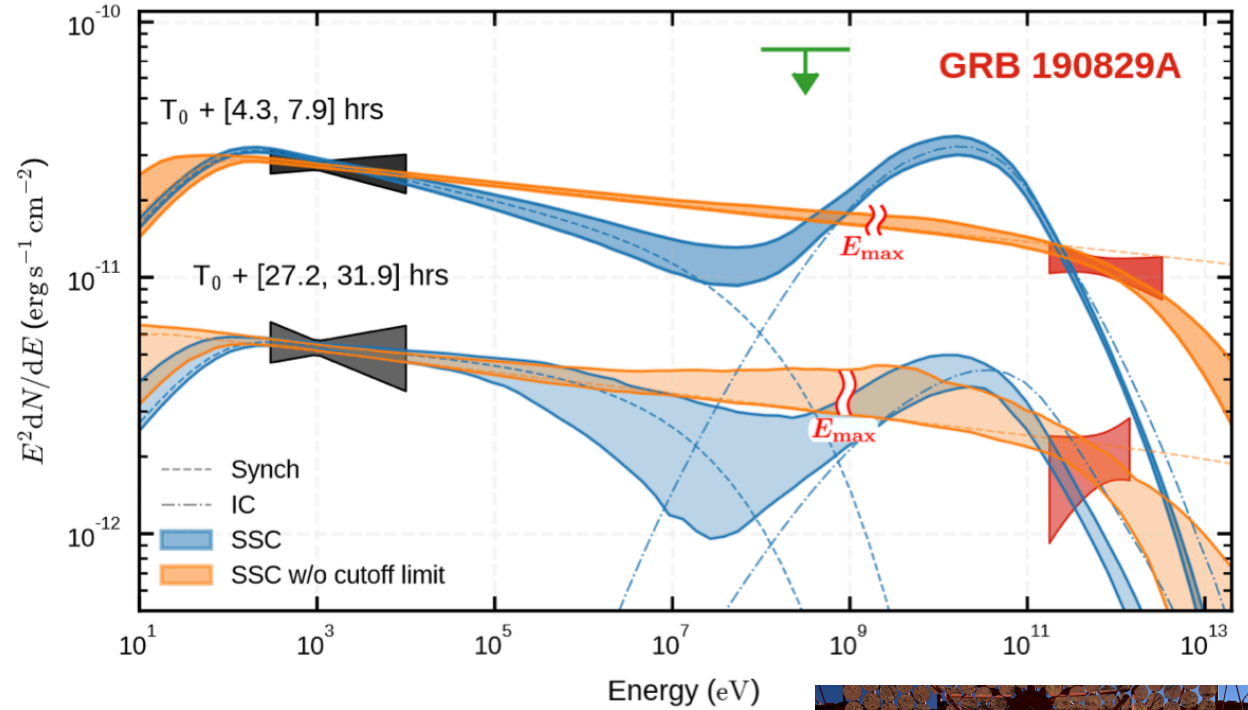
# Summary

- Jet are sources of GW
- The GW could provide excellent diagnostic of the acceleration process
- However, GRBs are most likely too distant to detect their jets even with planned detectors.
- However, hidden jets in SNe might bring us back to detection of SNe from ~10 Mpc

# TeV



190114C-Magic

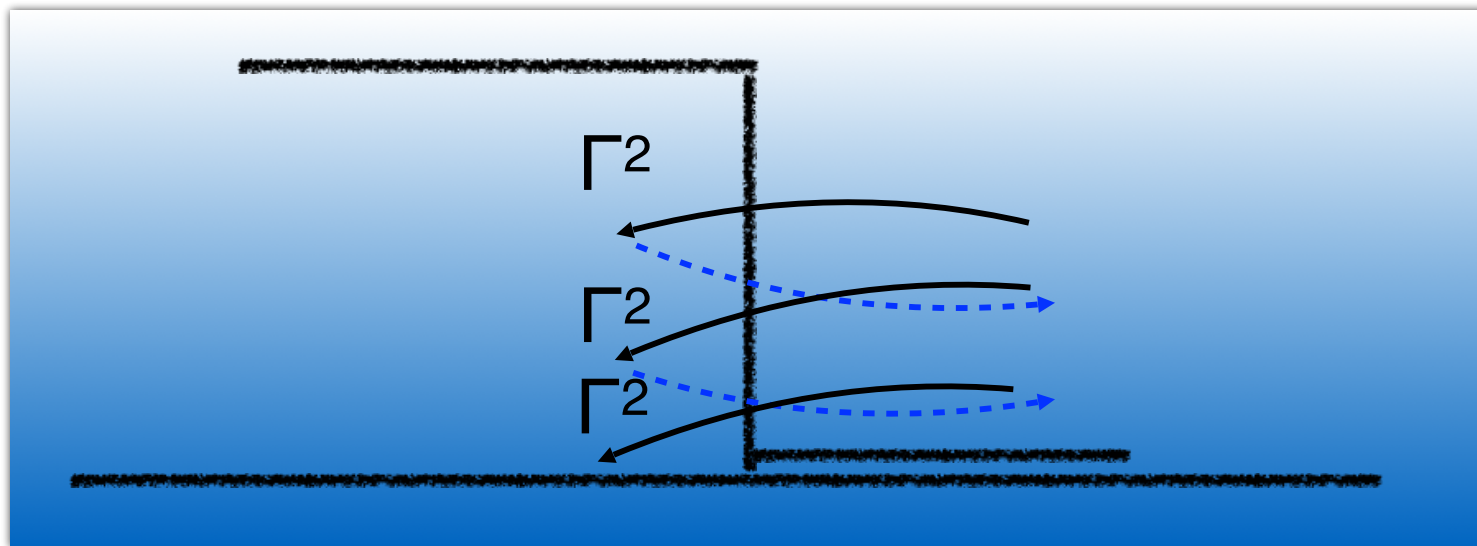
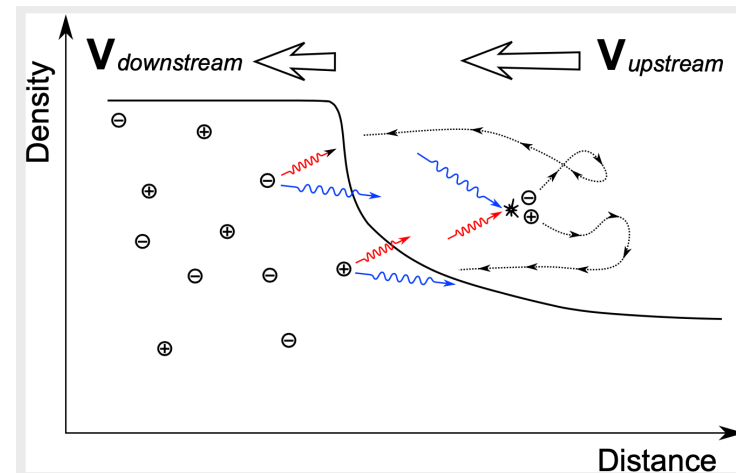


1900829A-H.E.S.S.

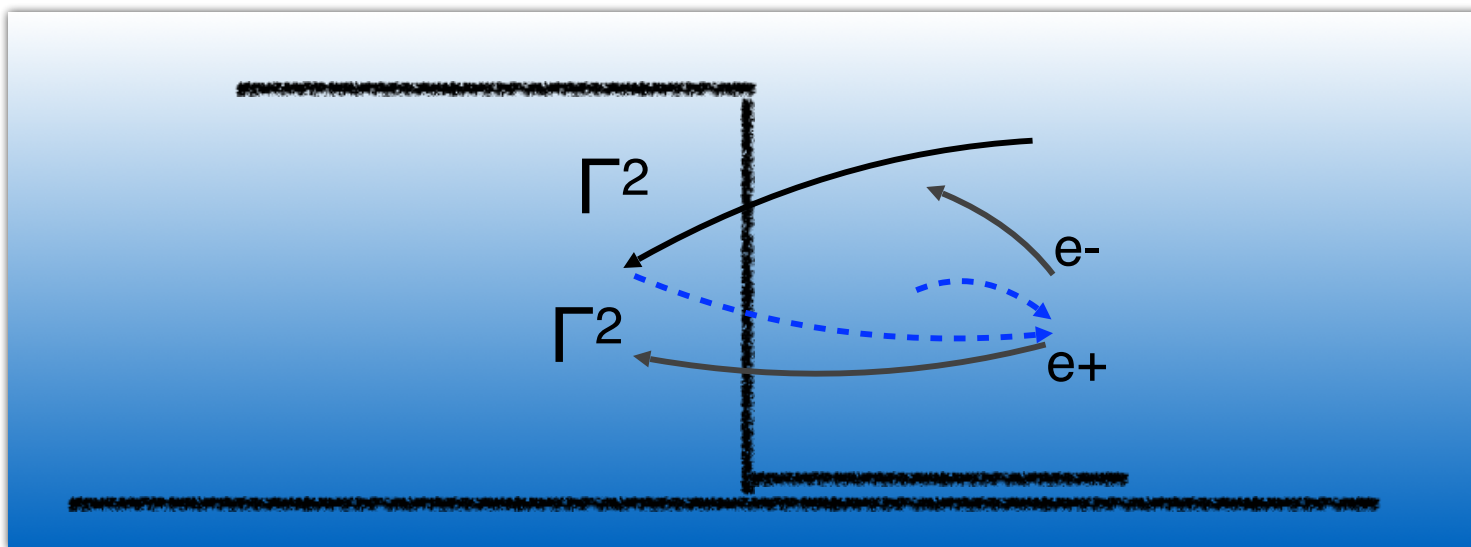
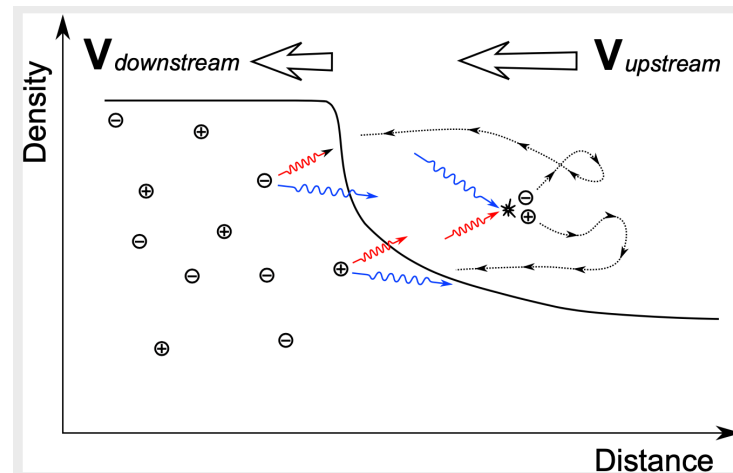


# The Pair Balance model

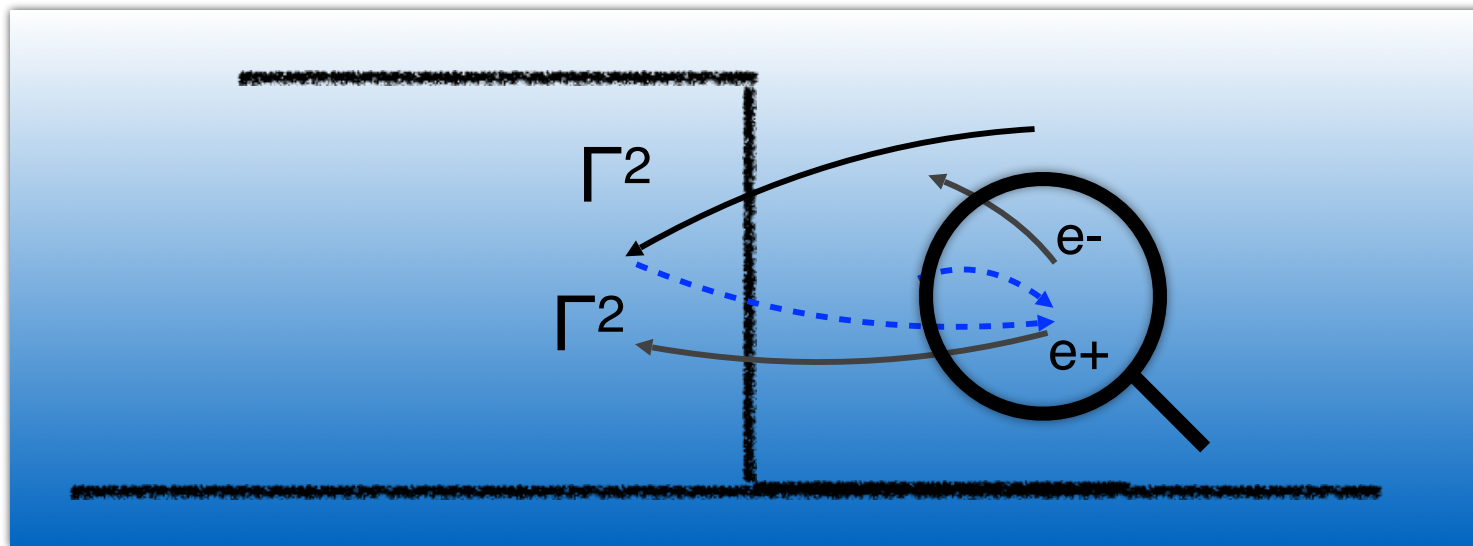
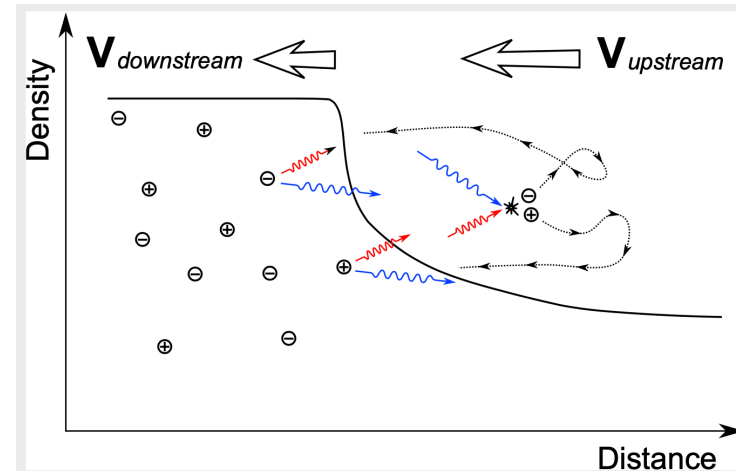
Derishev & TP 2016



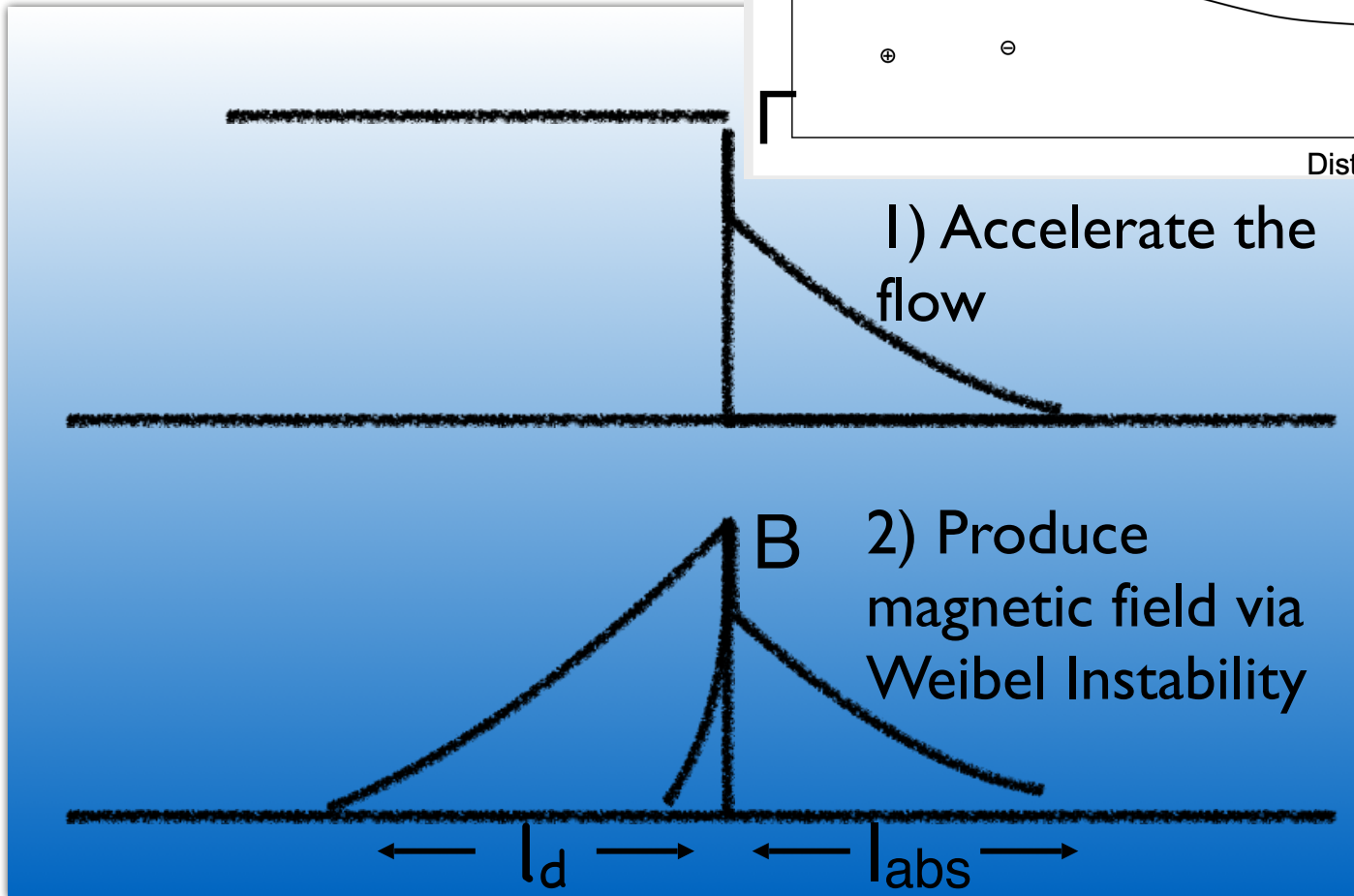
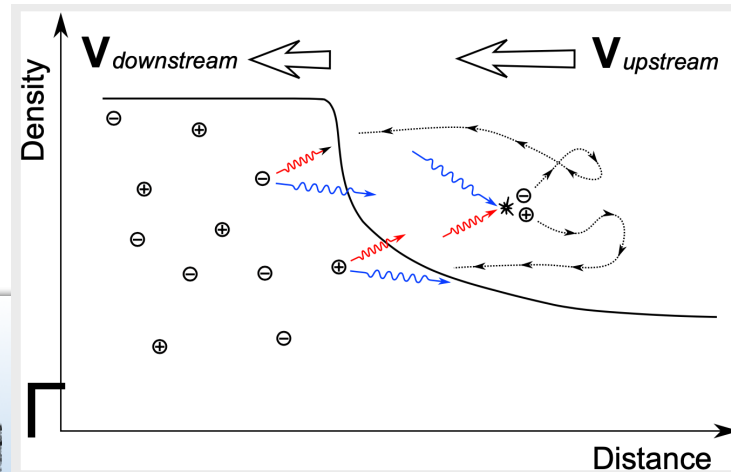
- 1) Pairs produced in the upstream
- 2) They are strongly accelerated once crossing the shock



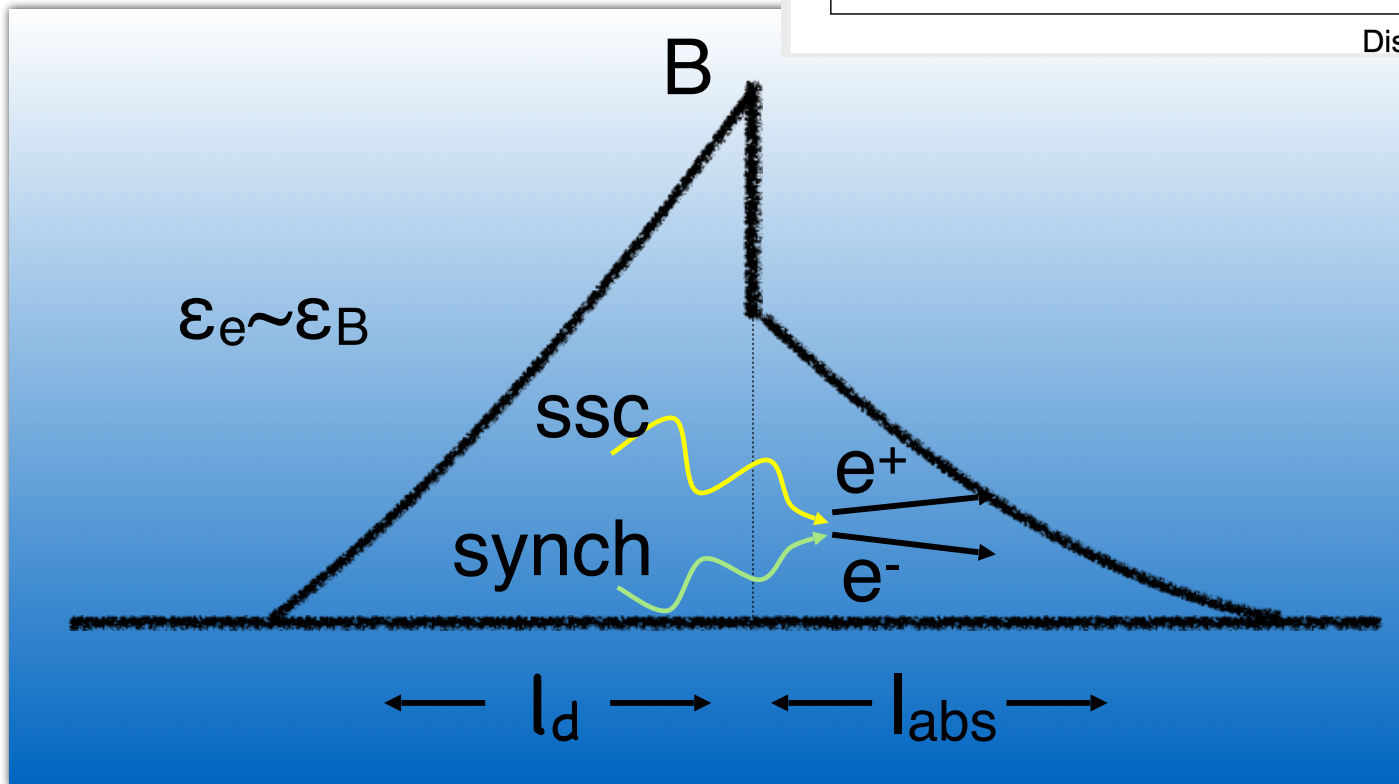
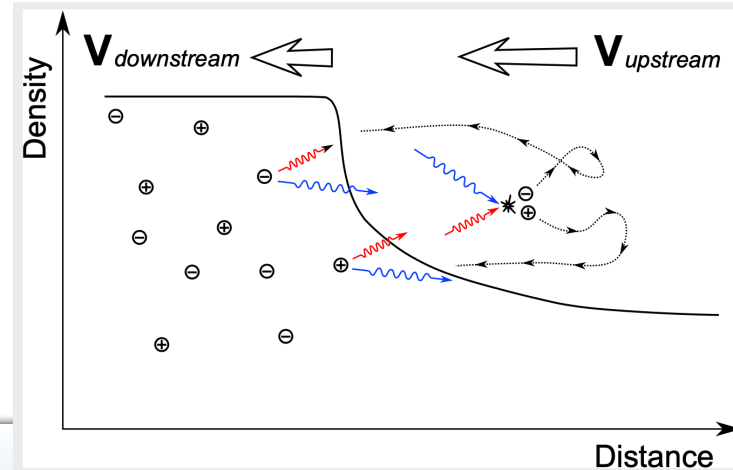
- 1) Accelerate the flow
  - 2) Produce magnetic field via Weibel Instability
- Instability



# Modified structure



Decaying magnetic field, in the downstream, accelerates particles



$$\epsilon_e \sim \epsilon_B$$

SSC

synch

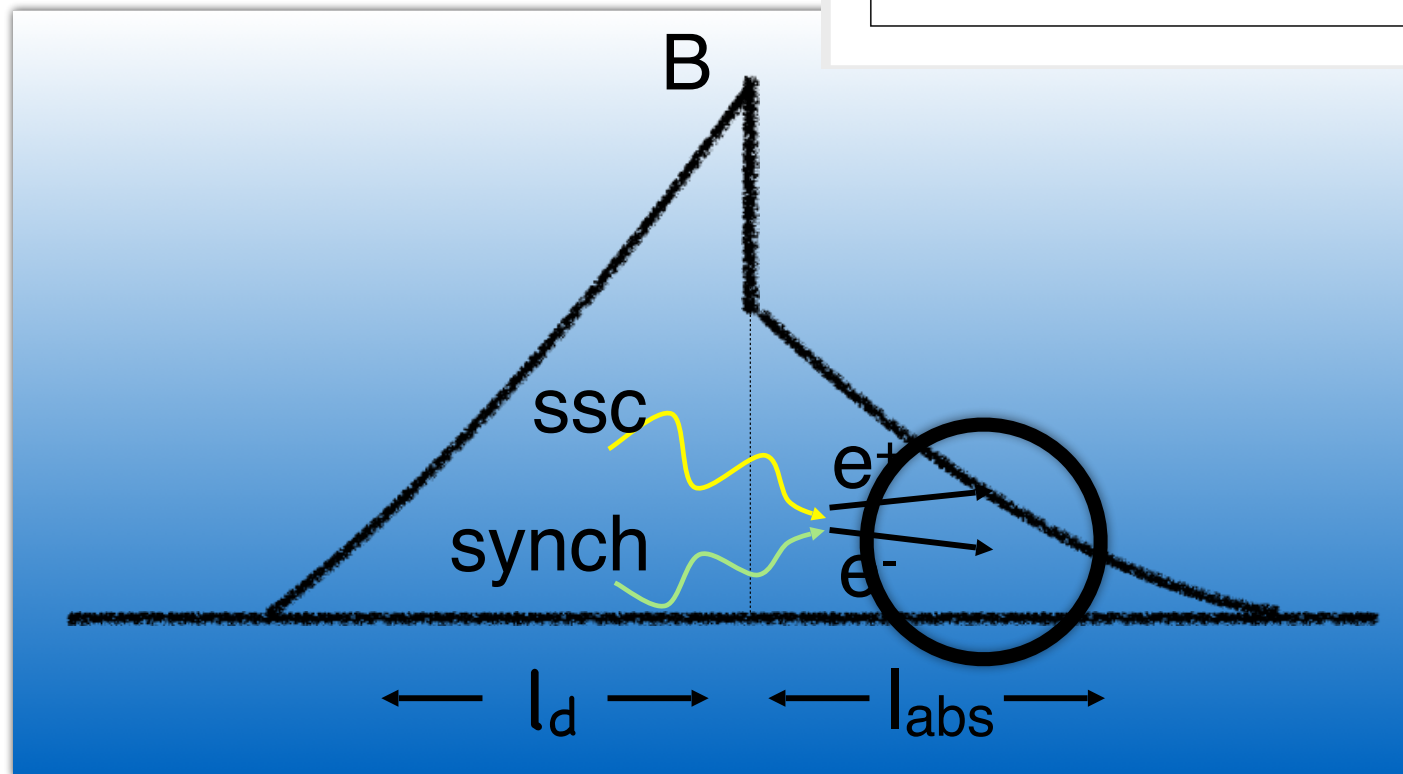
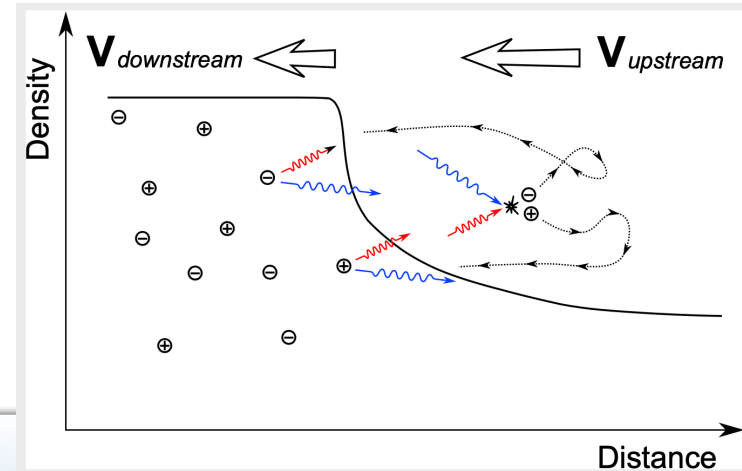
$e^+$

$e^-$

$l_d$

$l_{abs}$

Pairs from the upstream increase the multiplicity of the downstream





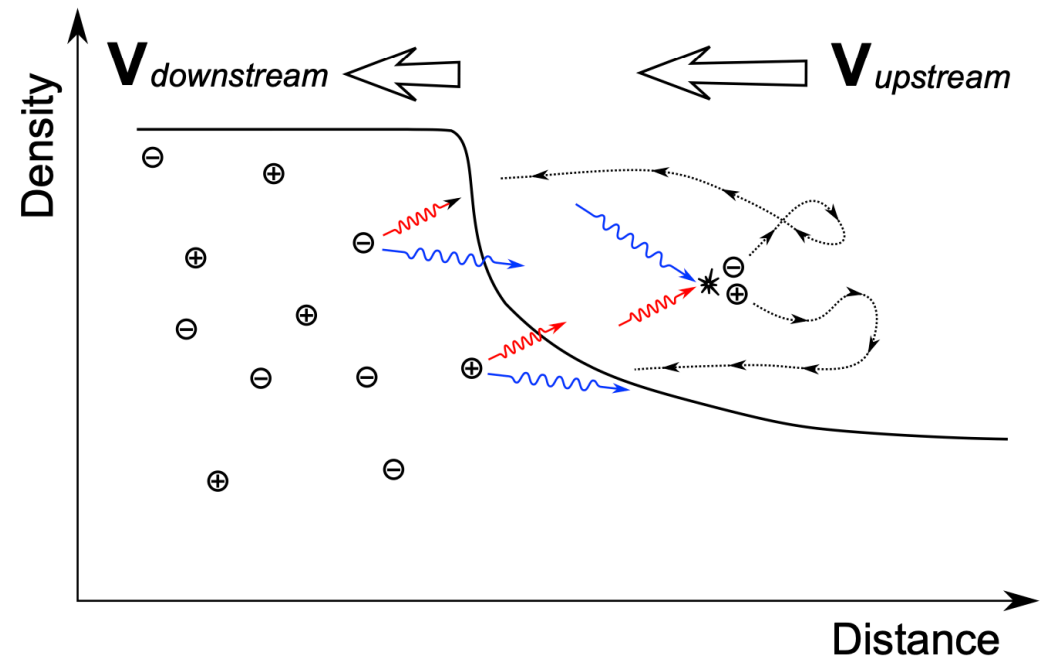
# Some basic features of the Pair-Balance model

Derishev & TP 2016

- Saturation at the Klein-Nishina limit  $\Rightarrow \gamma^3 B \approx B_{cr}$

$\Rightarrow \gamma_m \propto \Gamma$  doesn't hold

- $\tau_{\gamma\gamma} \lesssim 1$  for the IC photons



# One zone modeling

Sari, TP, Narayan 98

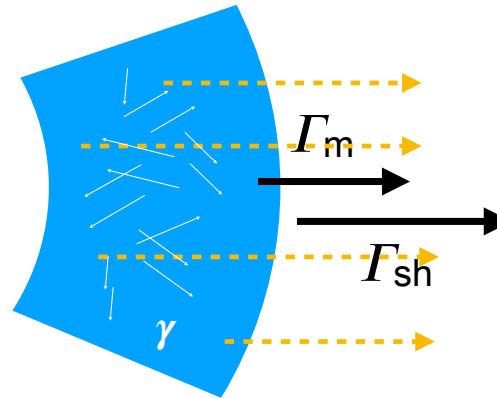
$$R = C_R \Gamma^2 c t_{\text{obs}} / (1 + z)$$

$$E_{\text{kin}} = C_E \Gamma^2 M c^2$$

$$t_{\text{eff}} = C_t \Gamma t_{\text{obs}} / (1 + z)$$

$$h\nu_{\text{obs}} = C_\Gamma \Gamma h\nu / (1 + z)$$

$$L = C_L \epsilon_r (1 + z) E_{\text{kin}} / t_{\text{obs}} ,$$



Reference	density profile	
	Wind	ISM
Sari et al. (1998), “SPN98 coefficients” hereafter		$C_R = 2, C_t = 1/\sqrt{2}, C_\Gamma = 1/\sqrt{2}, C_L = 17/12^*$
Panaitescu & Mészáros (1998b)	$C_R \approx 3.1$	$C_R \approx 6.5$
Nava et al. (2013)	$C_R = 4/5$	$C_R = 8/9$
Dai & Lu (1998)	$C_R = 4, C_t = 8\sqrt{2}/3$	$C_R = 8, C_t = 16\sqrt{2}/5$
Derishev & Piran (2019)	$C_R = 4, C_E = 1, C_t = 4, C_\Gamma = 1$	$C_R = 8, C_E = 1, C_t = 8, C_\Gamma = 1$
Current work (from Derishev 2021), “effective coefficients” hereafter	$C_R \approx 2.45, C_E = 2/9, C_t \approx 1.16,$ $C_\Gamma \approx 0.64, C_L = 9/8$	$C_R \approx 5.55, C_E = 6/17, C_t \approx 0.96,$ $C_\Gamma \approx 0.87, C_L = 17/16$

# One Zone Coefficients

$$R = C_R \Gamma^2 c t_{\text{obs}} / (1 + z)$$

$$E_{\text{kin}} = C_E \Gamma^2 M c^2$$

$$t_{\text{eff}} = C_t \Gamma t_{\text{obs}} / (1 + z)$$

$$h\nu_{\text{obs}} = C_\Gamma \Gamma h\nu / (1 + z)$$

$$L = C_L \epsilon_r (1 + z) E_{\text{kin}} / t_{\text{obs}}$$

Reference	density profile	
	Wind	ISM
Sari et al. (1998), “SPN98 coefficients” hereafter		$C_R = 2, C_t = 1/\sqrt{2}, C_\Gamma = 1/\sqrt{2}, C_L = 17/12^*$
Panaitescu & Mészáros (1998b)	$C_R \approx 3.1$	$C_R \approx 6.5$
Nava et al. (2013)	$C_R = 4/5$	$C_R = 8/9$
Dai & Lu (1998)	$C_R = 4, C_t = 8\sqrt{2}/3$	$C_R = 8, C_t = 16\sqrt{2}/5$
Derishev & Piran (2019)	$C_R = 4, C_E = 1, C_t = 4, C_\Gamma = 1$	$C_R = 8, C_E = 1, C_t = 8, C_\Gamma = 1$
Current work (from Derishev 2021), “effective coefficients” hereafter	$C_R \approx 2.45, C_E = 2/9, C_t \approx 1.16,$ $C_\Gamma \approx 0.64, C_L = 9/8$	$C_R \approx 5.55, C_E = 6/17, C_t \approx 0.96,$ $C_\Gamma \approx 0.87, C_L = 17/16$

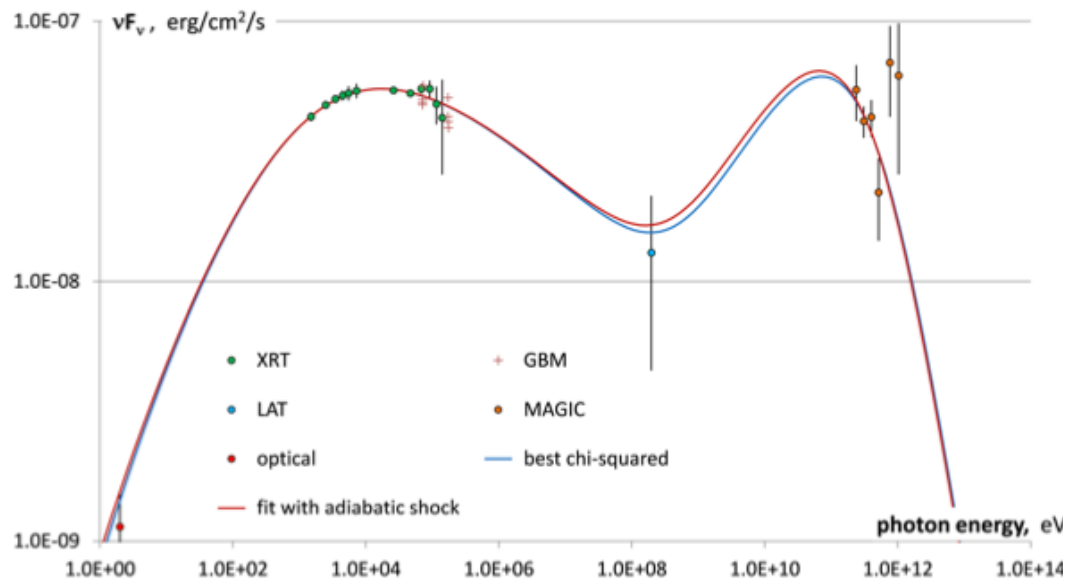
# First Guesses 190114c

- $\gamma\Gamma m_e c^2 > E_{IC} \Rightarrow \gamma\Gamma \approx 10^6$
- @ 70 sec and longer  $\Gamma$  cannot be too large  $\Rightarrow \gamma \gtrsim 10^4$
- $\Rightarrow$  Tev is Inverse Compton of X-rays (Consistent with a comparable X-ray luminosity)

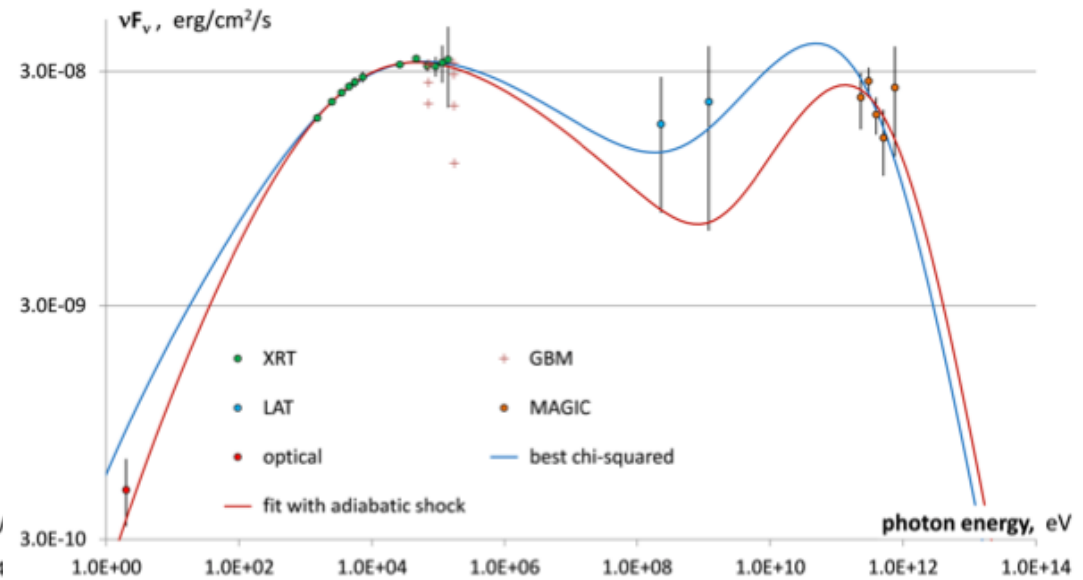
# Detailed modeling

(Derishev & TP 2021)

- Conditions at the emitting region are determined by  $\Gamma$ ,  $B$ ,  $\gamma_m$ ,  $\epsilon_e/\epsilon_B$



Early - 90 sec



late - 145 sec

# Best Fit Parameters

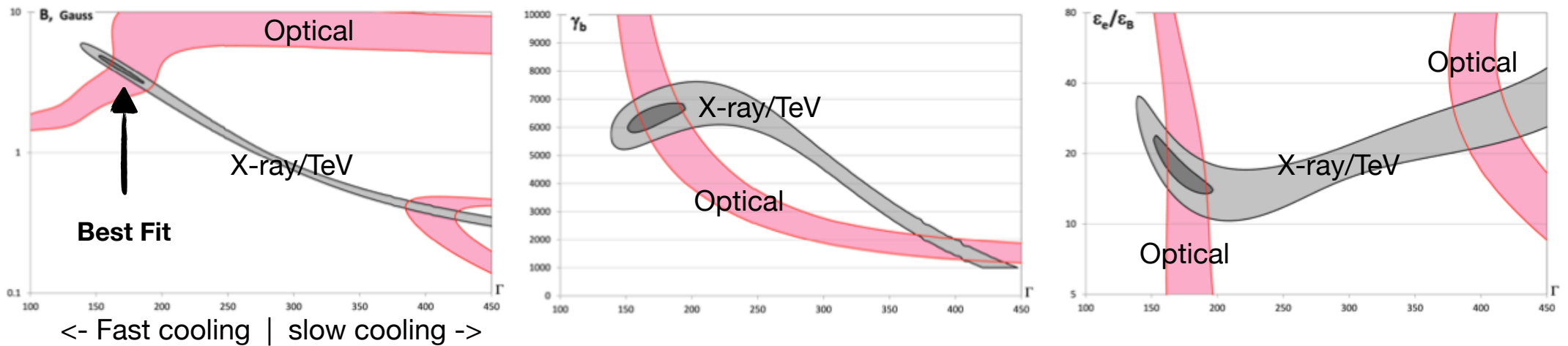
The fit didn't take into account the "pair balance" model however, the results are fully consistent with it and are inconsistent with standard afterglow modeling

parameter	$t_{\text{obs}} = 90 \text{ s}$	$t_{\text{obs}} = 145 \text{ s}$
$\Gamma$	161 (109) →	143 (91)
$B$	4.4 G (5.7 G)	2.0 G (3.1 G)
$\epsilon_e/\epsilon_B$	20 (21)	36 (41)
$\gamma_b$	6500 (5700) →	16700 (14400)
$p$	2.5	2.5
$E_{\text{kin}}$	$3 \times 10^{53} \text{ erg}$	$3 \times 10^{53} \text{ erg}$
$\epsilon_B$	0.0061 (0.0062) →	0.0027 (0.0026)
$\epsilon_e$	0.12 (0.13)	0.096 (0.107)
$\dot{M}$ (wind)	$1.4 \times 10^{-6} \frac{V_w}{3000 \text{ km/s}} M_{\odot}/\text{yr}$	$1.4 \times 10^{-6} \frac{V_w}{3000 \text{ km/s}} M_{\odot}/\text{yr}$
$n$ (ISM)	$2 \text{ cm}^{-3}$	$2 \text{ cm}^{-3}$

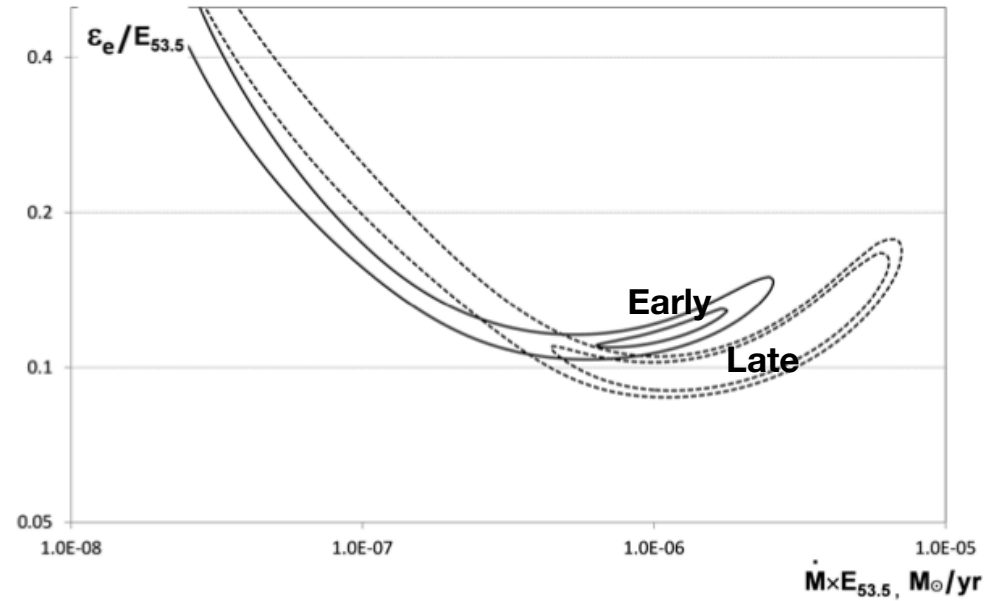
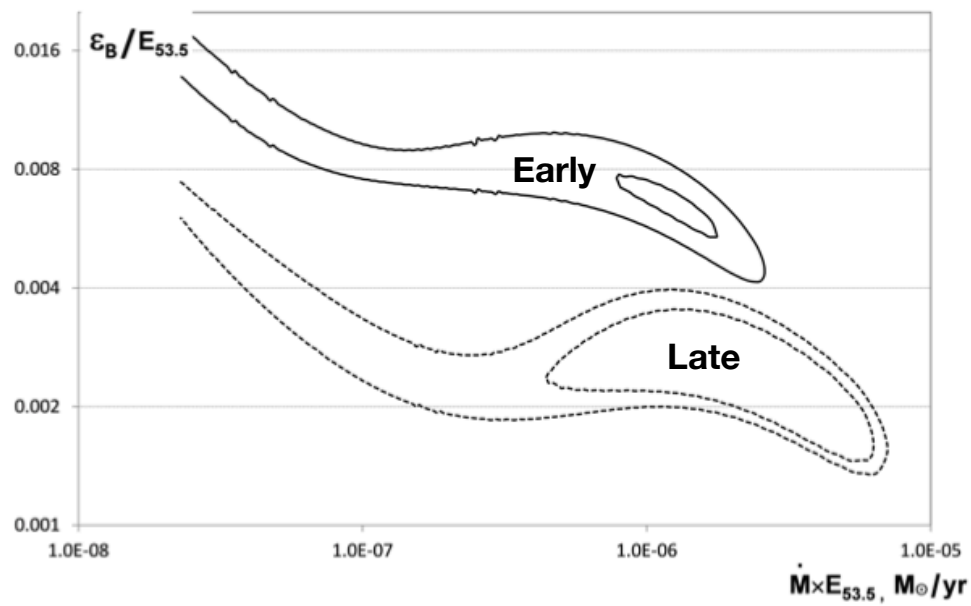
- Fast Cooling
- On the edge of KN regime
- $\gamma^3 B = (1.2 - 9) 10^{12}$   
 $\gamma_m \propto \Gamma$  doesn't hold
- $\tau_{\gamma\gamma} \approx 1$  for the IC photons (25% of IC power is self absorbed)
- $\epsilon_B = 0.006 \rightarrow 0.003$  (Varies)
- Somewhat surprisingly large  $\Gamma$  (large energy, low external density)

# Detailed modeling

- Conditions at the emitting region are determined by  $\Gamma$ ,  $B$ ,  $\gamma_m$ ,  $\epsilon_e/\epsilon_B$

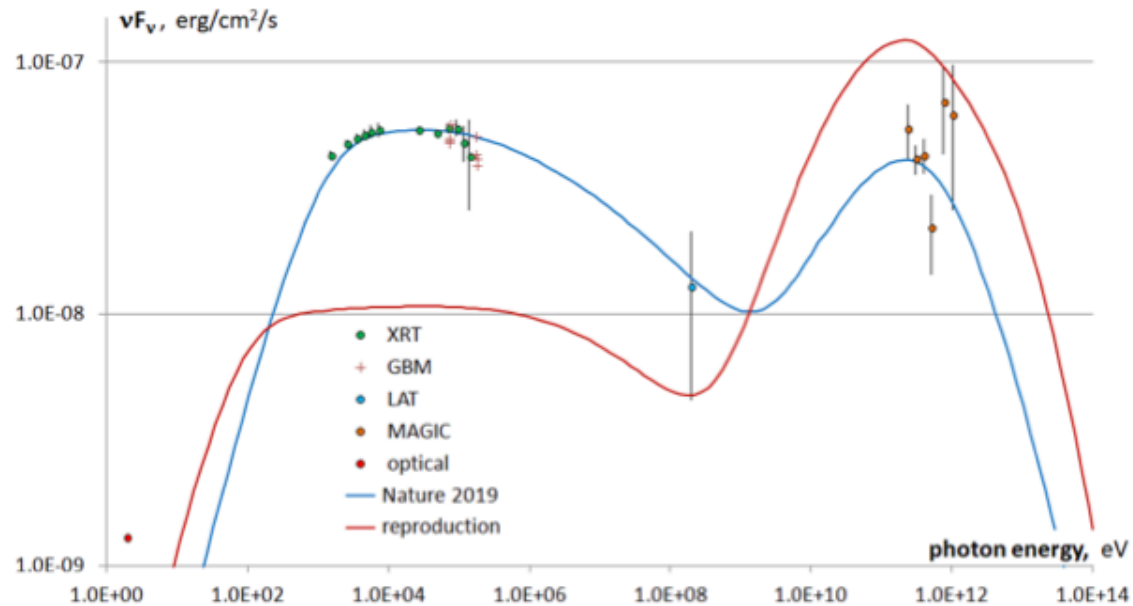


# $\epsilon_B$ must vary with time



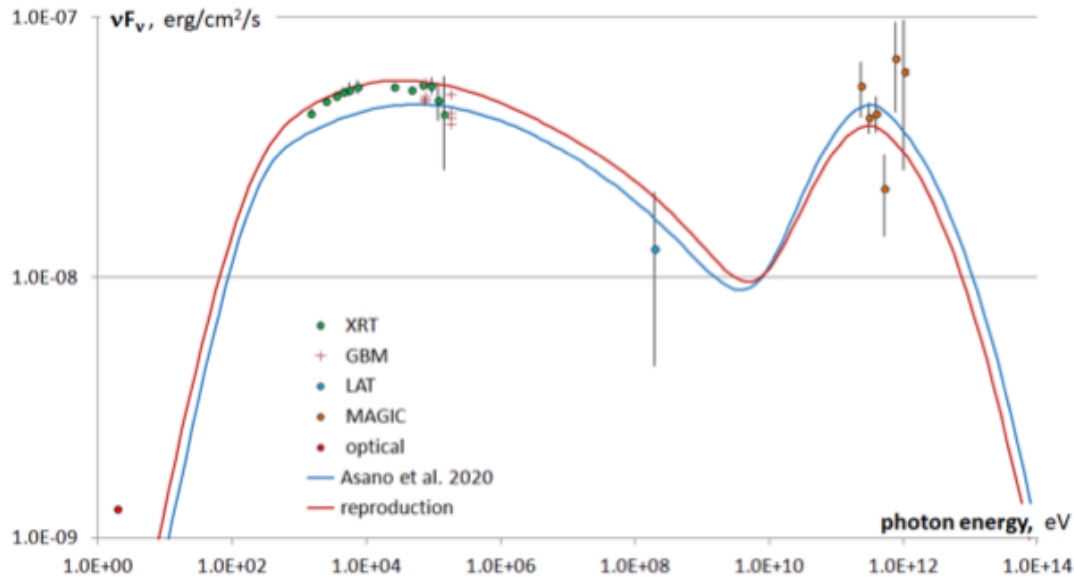


# Comparison with other work

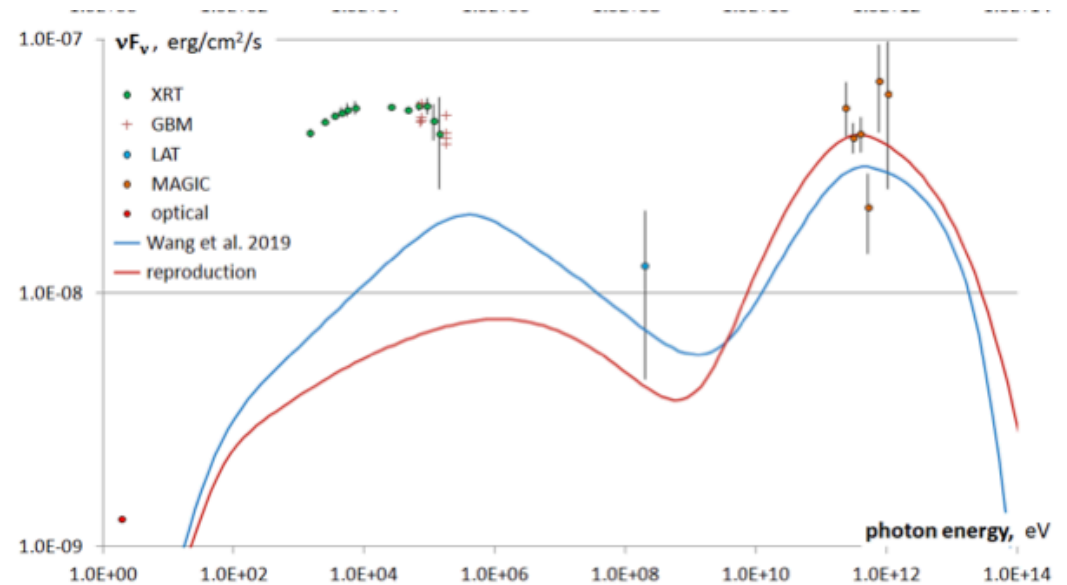


Magic 2019 ???

# Comparison with other work



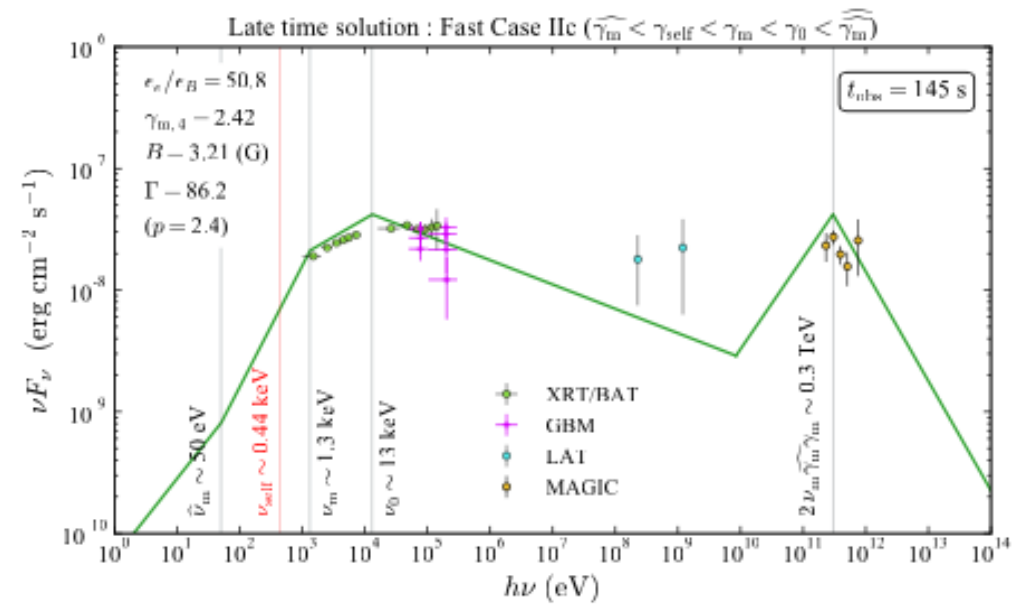
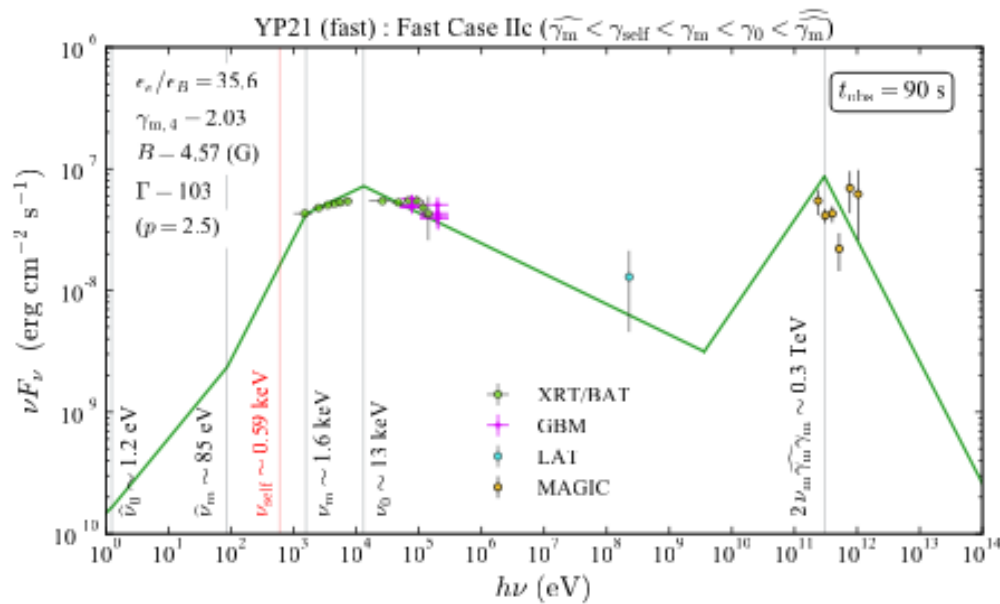
Asano & Murase 2020



Wang et al., 2019

# Analytic

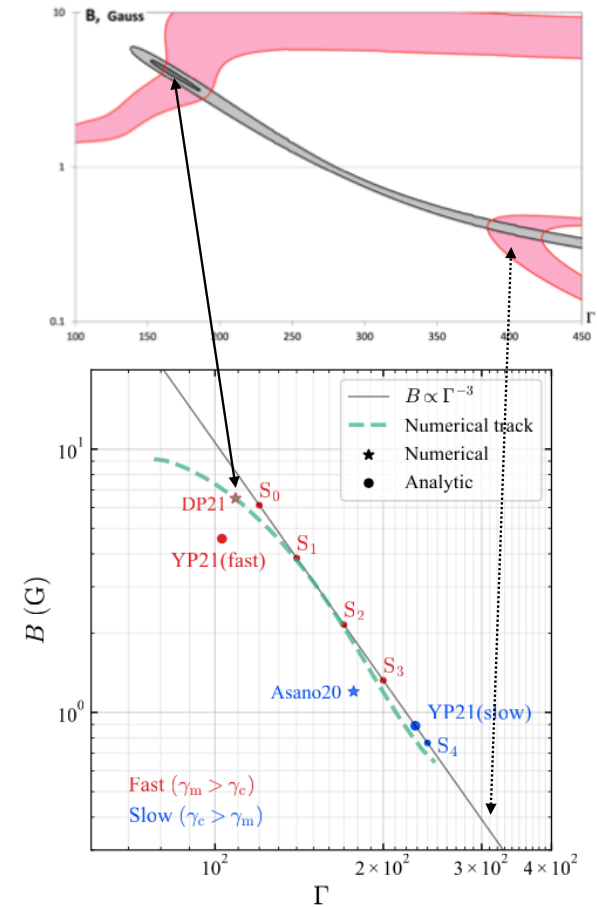
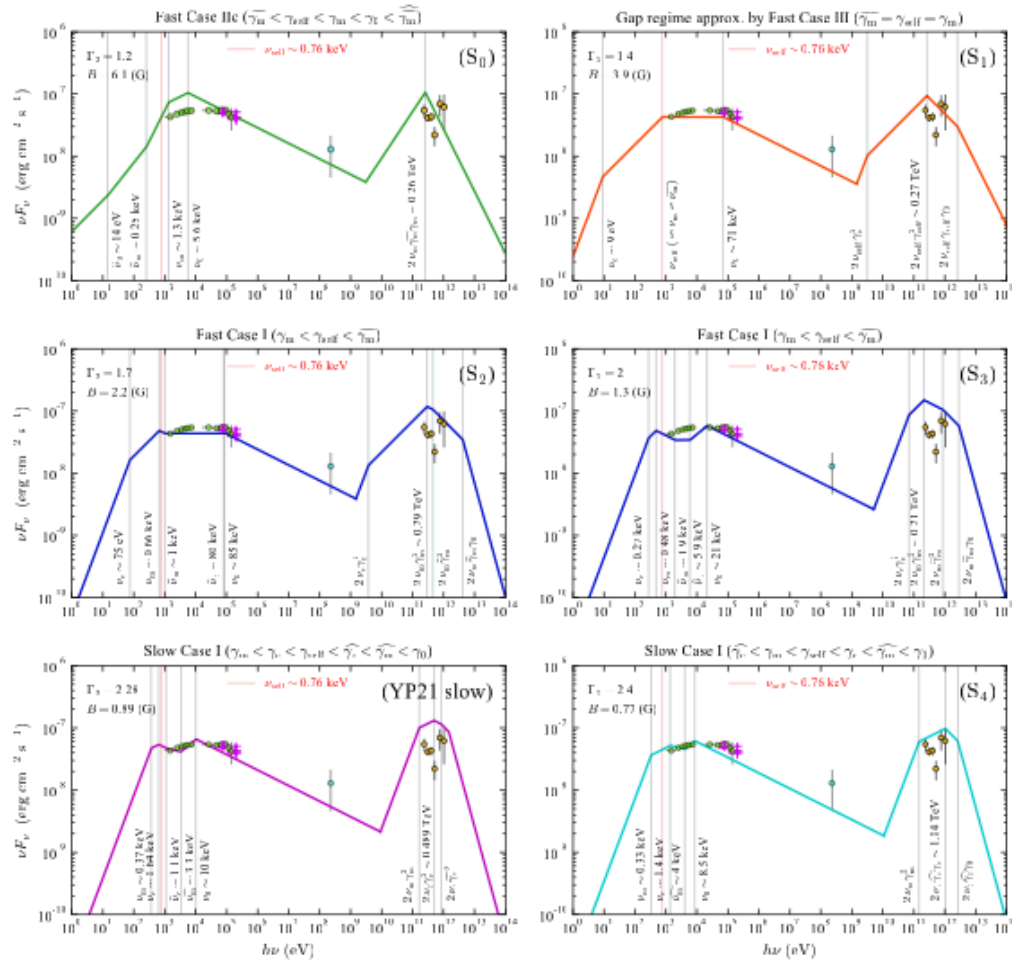
Yamasaki & Piran 2021 following Nakar et al., 2009\*



\* Sharp threshold for KN effect must be modified to start at  $\sim 100$  keV

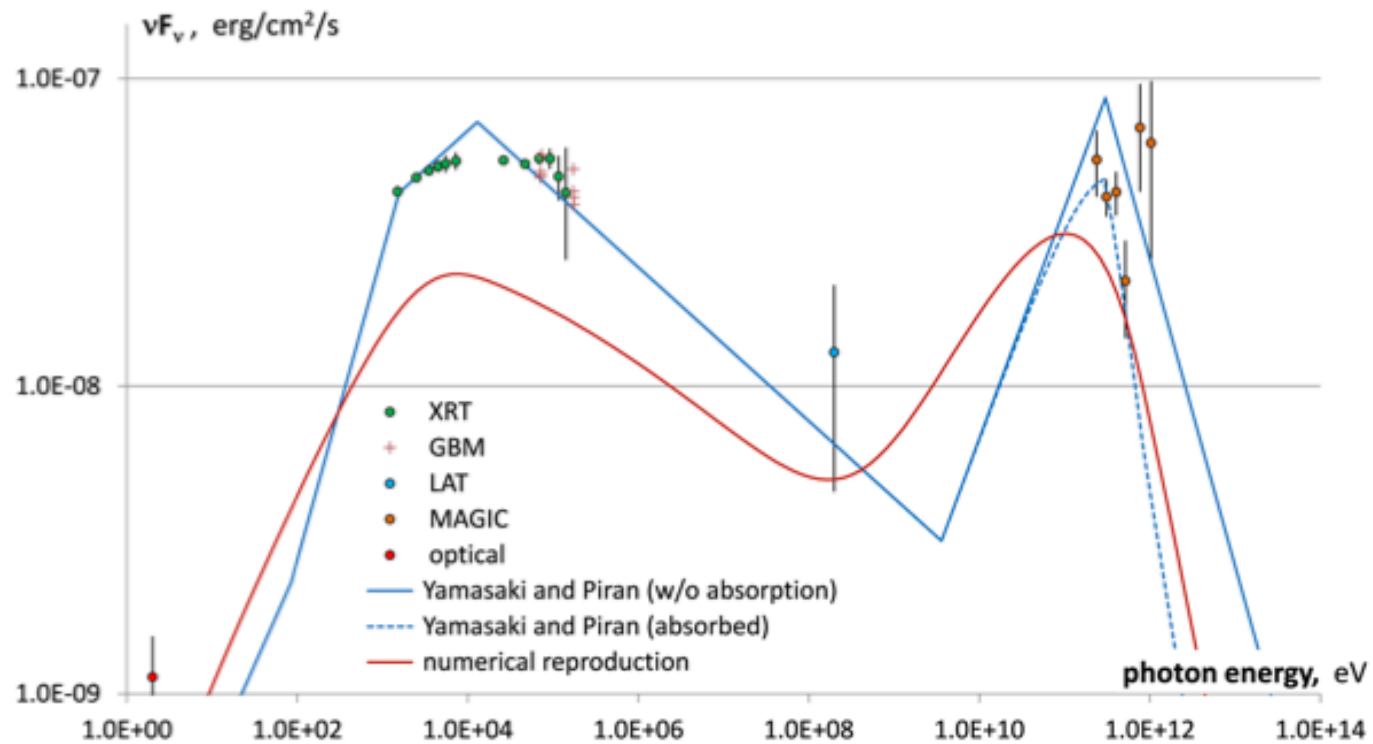
# Analytic

Yamasaki & Piran 2021 following Nakar et al., 2009

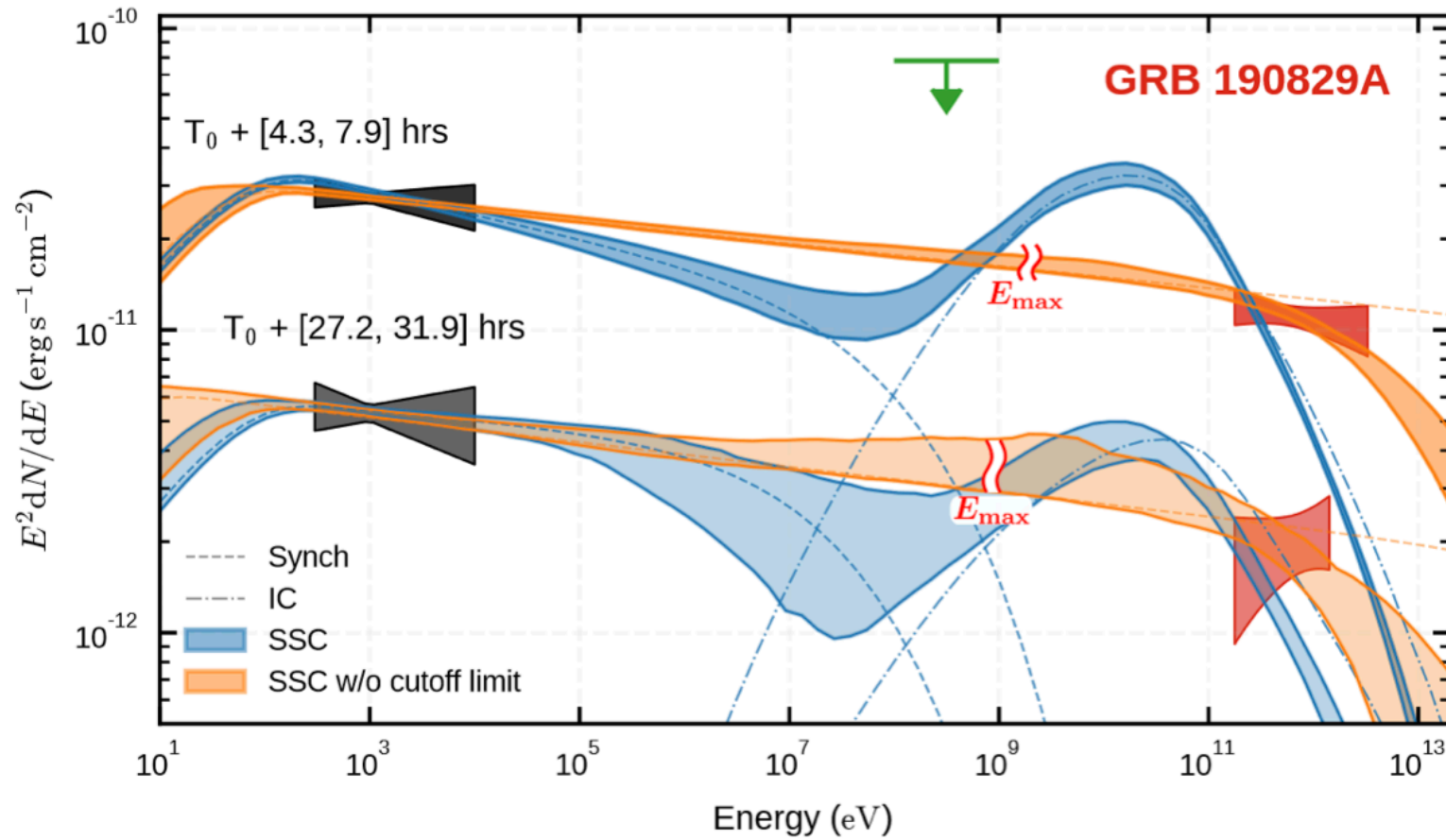


# Analytic

Yamasaki & Piran 2021 following Nakar et al., 2009

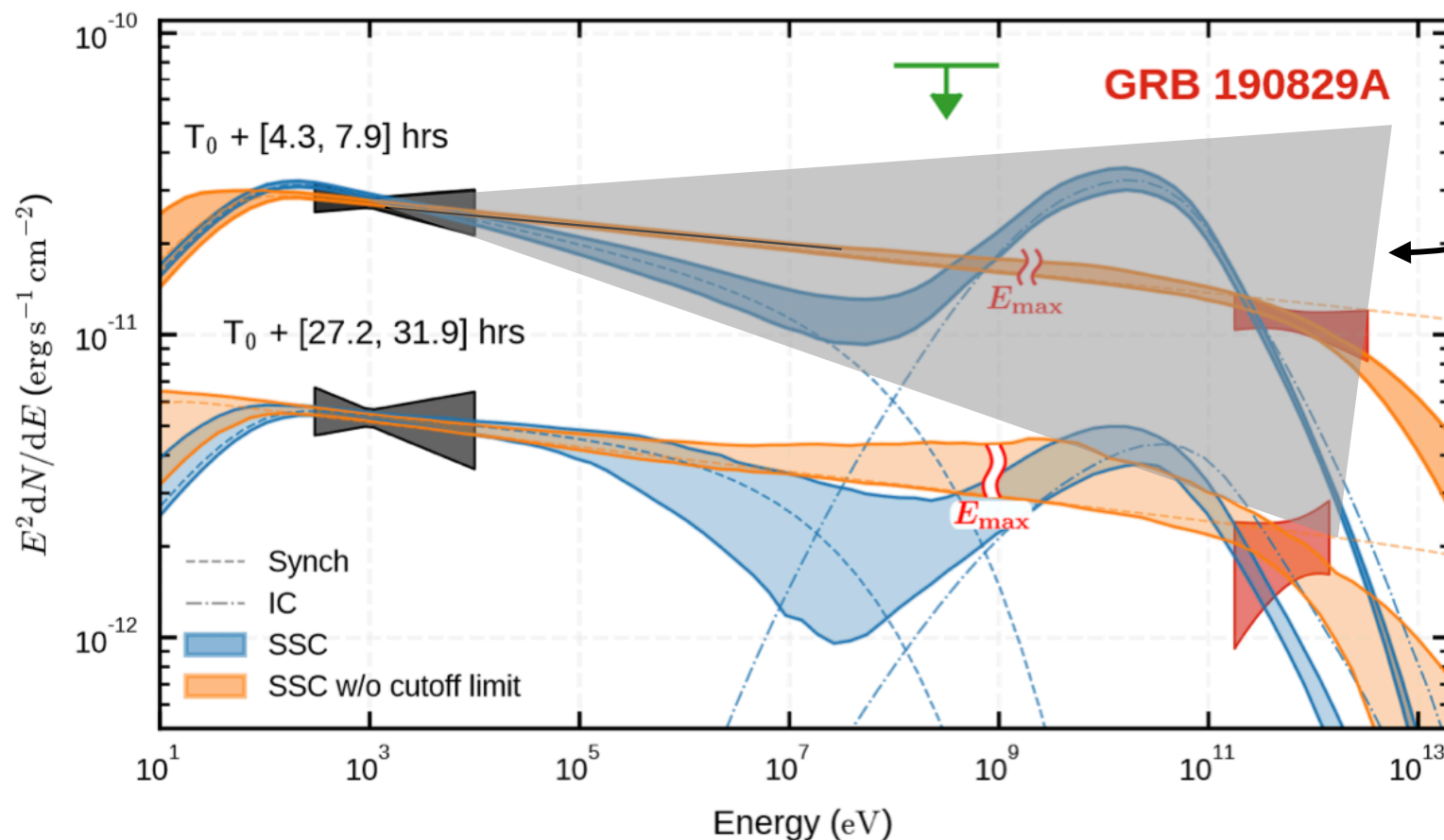


# 190829A

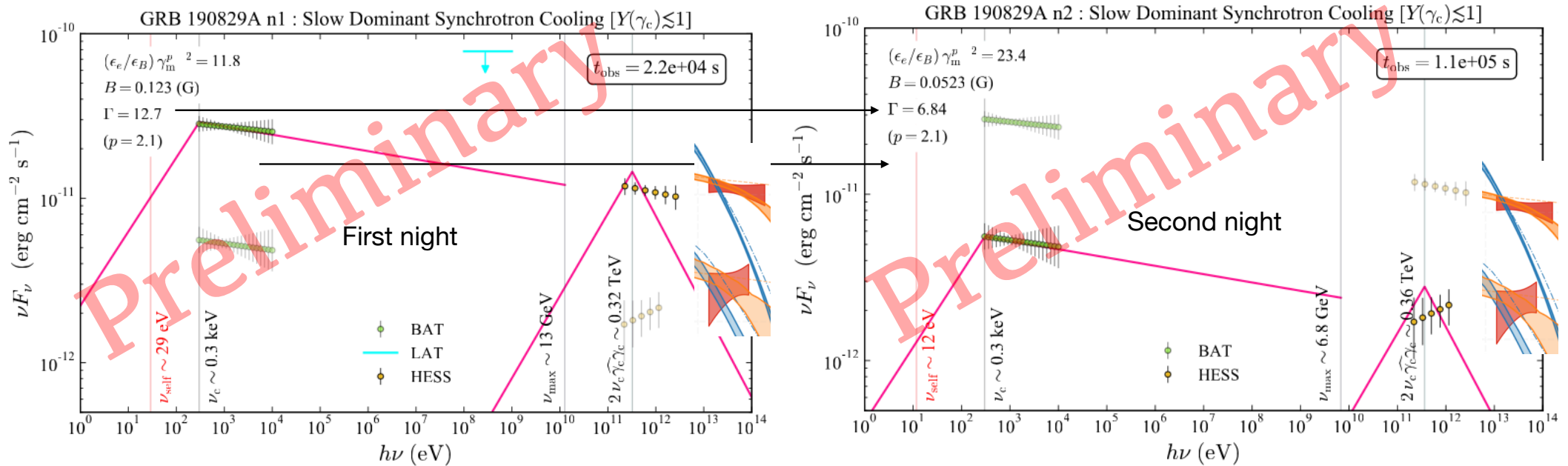


# A strange claim of a single power-law fit

With the error bars of the X-ray slope everything can fit a single power law to the TeV



# 190829A analytic modeling



Parameter	GRB 190829A	
	$t_{\text{obs}} = 22 \text{ ks}$	$t_{\text{obs}} = 110 \text{ ks}$
$\Gamma$	12.7	6.84
$B$	0.123 G	0.0523 G
$\gamma_m$	$< 1.3 \times 10^5$	$< 2.8 \times 10^5$
$\epsilon_e/\epsilon_B$	$> 3.7$	$> 6.9$
$p$	2.1	
Spectral regime	Slow DSC	

$\bullet \gamma^3 B \approx 10^{14}$

Yamasaki & TP 2021



# Summary

- TeV observations of both 190114c and 190829A seems to require modification of the simple afterglow model.
- A model independent fits for both bursts lead to parameters and evolutionary behavior that are (surprisingly) consistent with the Pair Balance model.

**Thanks for the attention**